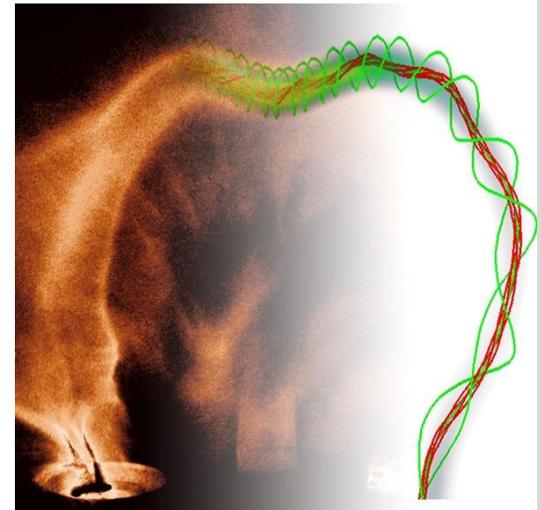


RUHR-UNIVERSITÄT BOCHUM

The FlareLab-Experiment: Laboratory Simulation of Arched Solar Prominences

Henning Soltwisch

AG Laser- und
Plasmaphysik 



FlareLab:

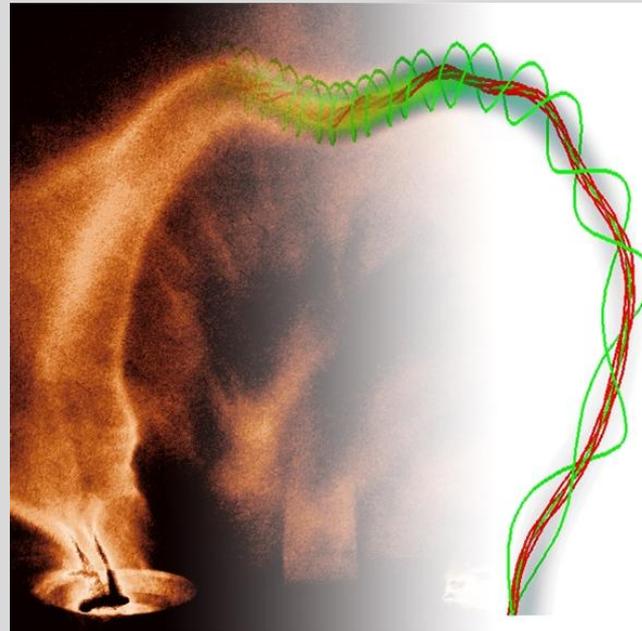
collaborative effort between laboratory investigations and numerical simulations of structures similar to solar prominences

Felix Mackel
Sascha Ridder
Jan Tenfelde
Philipp Kempkes^{a)}
Henning Soltwisch

*Institut für
Experimentalphysik*

AG Laser- und
Plasmaphysik und

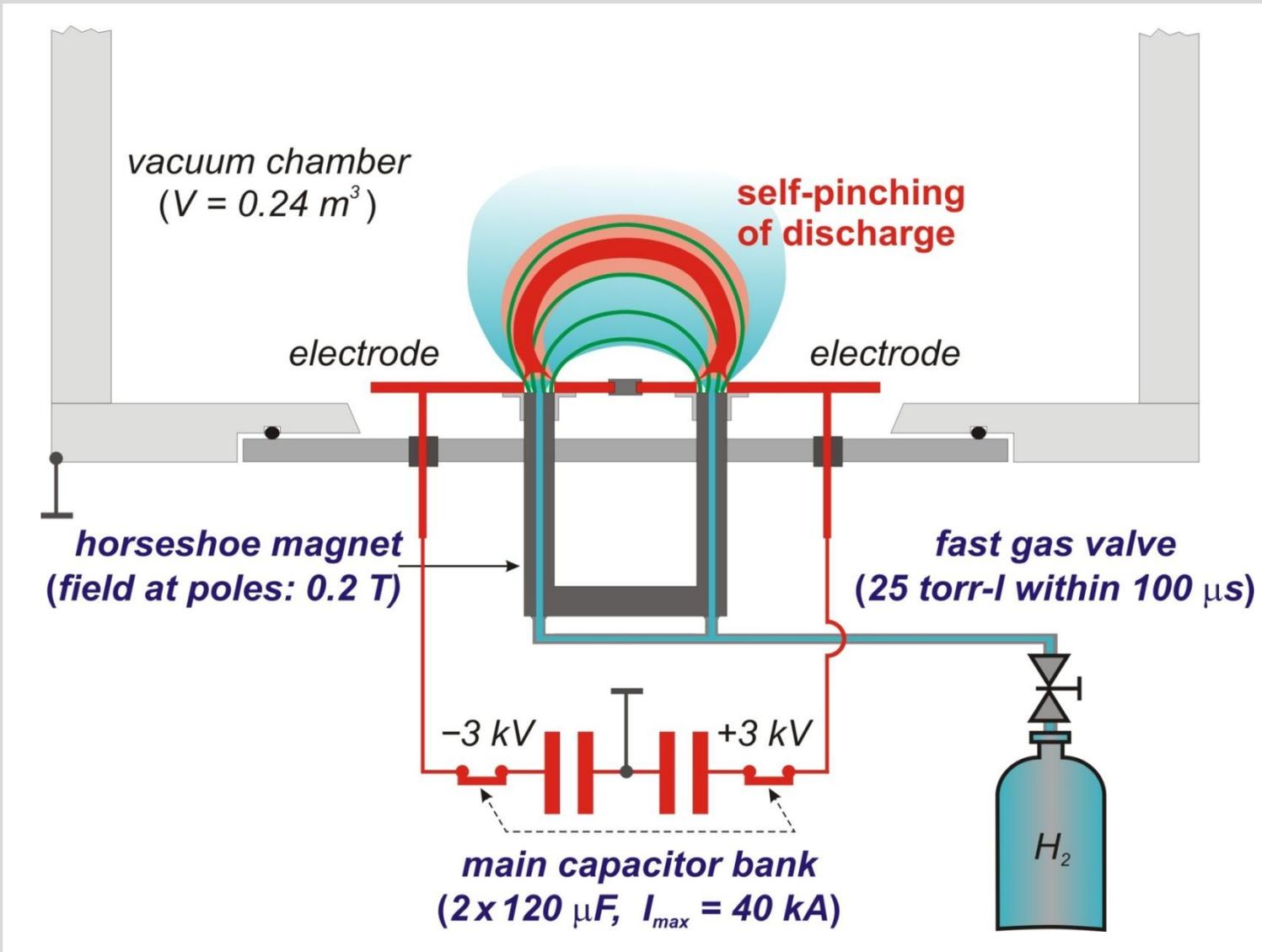
^{a)} MPI and University
Greifswald, Germany

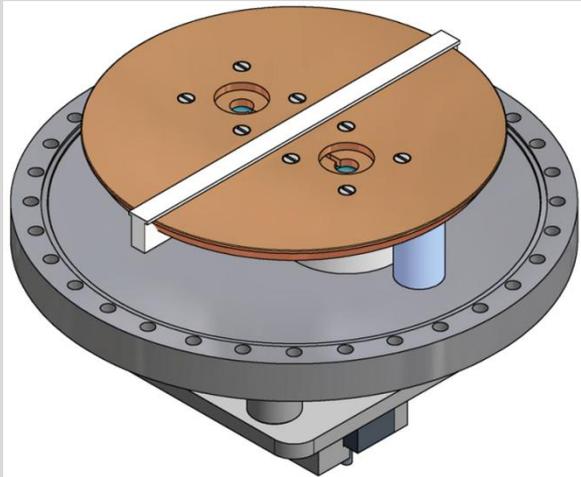


Thomas Tacke
Jürgen Dreher
Rainer Grauer

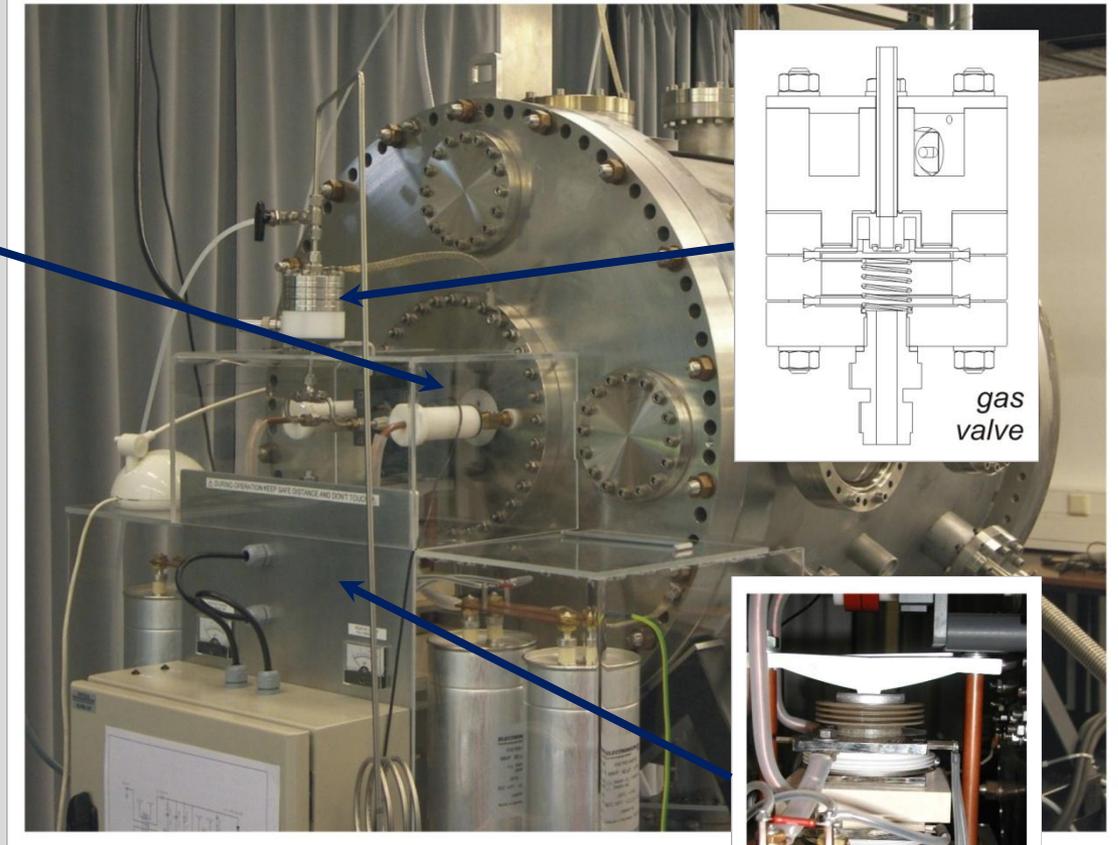
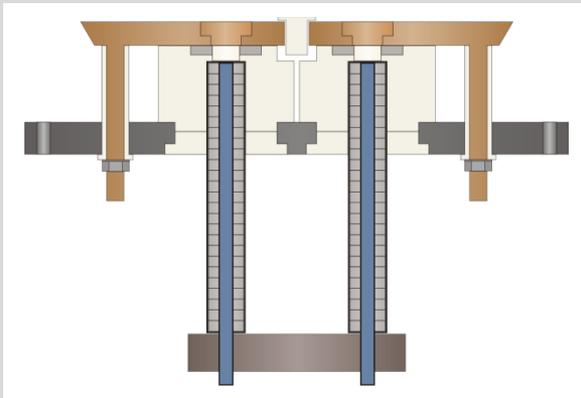
*Institut für
Theoretische Physik*

Lehrstuhl für
Plasma-, Laser-
und Atomphysik

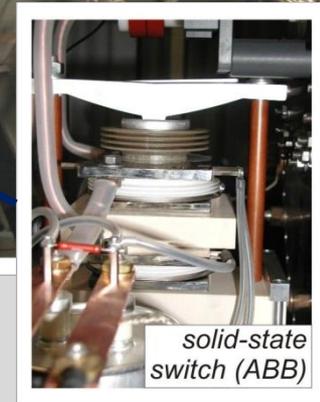
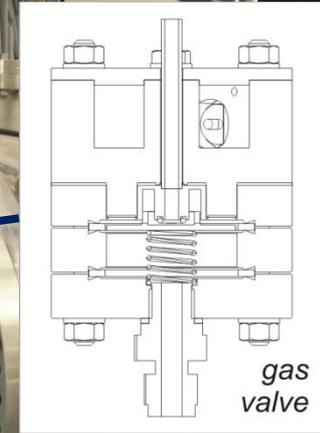


