



Cascade and dissipation of solar wind turbulence from MHD to electron scales

Fouad Sahraoui¹

G. Belmont¹, S. Huang³, M. Goldstein², J.L. Pinçon³, A. Retino¹, Y. Narita⁴, K. Kiyani⁶

1 LPP, CNRS-Ecole Polytechnique, France

2 NASA/GSFC, Maryland, USA

3 University of Wuhan, China

4 LPC2E, Orléans, France

5 Max Plank Institute, Allemagne

6 Imperial college, UK

Outline

1. Solar wind turbulence vs heating
2. The problem of measuring spatial properties of space plasma turbulence
3. 3D spatial spectra and anisotropies of MHD turbulence in the solar wind (Cluster data)
4. Kinetic scales in the SW: Some hotly debated question vs Cluster observations
 - *Cascade or dissipation below ρ_i ?*
 - *The scaling: power-law? Exponential? Others?*
 - *The nature of the cascade: KAW? whistler? Others?*
 - *The nature of the dissipation: wave-particle interactions? Current sheets/Reconnection?*
 - *Weak vs strong turbulence? Monofractality vs multifractality?*
5. Conclusions & perspectives (turbulence & the future space missions)

Turbulence in the Univers



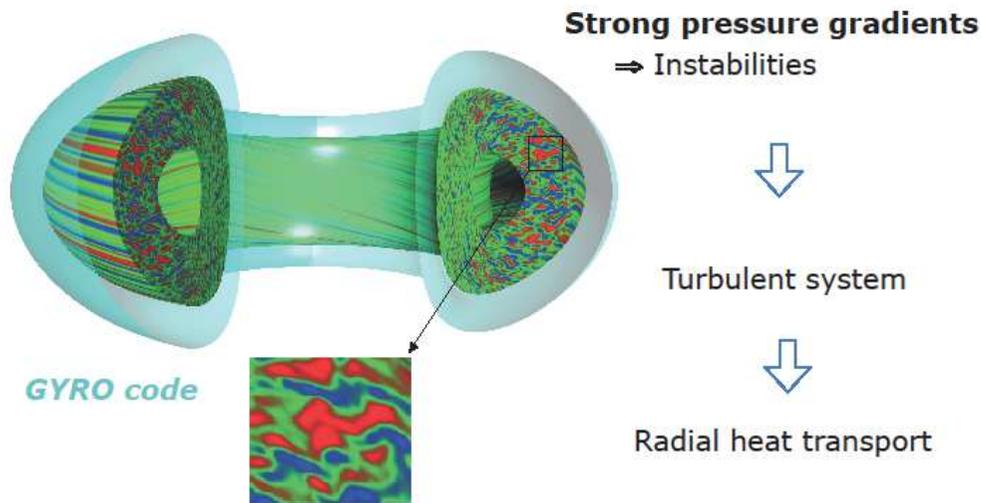
M100 galaxy $10^{23} m$

$N \sim 10^6 \text{ cm}^{-3}$
 $T_i \sim 10^{12} \text{ K}$
 $B \sim 10^6 \text{ nT}$