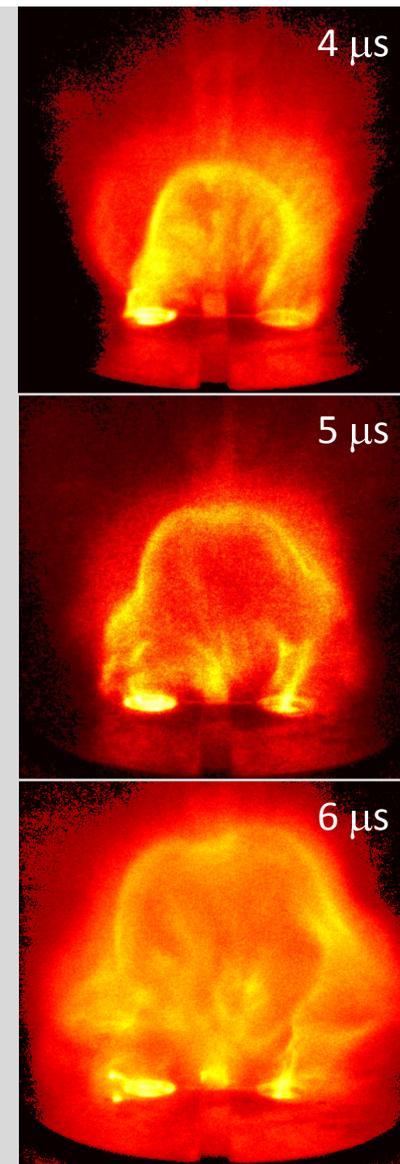
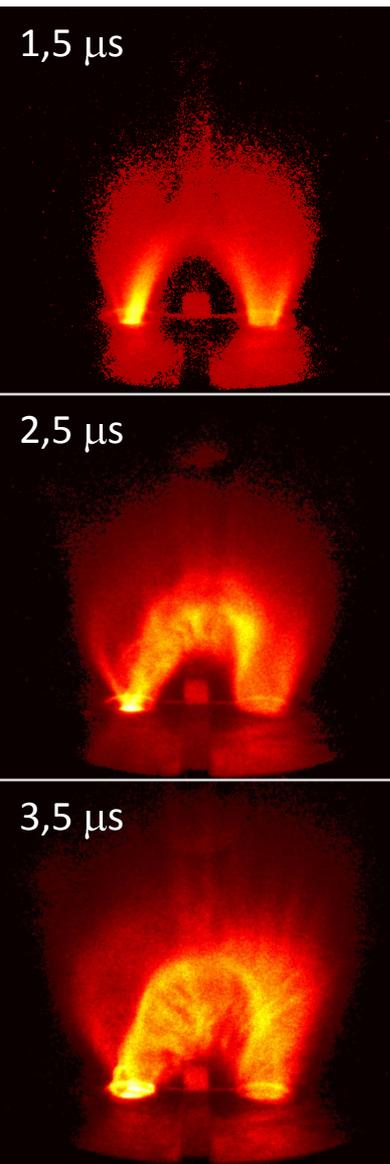


electrodes, gas inlet,
horseshoe magnet

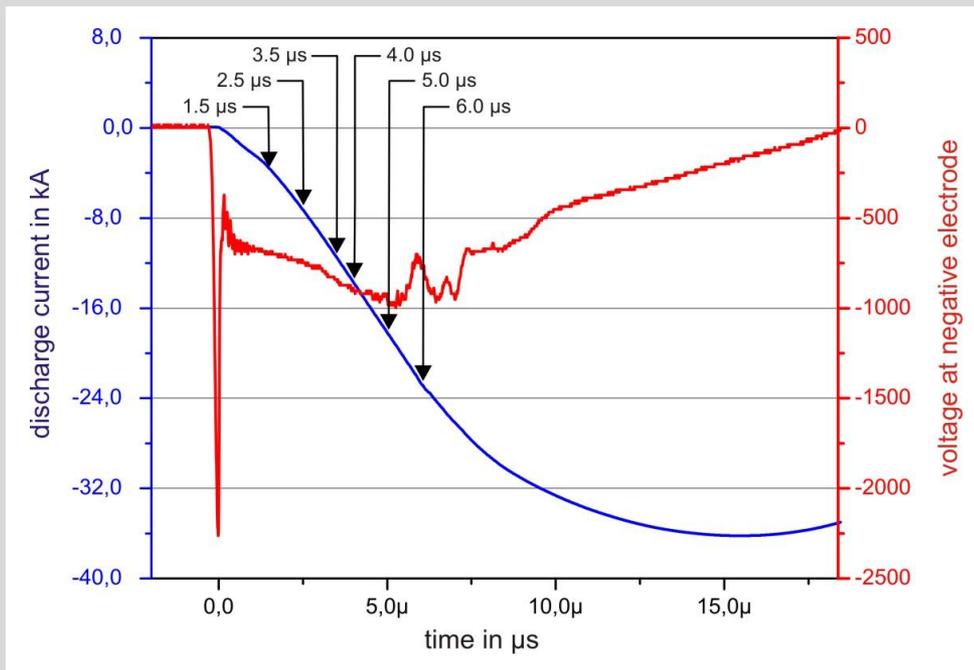


discharge chamber



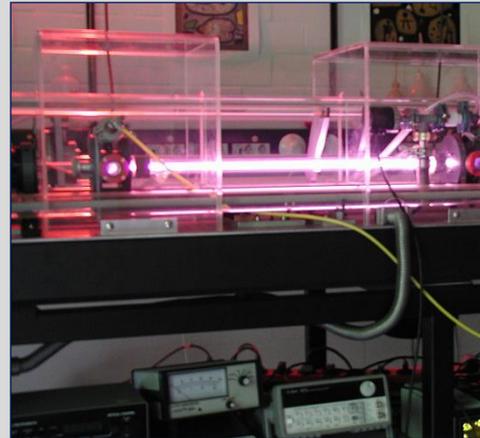


discharge current (**blue**) and
voltage at neg. electrode (**red**)

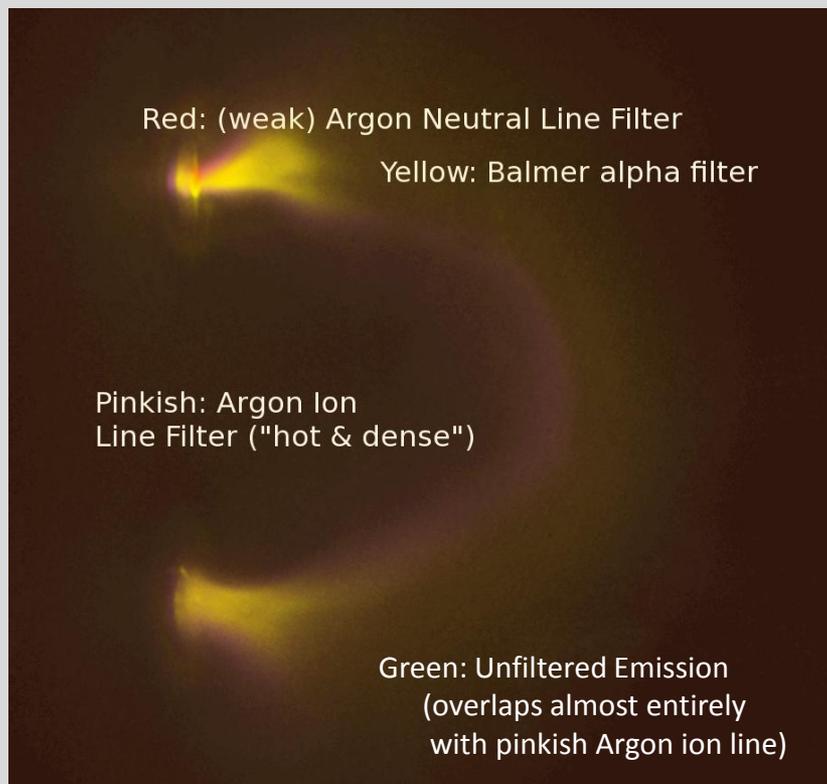


(exposure times: 50 ns)

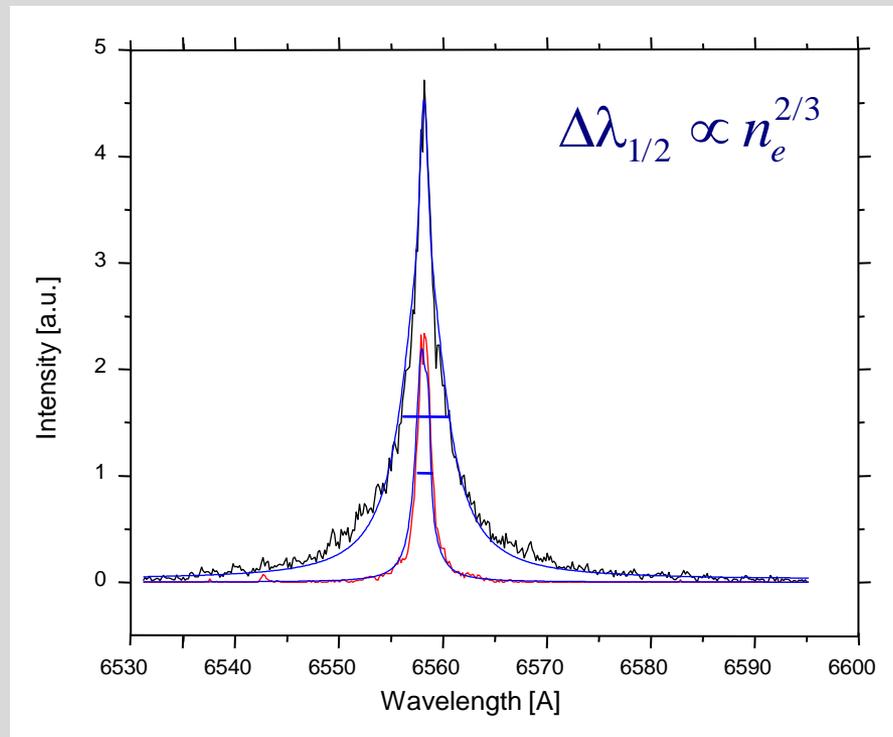
- **intensified CCD camera**
(discharge structure and dynamics)
- **HV probes and Rogowski coils**
(voltage and current measurements)
- **B-dot probes**
(local magnetic field measurements)
- **electrostatic triple probes**
(local measurements of n_e and T_e)
- **standard spectroscopy (visible)**
(plasma composition; line broadening)
- **XUV diode (metal filter foils)**
(generation of radiation by fast particles)
- **CO₂ laser interferometer**
(movable probe beam; line-averaged n_e)



- spectroscopic observations → plasma composition and spatial structure, n_e via Stark broadening

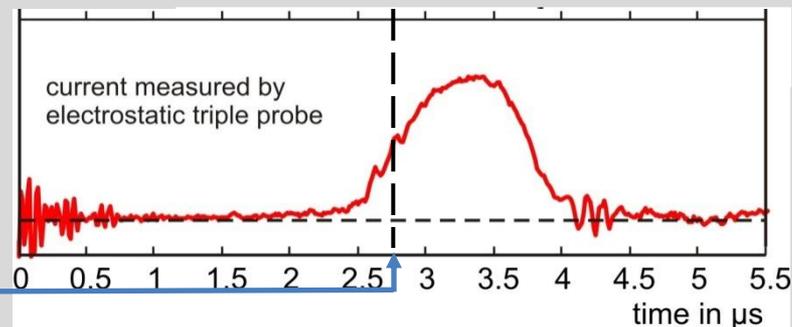
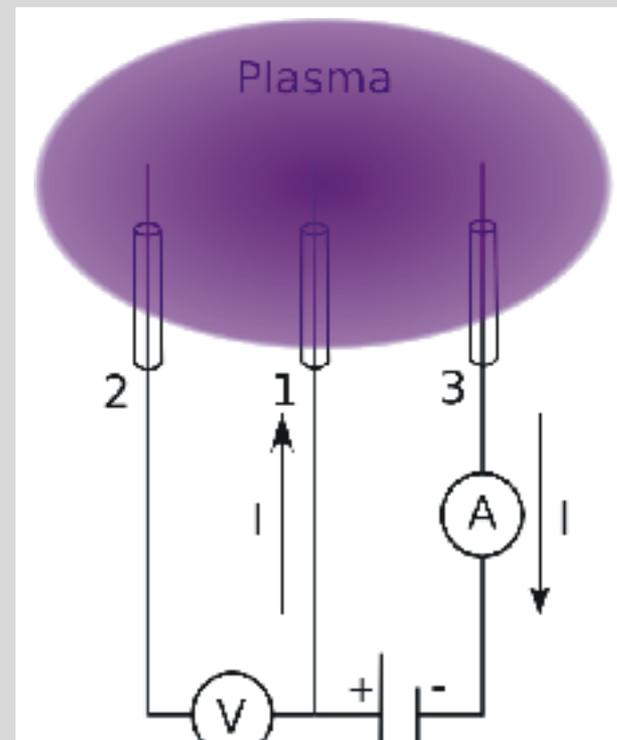
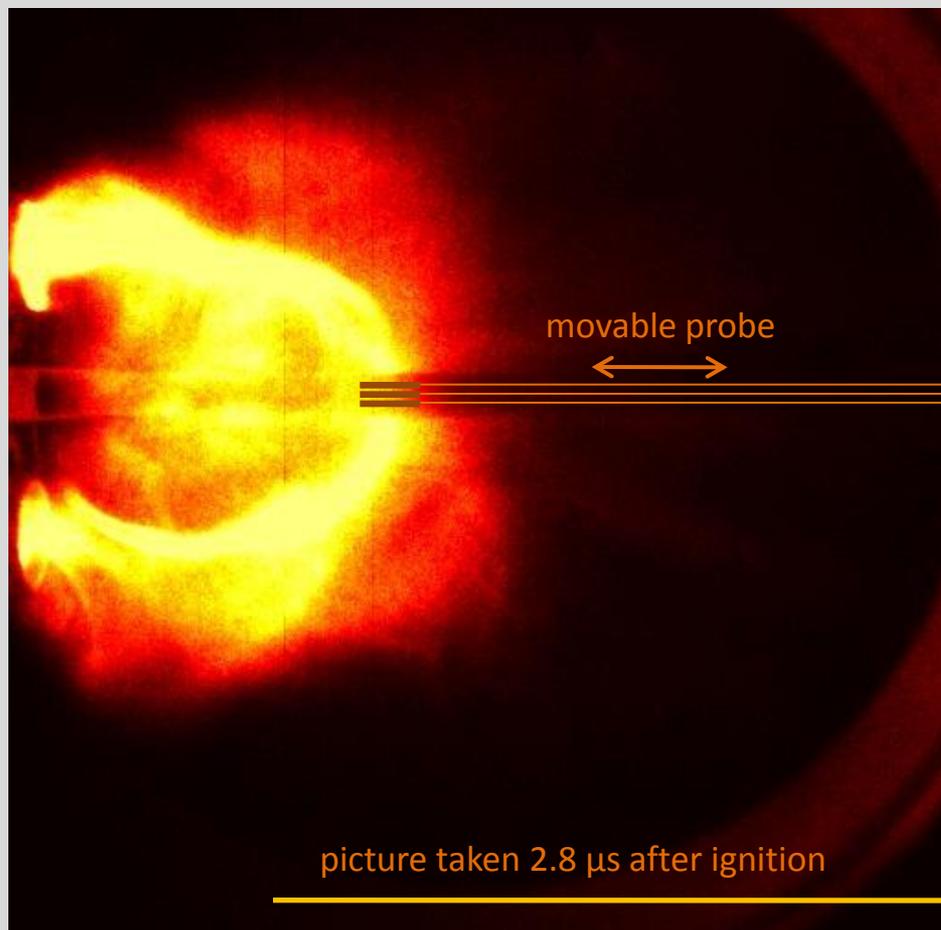


Superposition of discharge images using narrow-band optical filters

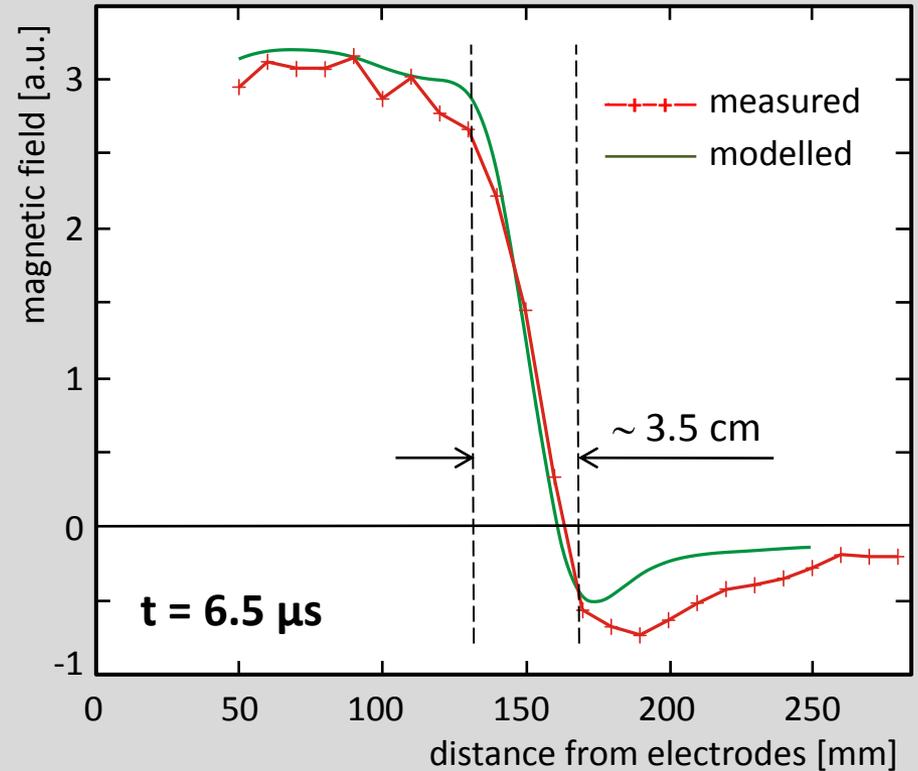
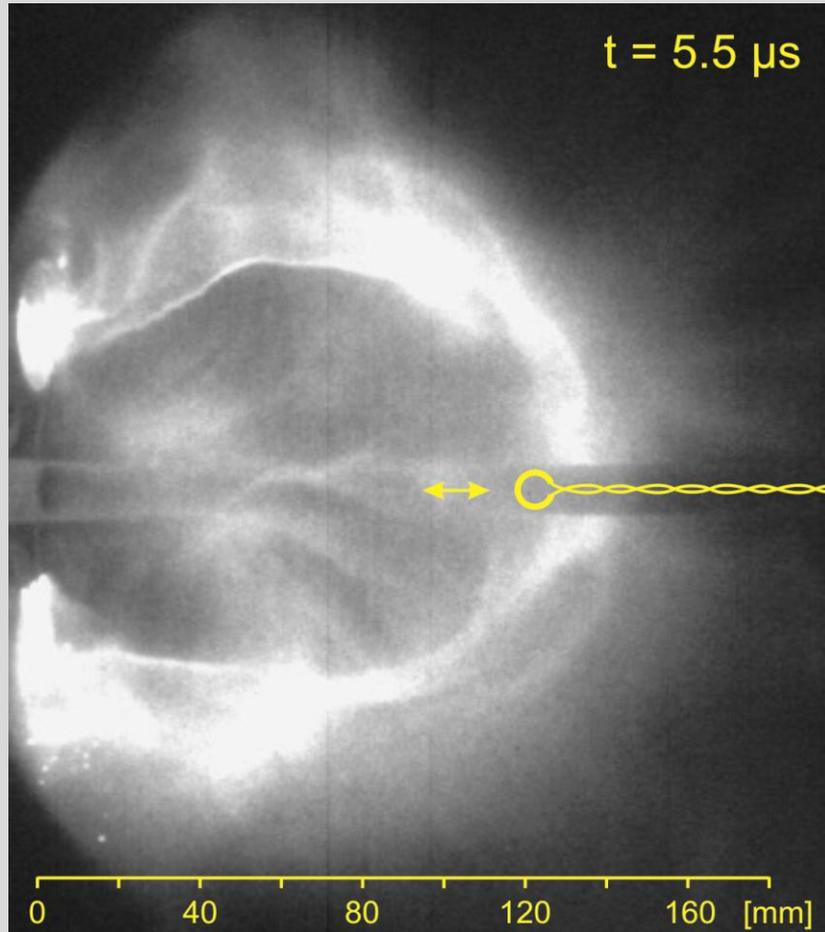


Hydrogen H_{α} line profiles emitted from plasmas of different electron densities

➤ movable electrostatic triple probe → n_e, T_e

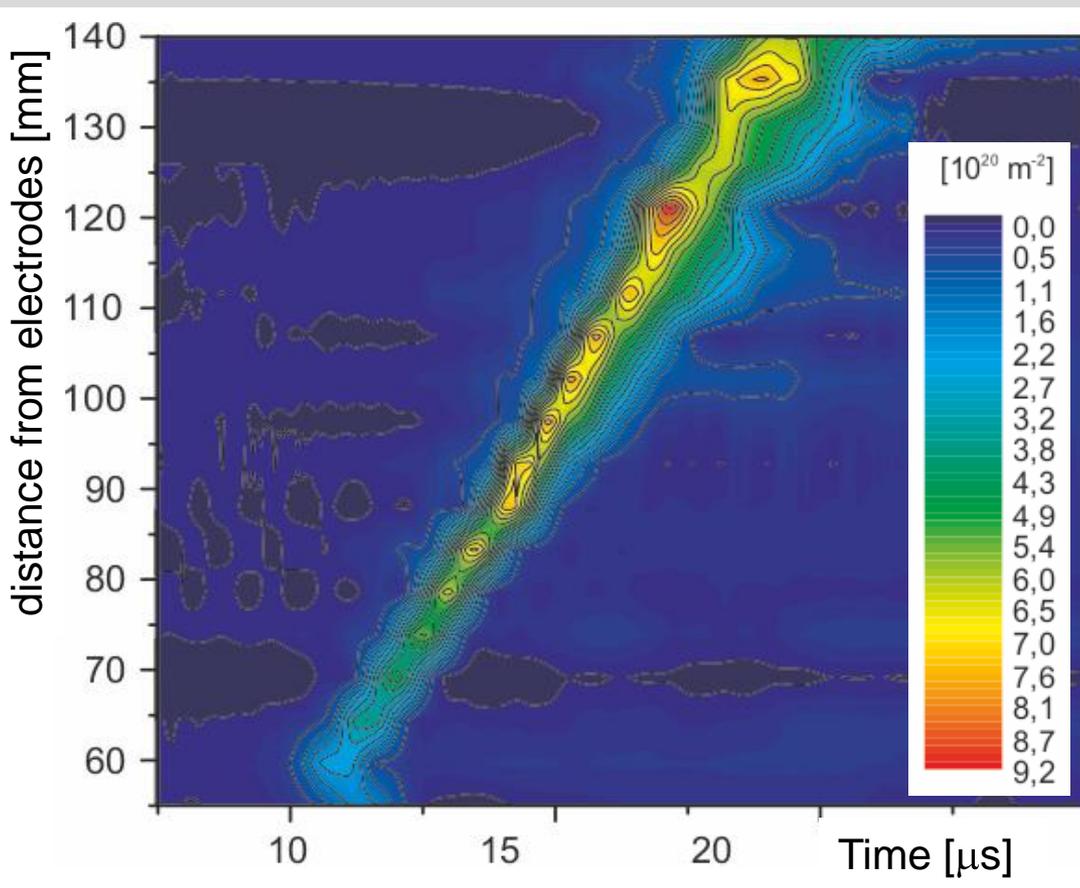
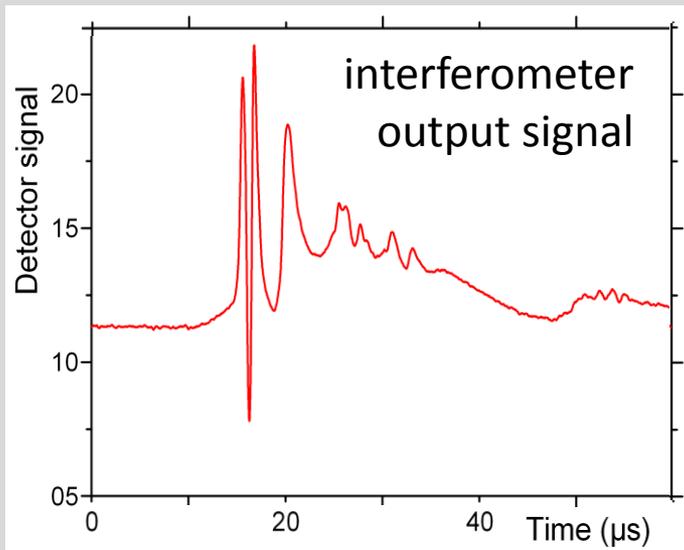
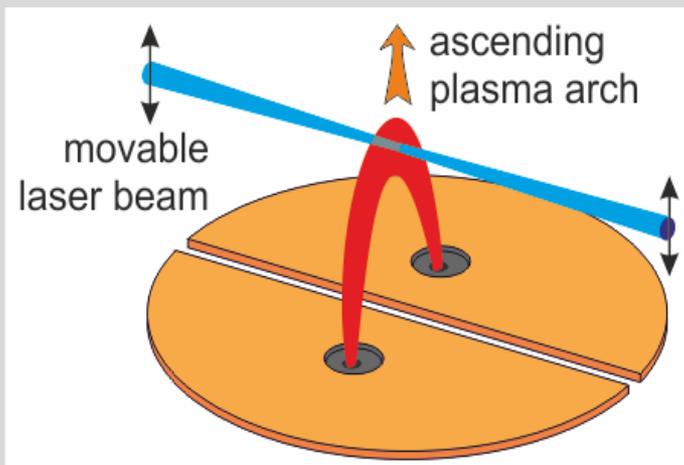


➤ movable magnetic pick-up coils exploiting Faraday's law → B_θ , B_ϕ



out-of-plane magnetic field component B_θ
as a function of distance from electrodes

➤ interferometer utilizing a CO₂ laser ($\lambda = 9.4 \mu\text{m}$) $\rightarrow \int_L n_e ds$



Phase shift in units of line-integrated electron density as a function of time and distance from the electrodes

plasma diameter	$d = 3 \times 10^{-2} \text{ m}$
plasma length	$L = 30 \times 10^{-2} \text{ m}$
electron density	$\bar{n}_e = 1 \times 10^{22} \text{ m}^{-3}$
electron temp.	$\bar{T}_e = 5 \text{ eV}$
current density	$j = 2,8 \times 10^7 \text{ A/m}^2$
magn. flux density	$B_p = 0.15 \text{ T}$
rate of current rise	$dI / dt = 4 \text{ kA}/\mu\text{s}$
injected gases	H_2, D_2, He, Ar

plasma beta:

$$\beta = \frac{n_e (T_e + T_i)}{B^2 / 2\mu_0} \approx 0.8$$

Lundquist number:

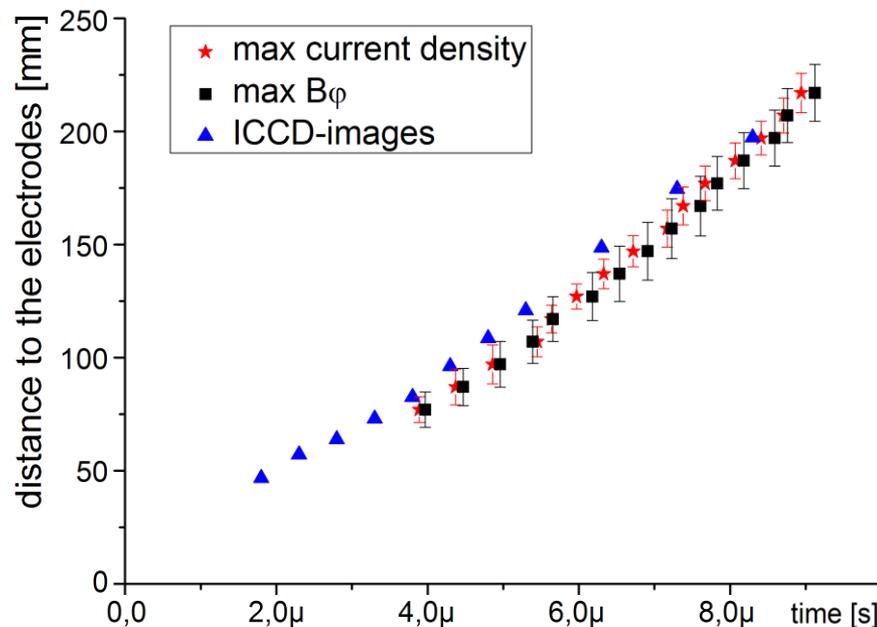
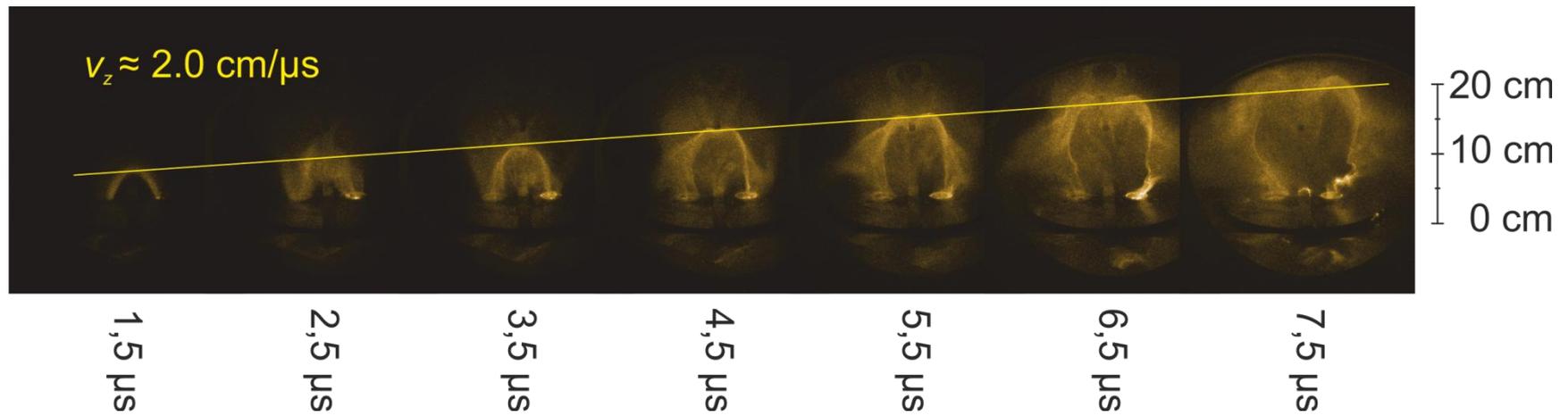
$$S = \mu_0 v_A \sigma \cdot L$$