Eagle nebula $10^{18} m$

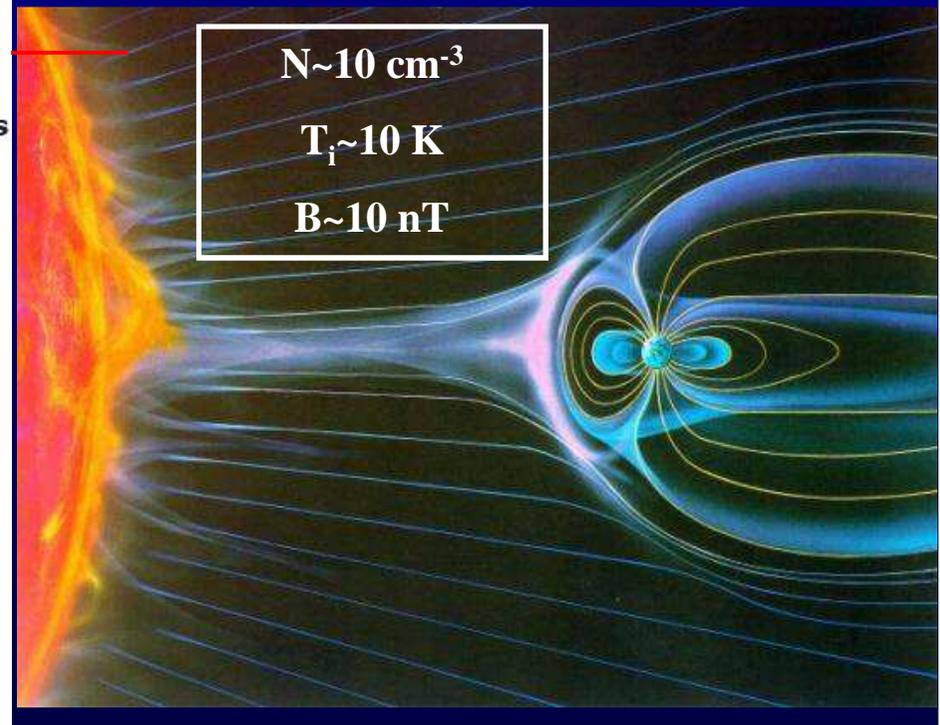
Performances limited by plasma turbulence



⇒ Prediction & control of turbulent transport

It is observed from quantum to cosmological scales!

It controls mass transport, energy transfers & heating, *magnetic reconnection in plasmas*, ...



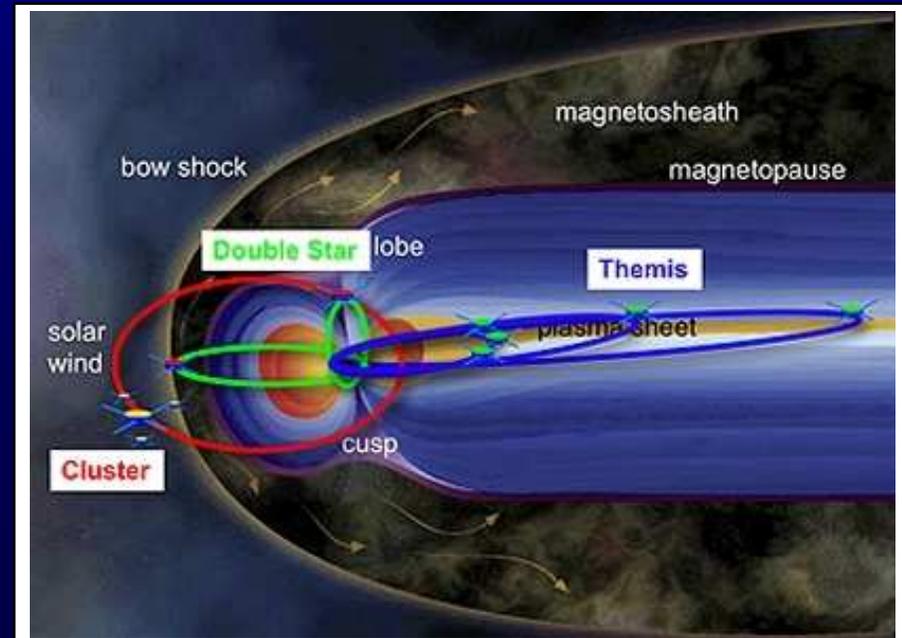
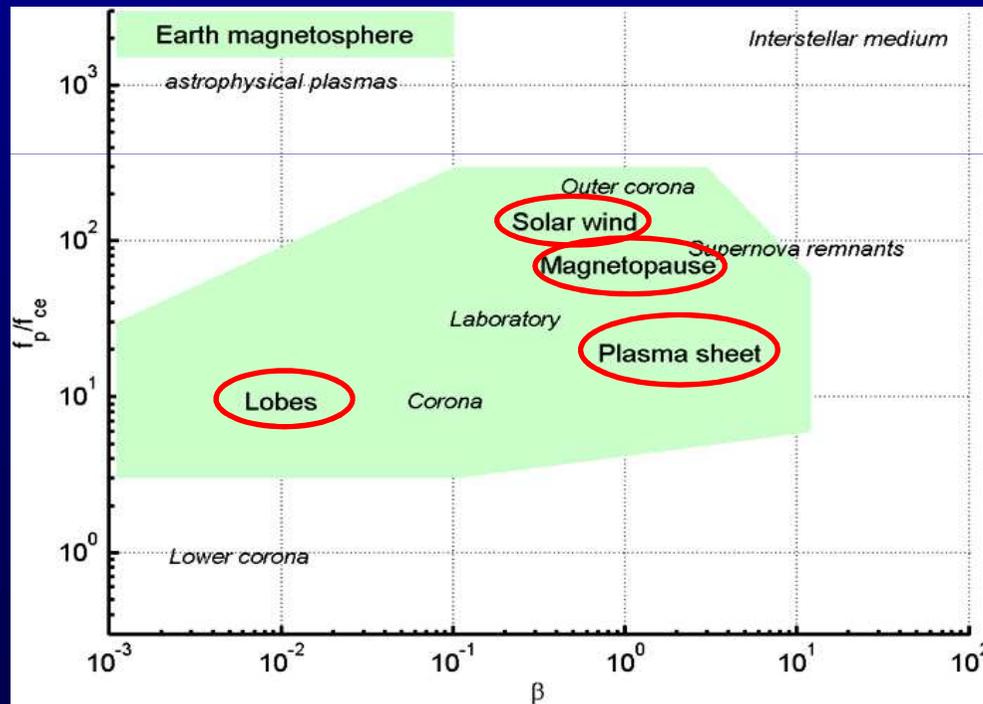
Near-Earth space plasmas