$$\approx 190 \quad (L = 0.3 \text{ m})$$

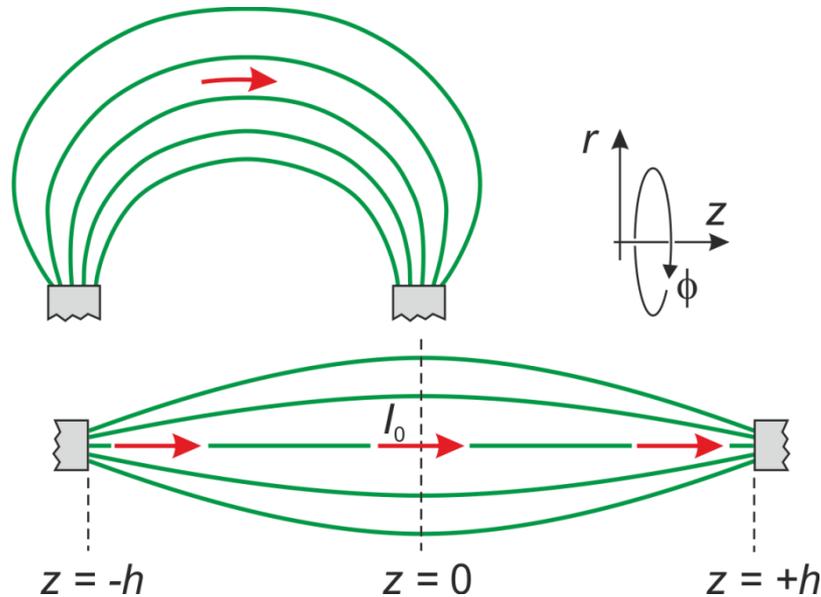
(res. skin time) : (plasma duration) $\tau_{\text{res}} : \tau_{\text{exp}} \cong 20\mu\text{s} : 8\mu\text{s} = 2.5$

Alfvén velocity \cong **electron drift velocity** \cong **expansion velocity** $\cong 2 \times 10^4 \text{ m/s}$

ion Larmor radius $\rho_{\text{ion}} \cong 1.5 \text{ mm}$; **ion skin depth** $c/\omega_{\text{pi}} \cong 3 \text{ mm}$



- **almost constant rate of ascend in the apex of the arch**
(velocity depends on gas species and on the rate of current rise)
- **almost constant width of the current channel**
(well-collimated flux tube with uniform cross section is formed shortly after gas breakdown)



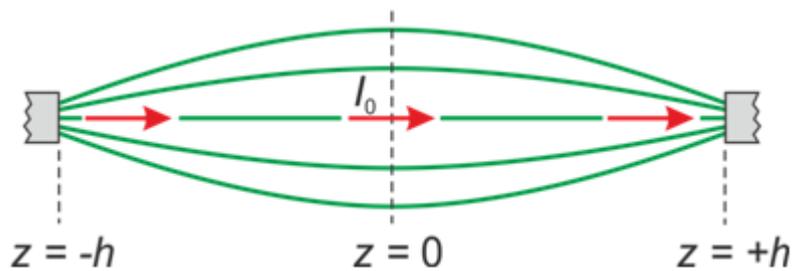
P. M. Bellan,
Why current-carrying magnetic flux tubes gobble up plasma and become thin as a result
 Phys. Plasmas **10**, 1999 (2003)

→ flows from both foot points towards the center due to a total axial force:

$$F_z = (\mathbf{J} \times \mathbf{B})_z - \frac{\partial P}{\partial z} = \frac{\mu_0 I_0^2}{2\pi^2 a^3} \left(1 - \frac{r^2}{a^2} \right) \frac{\partial a}{\partial z}$$

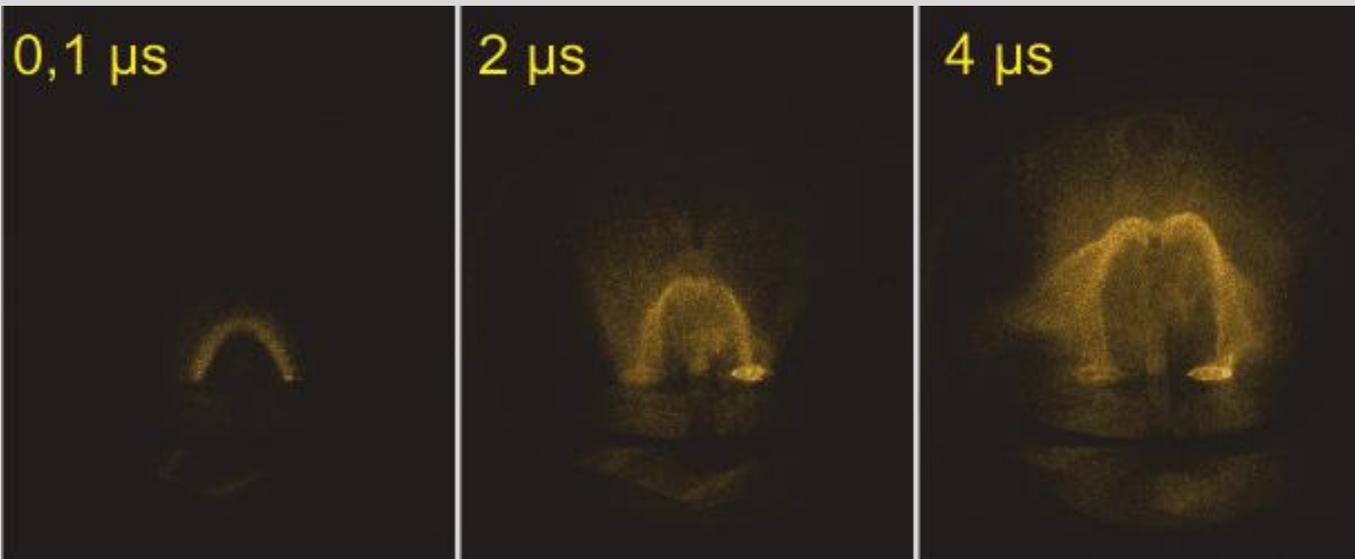
Predictions:

- the mechanism causes a strong increase of the plasma density by ingesting particles from the electrode regions into the bulged flux tube;
- the ingested plasma convects frozen-in magnetic flux from $z = \pm h$ to $z = 0$, which increases the local azimuthal field B_ϕ and hence the pinch force, thus reducing the tube radius $a(0)$ and forming a uniform cross-section.



$$F_z = (\mathbf{J} \times \mathbf{B})_z - \frac{\partial P}{\partial z} = \frac{\mu_0 I_0^2}{2\pi^2 a^3} \left(1 - \frac{r^2}{a^2} \right) \frac{\partial a}{\partial z}$$

$\frac{\partial a}{\partial z} = 0 \Rightarrow F_z = 0 \Rightarrow$ mechanism is not effective in a collimated flux tube

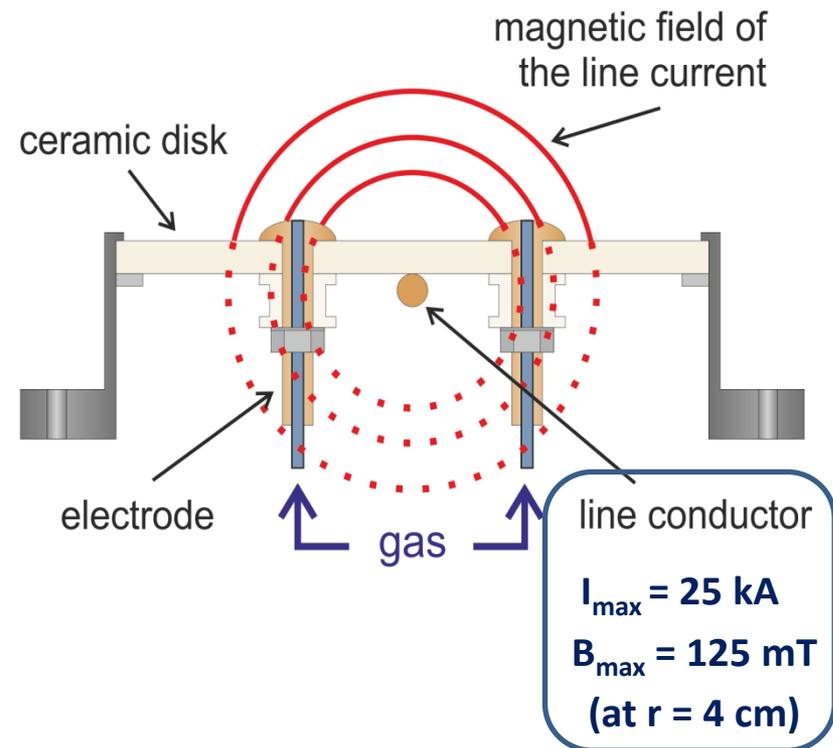
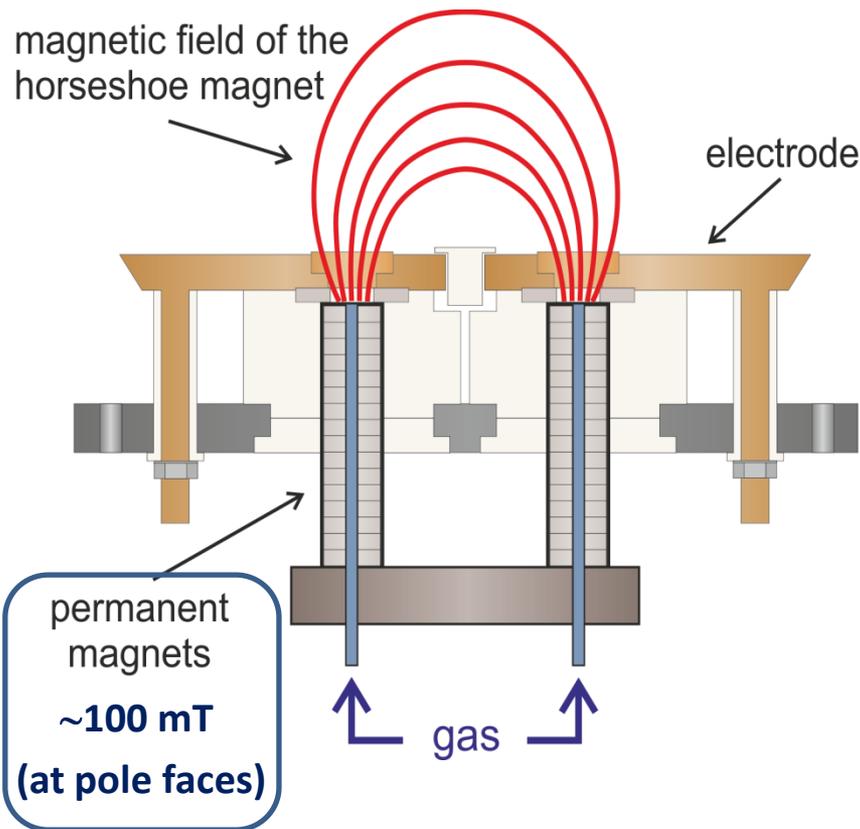


Initial discharge channel is thin and follows an arched path between the orifices in the electrodes;
 discharge remains well collimated during expansion.

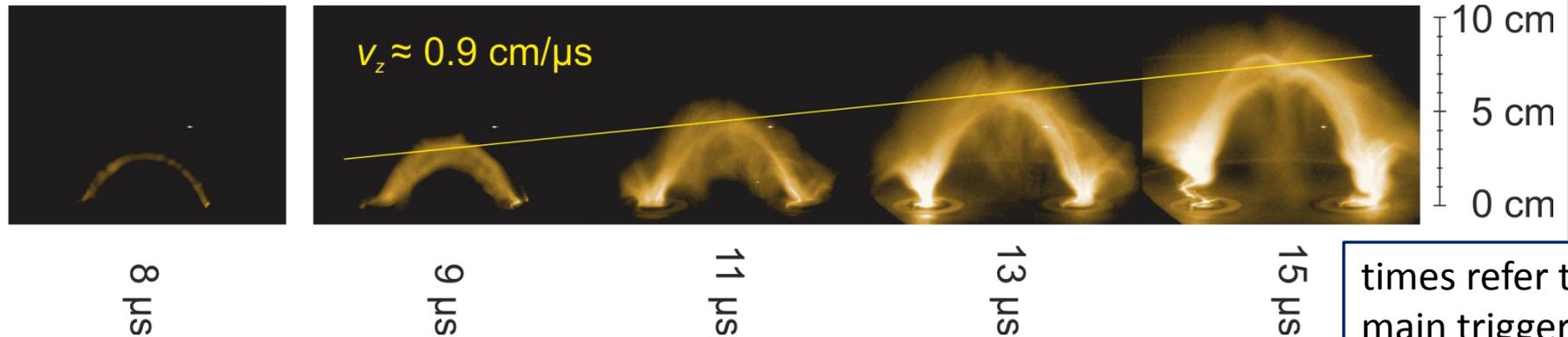
original plasma source
(bulged dipole field lines)



new line current source
(circular concentric field lines)



gas species: Argon; line current: 23 kA; main capacitor bank: ± 3 kV

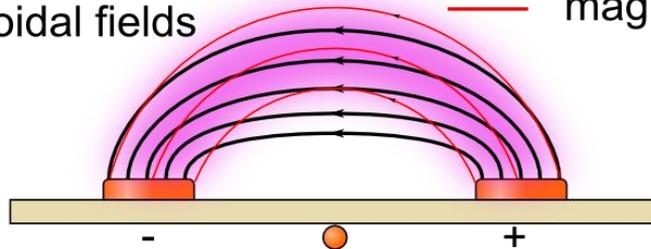


times refer to main trigger 8 μs before ignition.

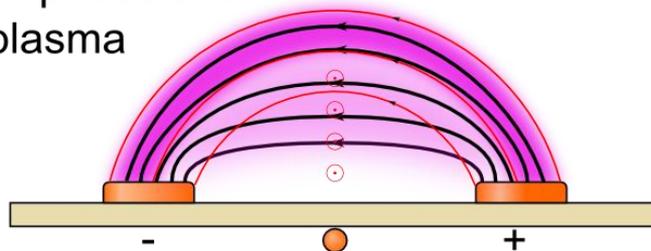
- **smaller rate of ascend because of heavier gas**
(Argon instead of Hydrogen-Helium mixtures as used in previous investigations)
- **general evolution of the discharge practically unchanged**
(close similarity to the temporal development observed in the original setup including a kink-like deformation in the apex at a later time)
- **plasma density in the same order of magnitude as before**
(neutral gas injection and discharge currents as before lead to comparable electron densities, although the proposed MHD pumping mechanism cannot be effective)

plasma ignition
along toroidal fields

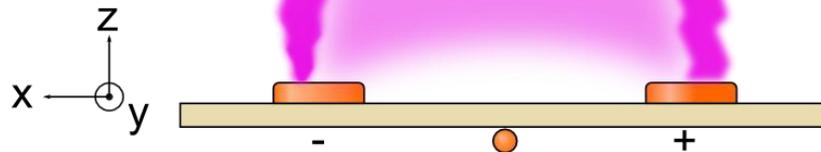
— electric fields
— magn. fields



collimation phase and
residual plasma



expansion phase
with constant drift velocity



the **electric field** between electrodes and the poloidal **magnetic field** due to the plasma current are perpendicular to each other and cause an **$E \times B$ -drift** in upward direction.

first step: **Ideal MHD**

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \vec{v}) + S$$

$$\frac{\partial \vec{v}}{\partial t} = -(\vec{v} \cdot \nabla) \vec{v} + \frac{\vec{j} \times \vec{B}}{\rho} - \frac{\nabla p}{\rho} + \nu \Delta \vec{v}$$

$$\frac{\partial \vec{B}}{\partial t} = \nabla \times (\vec{v} \times \vec{B})$$

$$\nabla \times \vec{B} = \mu_0 \vec{j}$$

ρ : mass density

\vec{v} : bulk velocity

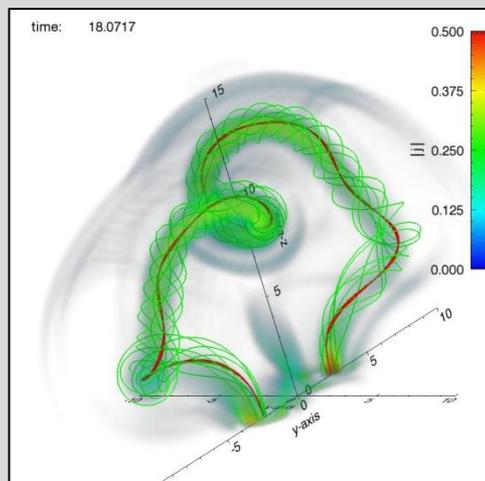
\vec{B} : magnetic field

p : thermal pressure

\vec{j} : current density

S : source term to implement different density models

$\nu \Delta \vec{v}$: viscosity term (formally included for numerical stability)

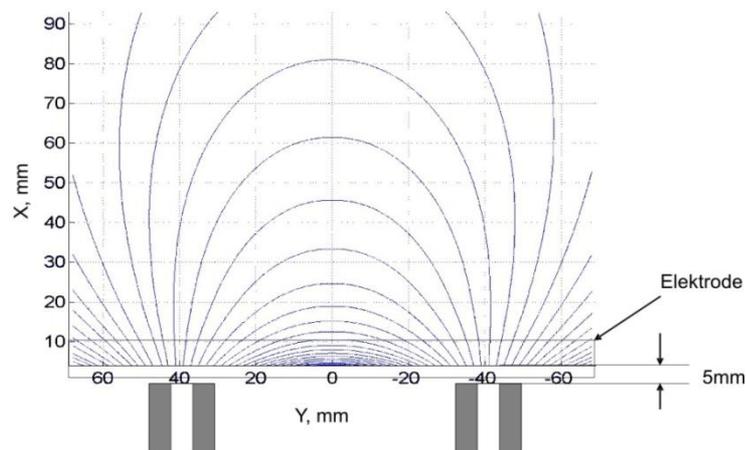


numerical implementation:

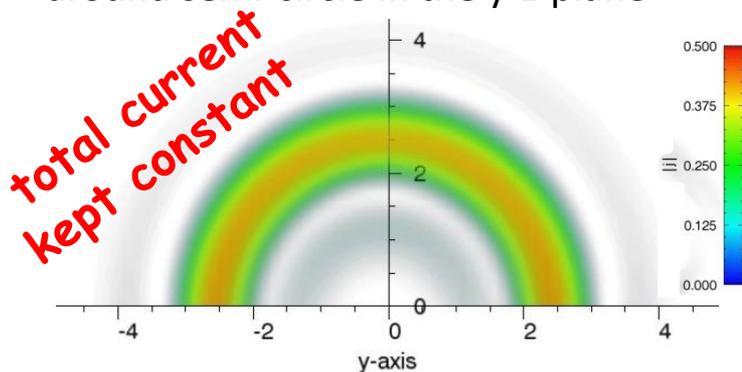
- finite differences on cartesian grid
- boundary condition in x-y-plane: $\vec{v} = 0$ (→ line tying)
- adaptive mesh refinement

initial conditions for the magnetic field

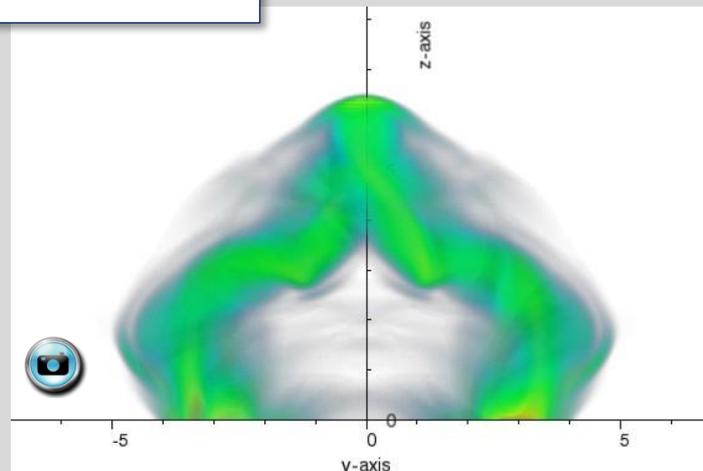
a) field of horseshoe magnet



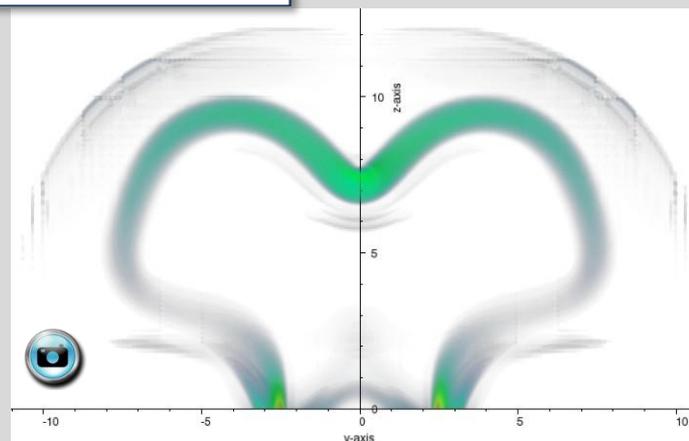
b) field of plasma current initially localized around semi-circle in the y-z-plane



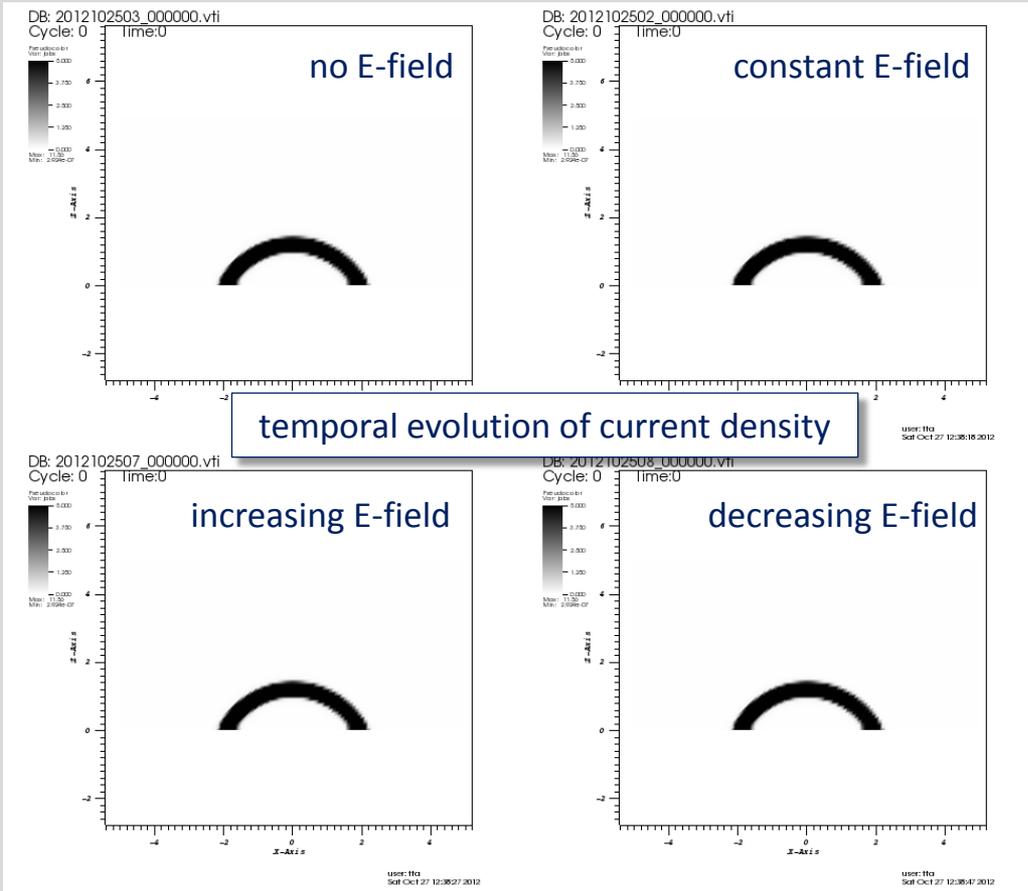
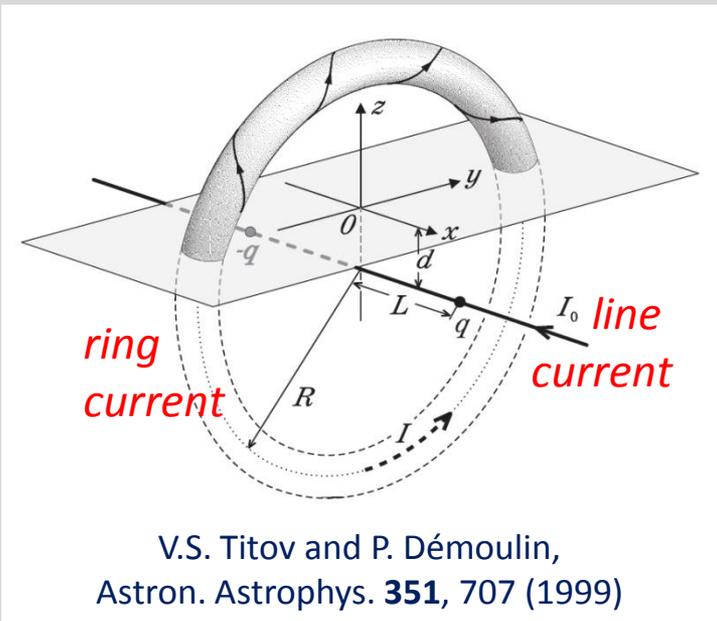
$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \vec{v}) \quad (\text{mass conservation})$$



$$\rho(\vec{x}, t) = \frac{|\vec{B}(\vec{x}, t)|^2}{v_A^2} \quad (\text{const. Alfvén velocity})$$



initial model configuration

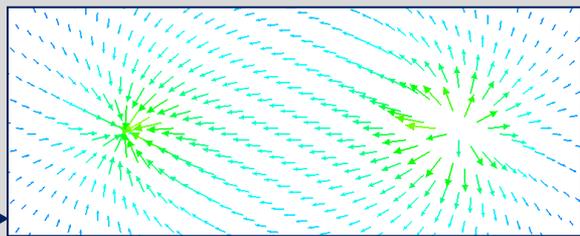


inclusion of finite resistivity

→ hinders the development of kinks

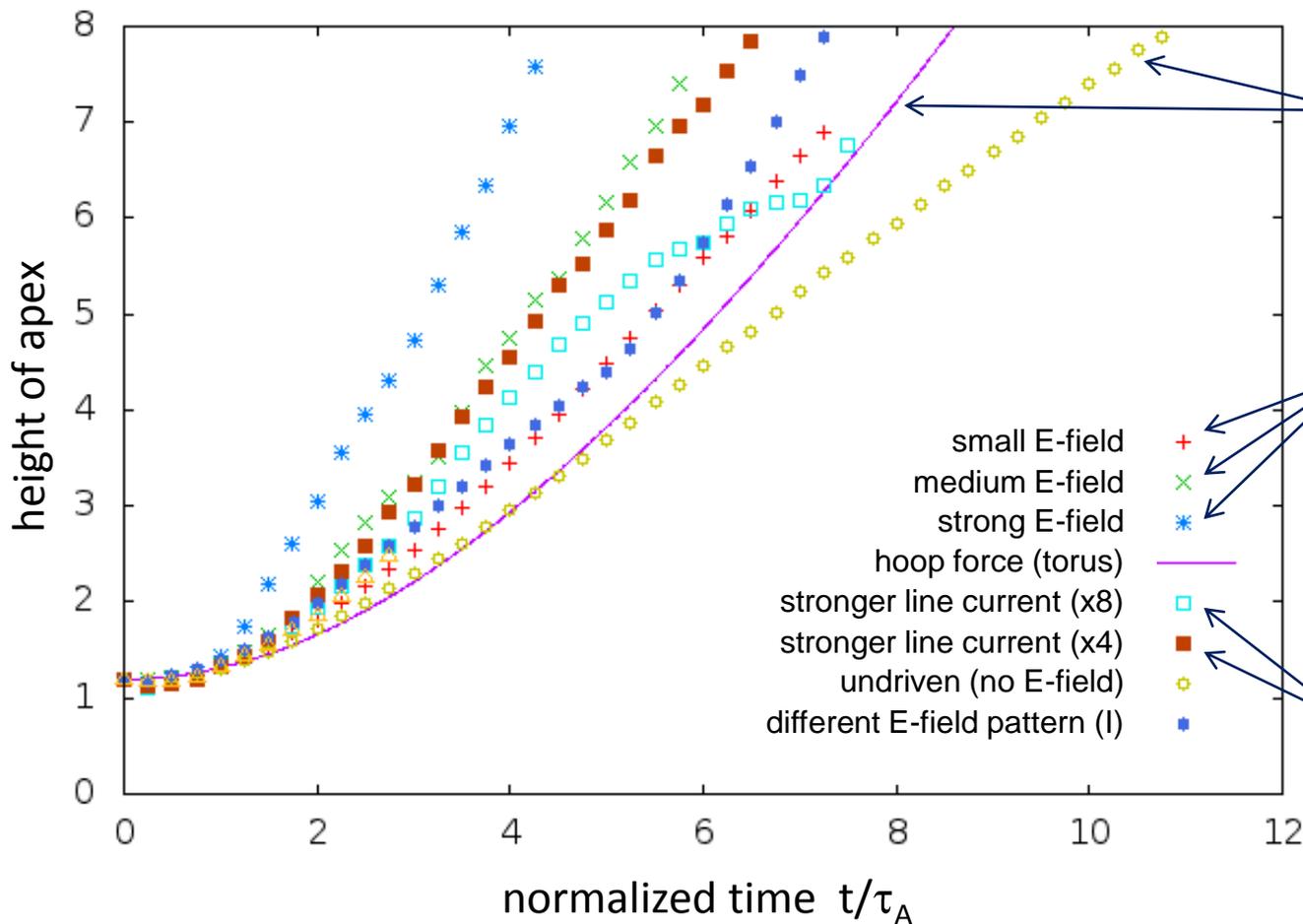
increase of the plasma current due to a prescribed E-field:

$$\mu_0 \frac{\partial}{\partial t} j_z = \frac{\partial}{\partial t} (\nabla \times \vec{B})_z = -(\nabla \times (\nabla \times \vec{E}))_z$$



tangential E-field at lower boundary reflecting the applied voltage between electrodes





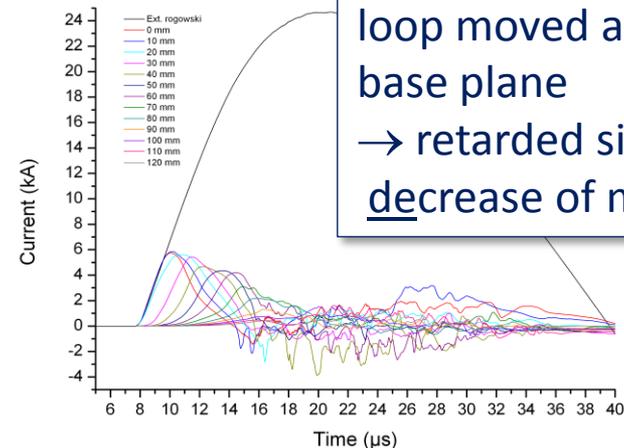
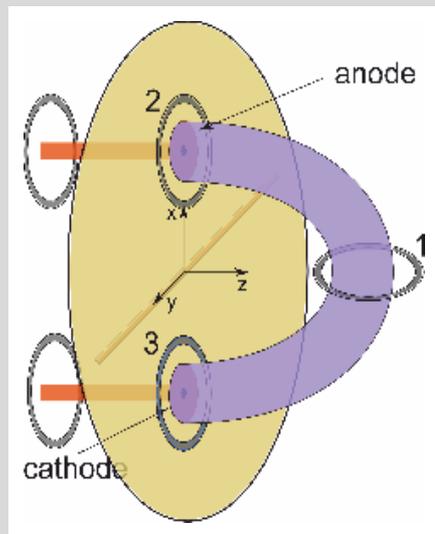
undriven case differs from toroidal hoop force expansion due to fixed foot points

driving enhances the rise velocity (as seen in the experiment)

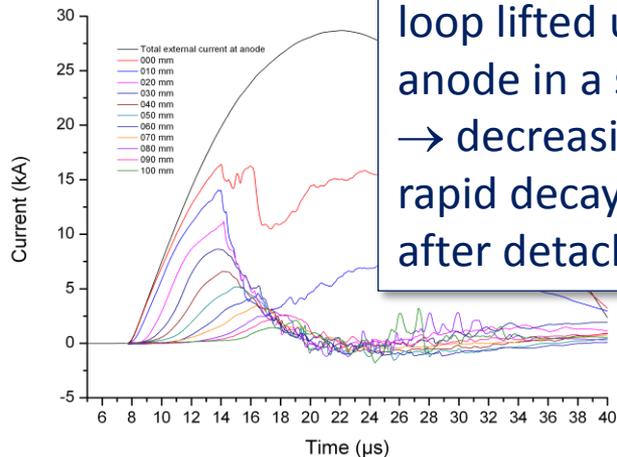
stronger line current (= less twist) shows similar behavior, but reduces the outward bulges at the flanks and the small-scale kinking in the apex

- early phase dominated by acceleration and expansion due to magnetic hoop force
- later phase characterized by nearly constant expansion velocity under varying conditions

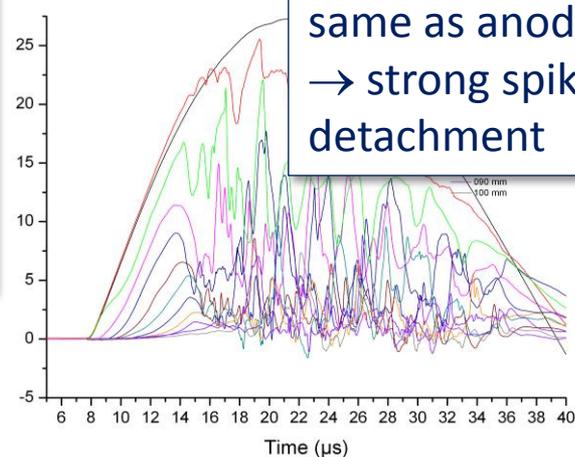
movable Rogowski coil (37 mm diameter) at various positions inside the vacuum vessel



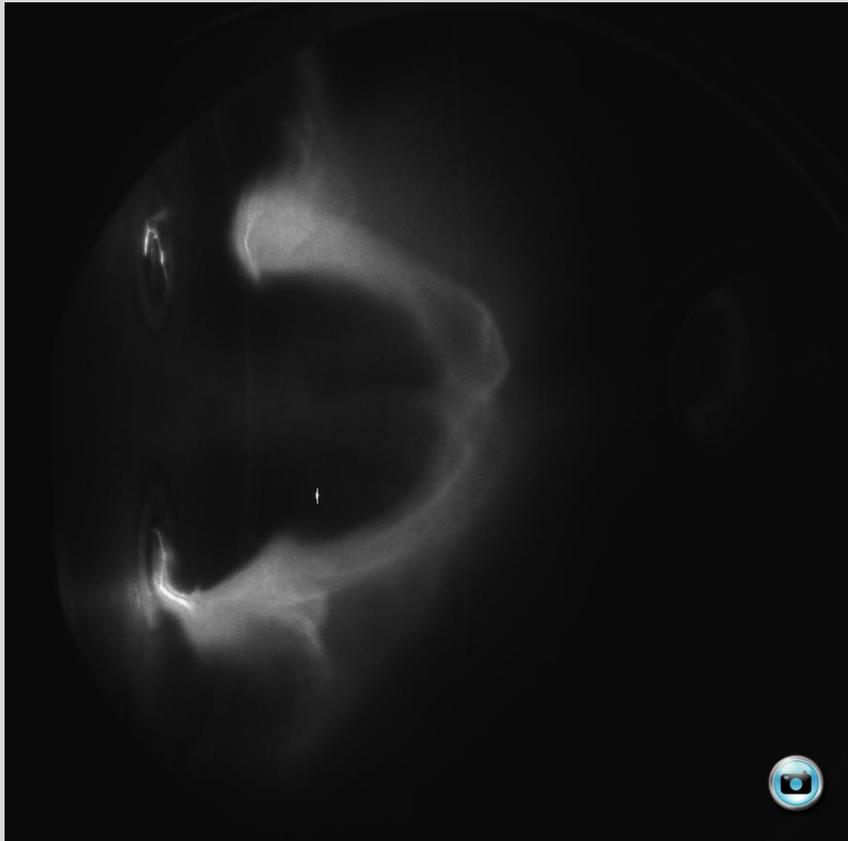
apex position:
loop moved away from base plane
→ retarded signal, but decrease of maximum



anode position:
loop lifted up from anode in a straight line
→ decreasing amplitude, rapid decay in each trace after detachment

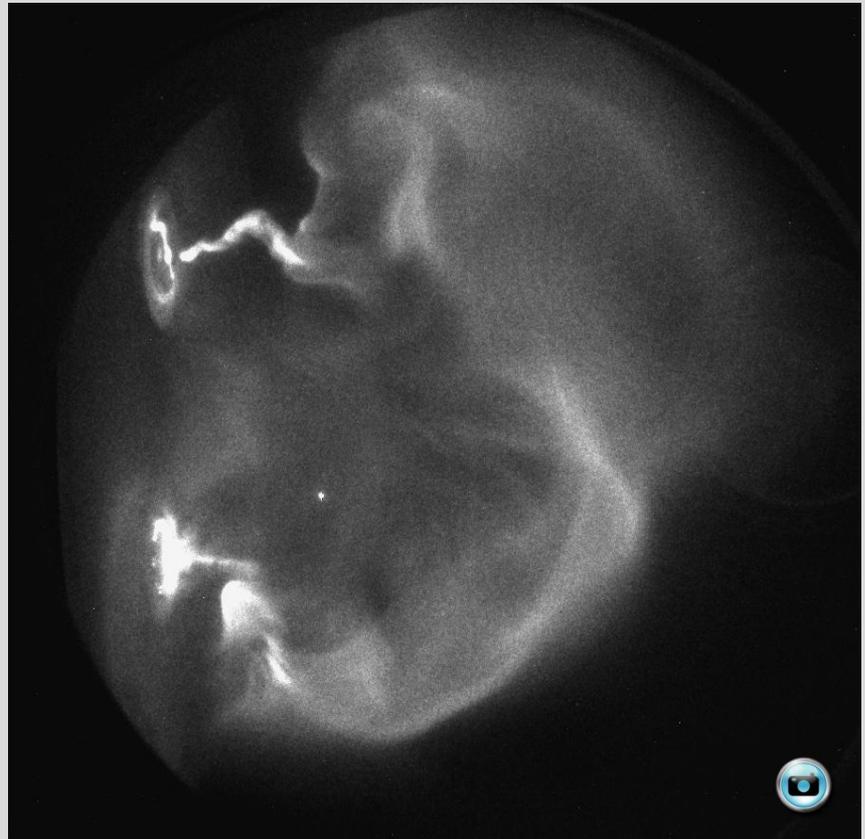


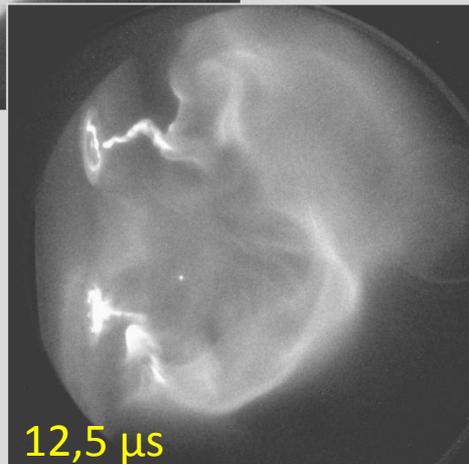
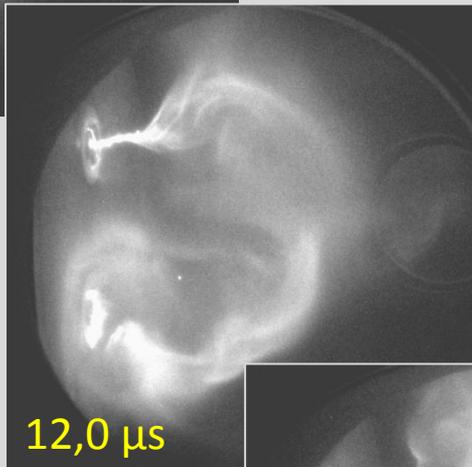
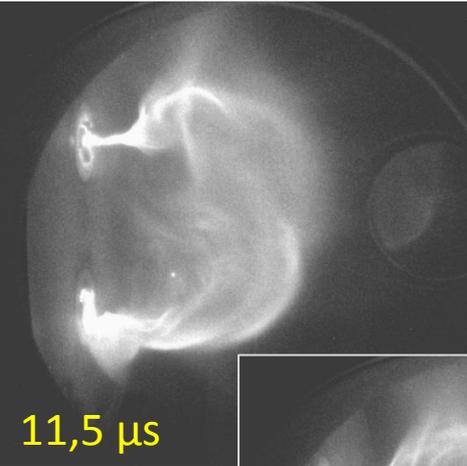
cathode position:
same as anode position
→ strong spiking after detachment



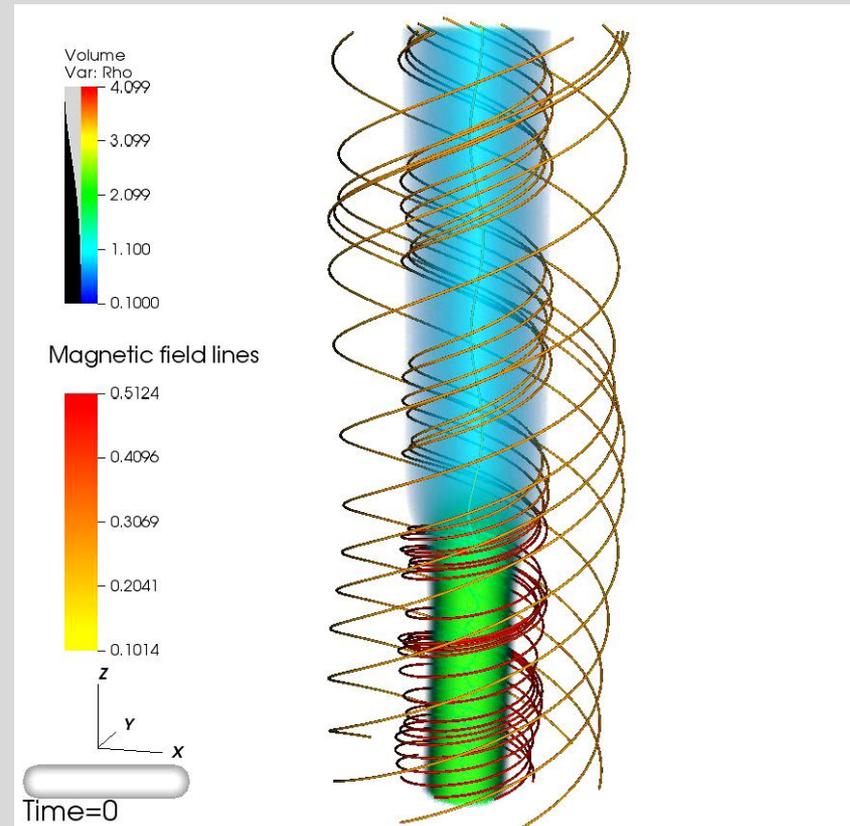
Discharge constricts and detaches from electrodes (starting at the anode); current continues to flow from the capacitor bank.

Contraction of the current channel at the electrodes and abrupt widening into a rather unshaped plasma bulk (similar to a tornado).





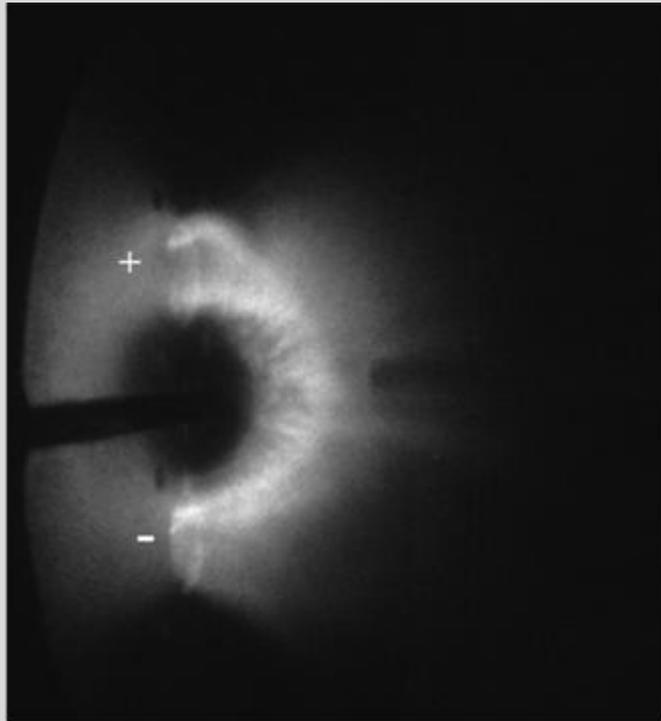
discharge in a He-H mixture;
 ± 3 kV main bank;
 24 kV aux. bank



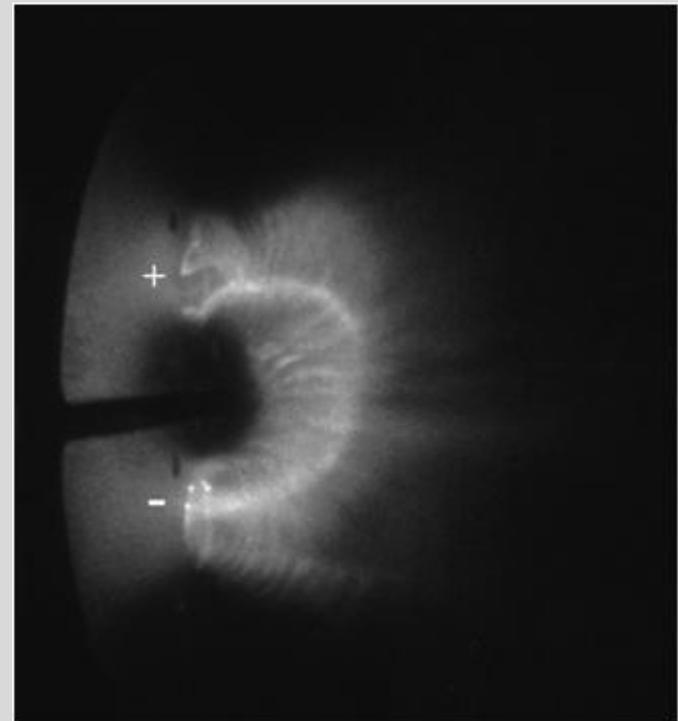
initial configuration is deformed and slightly widening in order to produce $j_r \times B_p$ force along the axis and to allow for growth of non-symmetric modes.

**Weak magnetic
guiding field:**

appearance of
persisting
periodic
structures
within the
arch-shaped
plasma column
and in its
surroundings

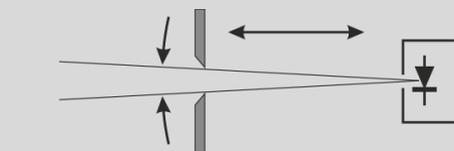
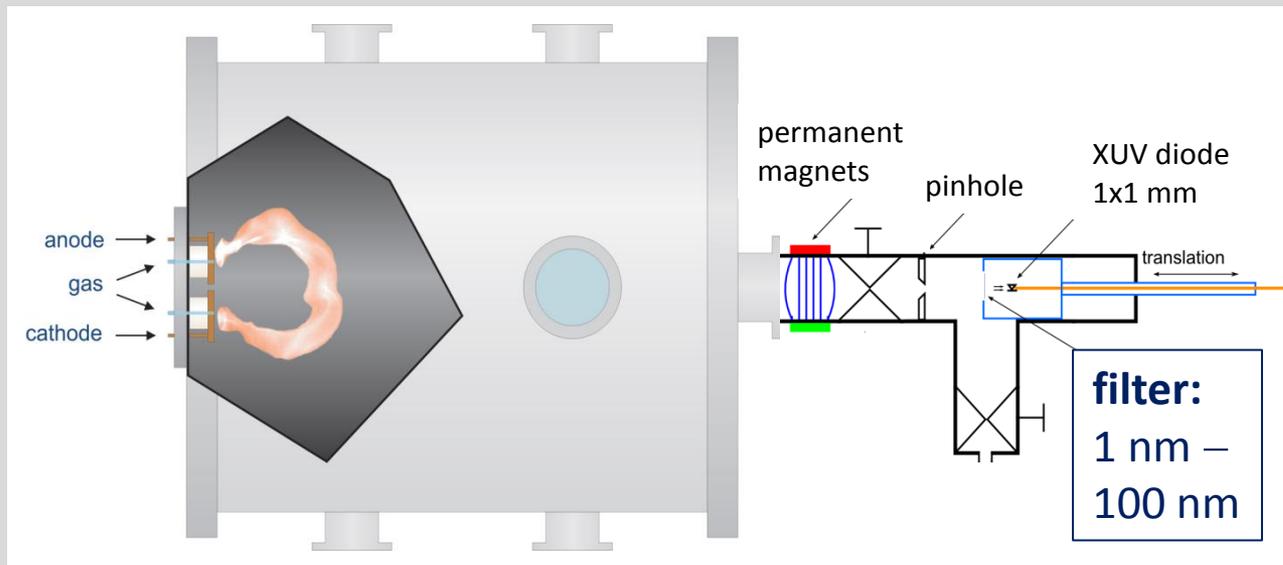
0.3 μ s after ignition

periodic modulation of
luminosity inside the arch

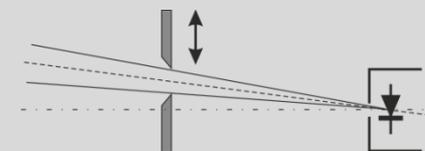
0.9 μ s after ignition

expansion of rib-like structure
underneath the arch and
appearance of ray-like tracers

→ Generation of fast particles ???

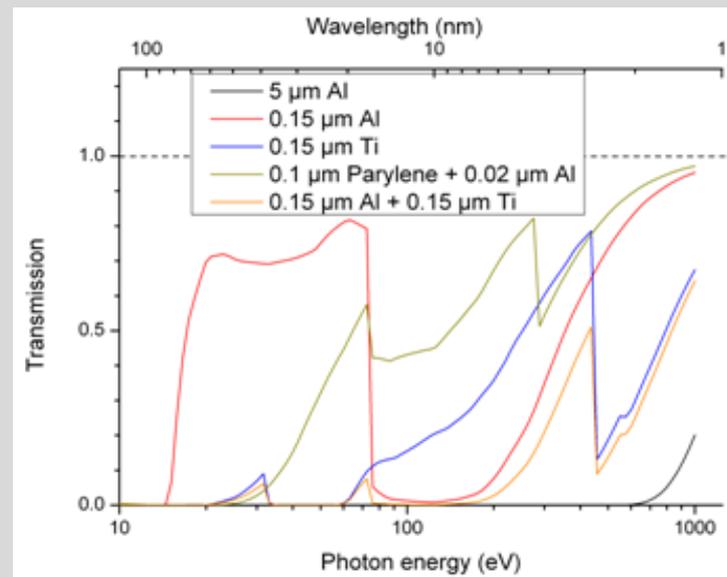


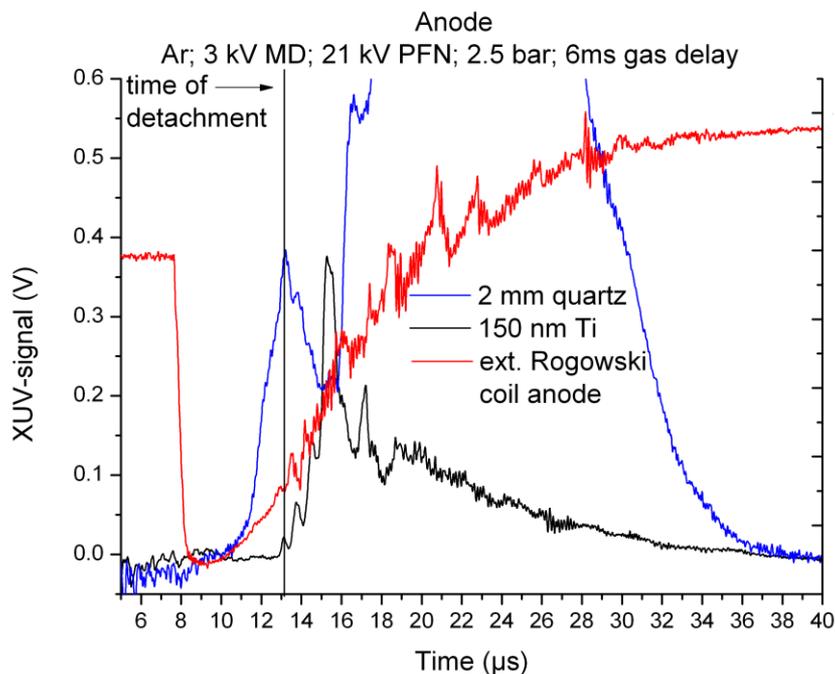
change of solid angle



change of viewing direction

fast photodiode sensitive to visible and XUV radiation is installed at the back of the vessel;
permanent magnets in front of the pinhole deflect fast particles before they can reach the detector;
 solid angle and viewing direction can be changed by means of a **movable pinhole**;
thin metal foils transmit energetic photons only.

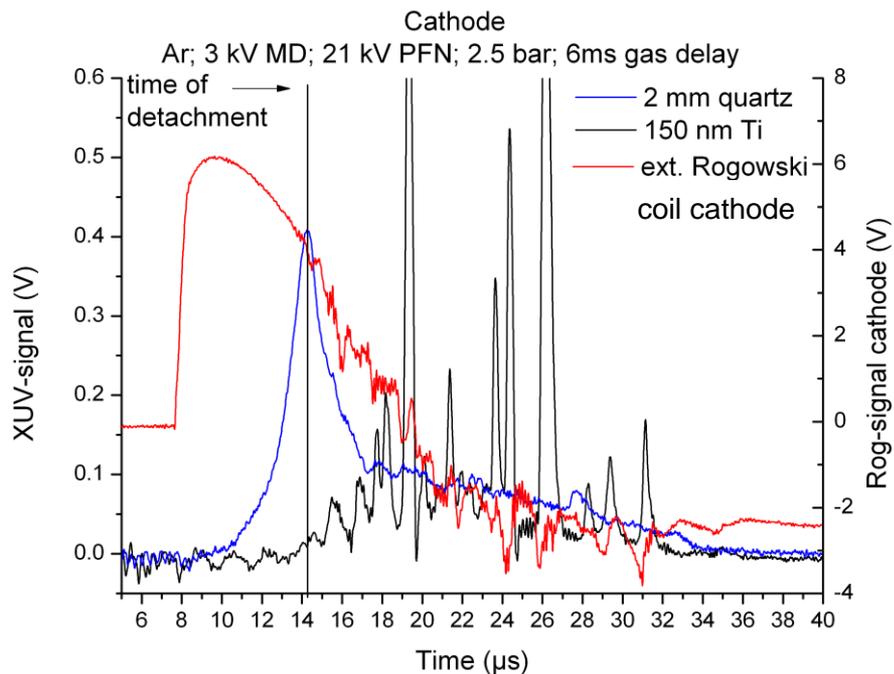


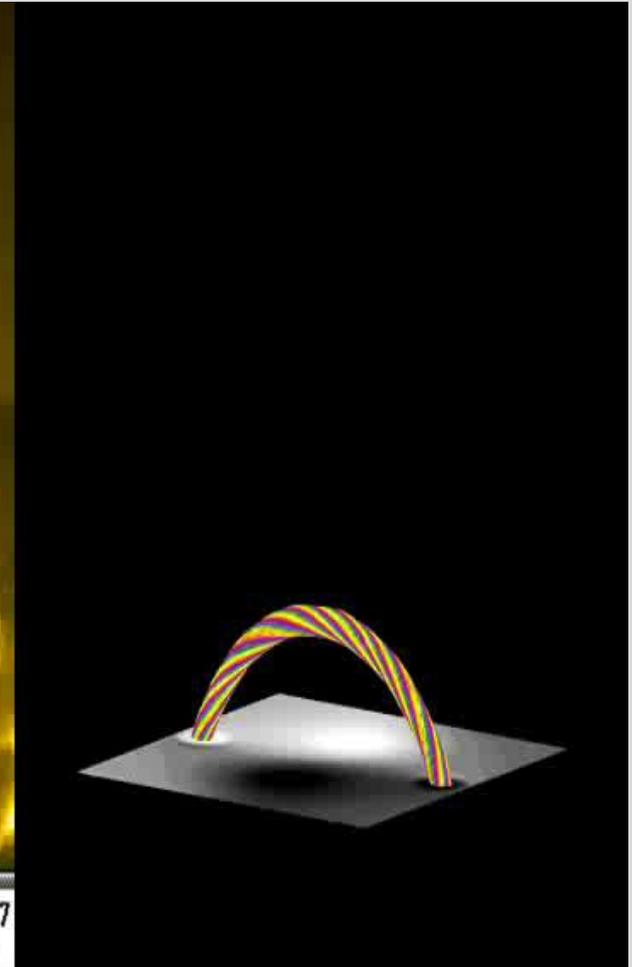
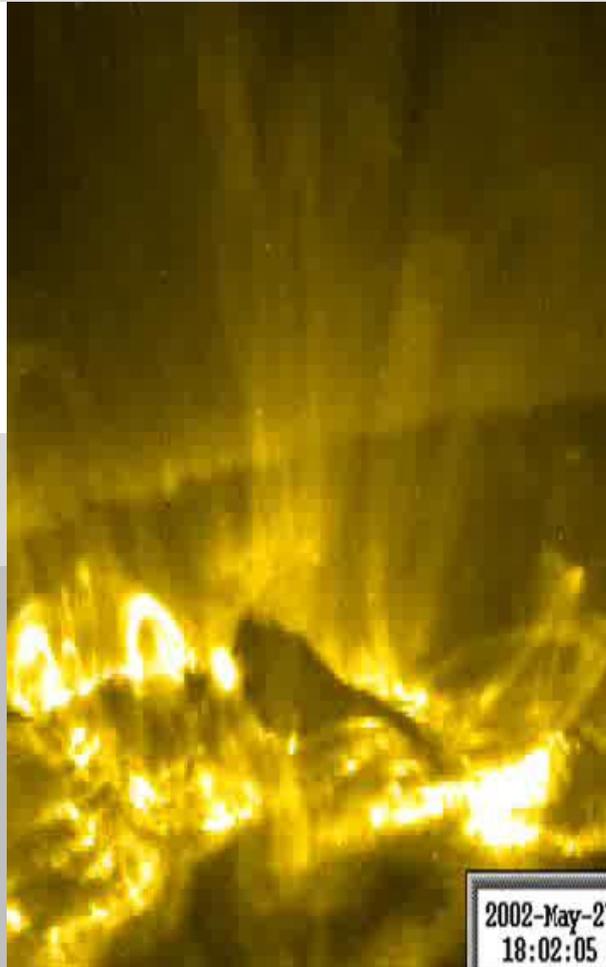
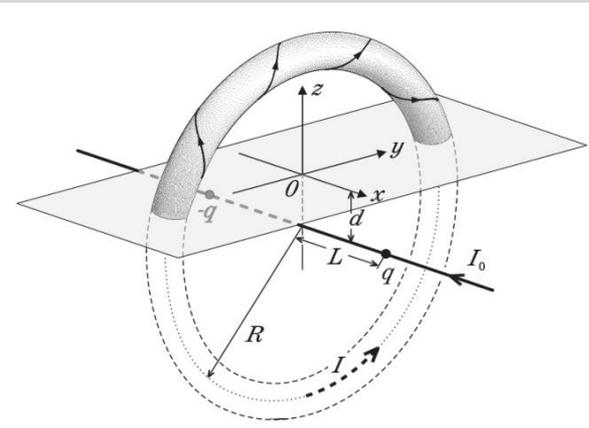


← **Anode region**

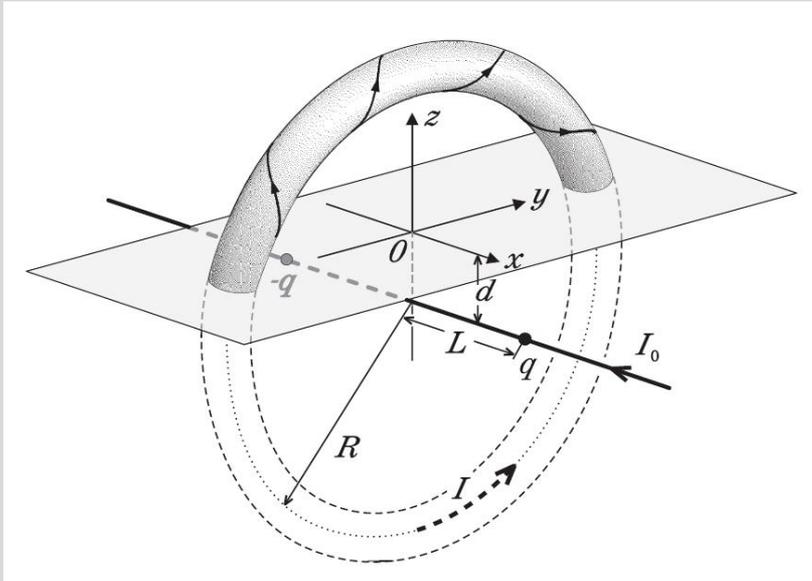
- visible light (blue curve) decreases at time of detachment and increases strongly when XUV decays
- XUV light (>60 eV, black curve) starts after detachment and decays fairly smoothly

- Cathode region →**
- visible light (blue curve) decreases at time of detachment and decays fairly smoothly
 - XUV light (>60 eV, black curve) increases slowly after detachment and shows strong spiking afterwards

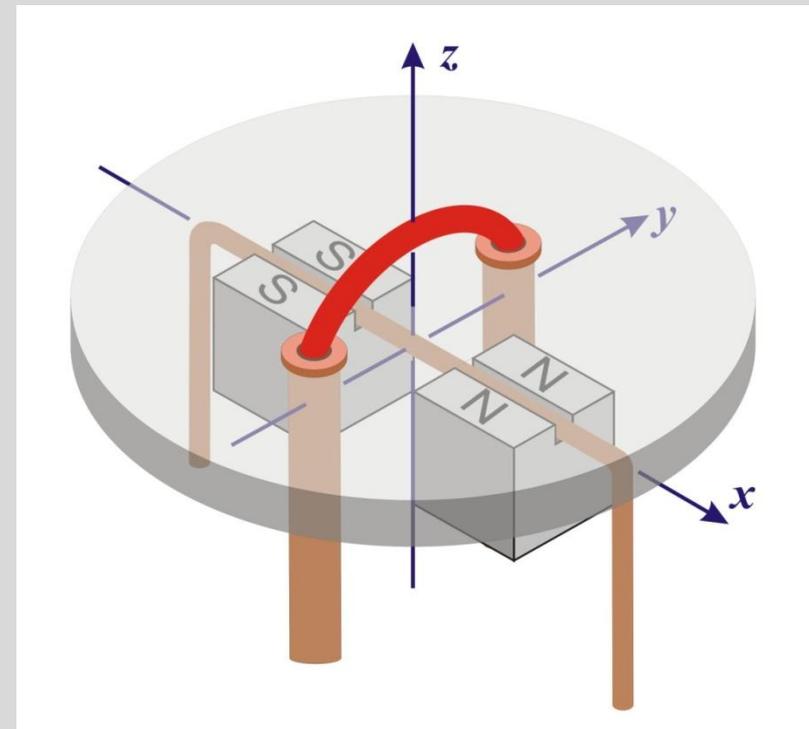




simulation by
 T. Török and B.Kliem,
 (Astrophys. Inst. Potsdam)
Astrophys. J. **630**,
 L97-L100 (2005)



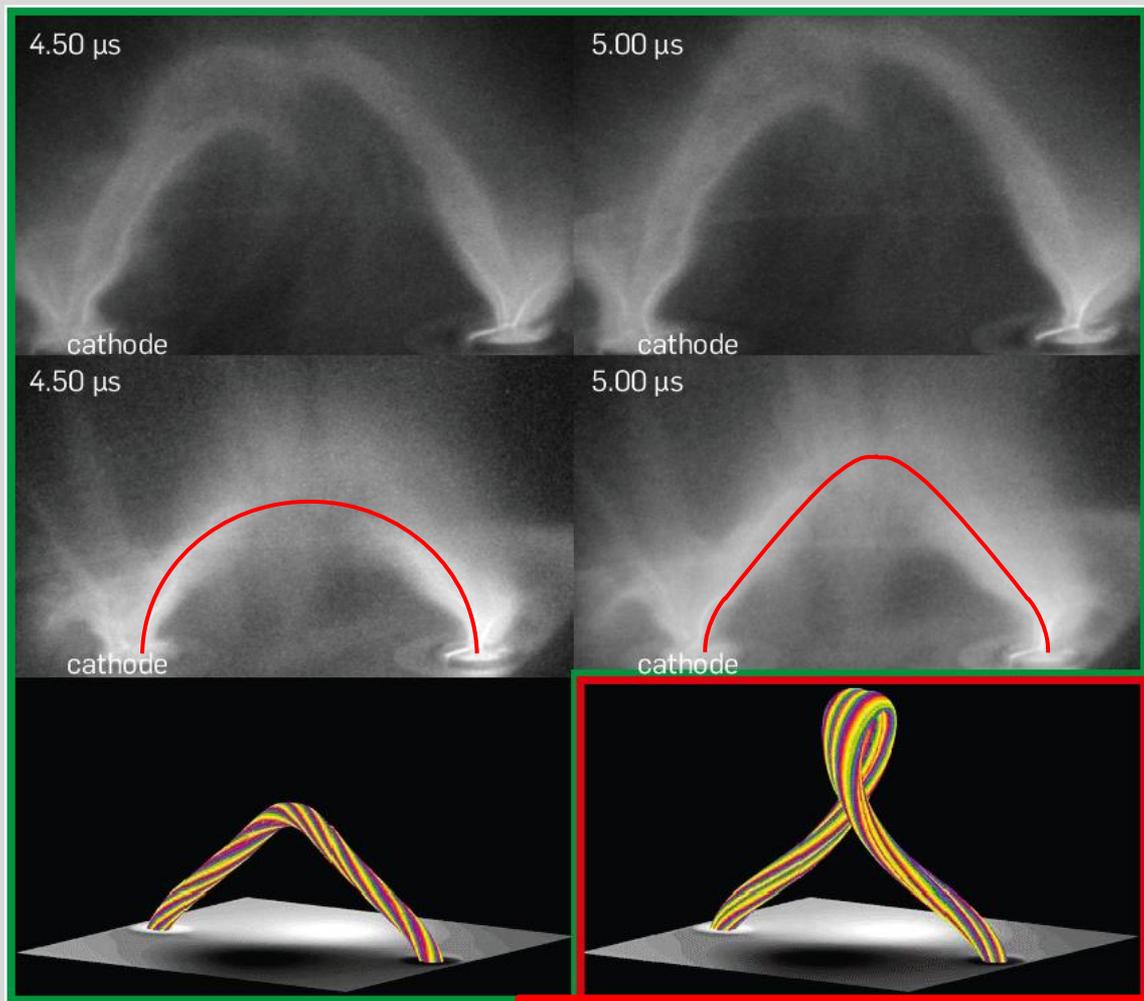
experimental realisation



$q, -q$: „magnetic charges“ to produce dipole field

I_0 : line current to produce toroidal magnetic field

I : ring current to produce poloidal magnetic field



without strapping field

plasma shows a relatively fast and uniform expansion

with strapping field

- flux tube expands more slowly,
- shape of flux tube evolves from semi-circular to triangular.

simulation by Török and Kliem

starting from a semi-circular configuration, the flux tube attains a triangular shape and develops a strong outward twist.

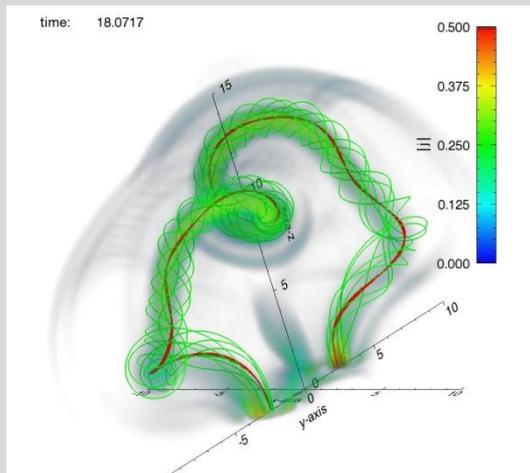
not yet observed in experiment

reliable and fairly reproducible plasma production for different configurations and operational conditions:

- high degree of ionization, but rather resistive;
- prone to macroscopic and small-scale instabilities;
- generation of fast particles (and energetic photons);
- phenomena near electrodes (“tornados”, detachment)



P.M. Bellan & J.F. Hansen,
Proc. of ISSS-6, 192 (2001)



extensive time-dependent 3-dimensional MHD calculations using adaptive mesh refinement:

- arch expansion for different density models;
- arch expansion using special drivers for $\Delta I_p > 0$;
- development of kink modes;
- simulation of “tornados” and detachment.

L. Arnold, J. Dreher et al.,
Phys. Plasmas **15**, 042106 (2008)