[Scheckochihin et al., ApJ, 2009]

$$\beta = \frac{\text{Pression thermique}}{\text{Pression magnétique}} \approx 0.4 \frac{NT}{B^2}$$

Representative Parameters for Astrophysical Plasmas

Parameter	Solar wind at 1 AU ^(a)	Warm ionized ISM ^(b)	Accretion flow near Sgr A* ^(c)	Galaxy clusters (cores) ^(d)
$n_e = n_i, \text{ cm}^{-3}$	30	0.5	10^6	6×10^{-2}
$T_e, \text{ K}$	$\sim T_i^{(e)}$	8000	10^{11}	3×10^7
$T_i, \text{ K}$	5×10^5	8000	$\sim 10^{12(f)}$	$?^{(e)}$
$B, \text{ G}$	10^{-4}	10^{-6}	30	7×10^{-6}
β_i	5	14	4	130
$v_{thi}, \text{ km s}^{-1}$	90	10	10^5	700
$v_A, \text{ km s}^{-1}$	40	3	7×10^4	60
$U, \text{ km s}^{-1(f)}$	~ 10	~ 10	$\sim 10^4$	$\sim 10^2$
$L, \text{ km}^{(f)}$	$\sim 10^5$	$\sim 10^{15}$	$\sim 10^8$	$\sim 10^{17}$
$(m_i/m_e)^{1/2} \lambda_{mfpi}, \text{ km}$	10^{10}	2×10^8	4×10^{10}	4×10^{16}
$\lambda_{mfpi}, \text{ km}^{(g)}$	3×10^8	6×10^6	10^9	10^{15}
$\rho_i, \text{ km}$	90	1000	0.4	10^4
$\rho_e, \text{ km}$	2	30	0.003	200

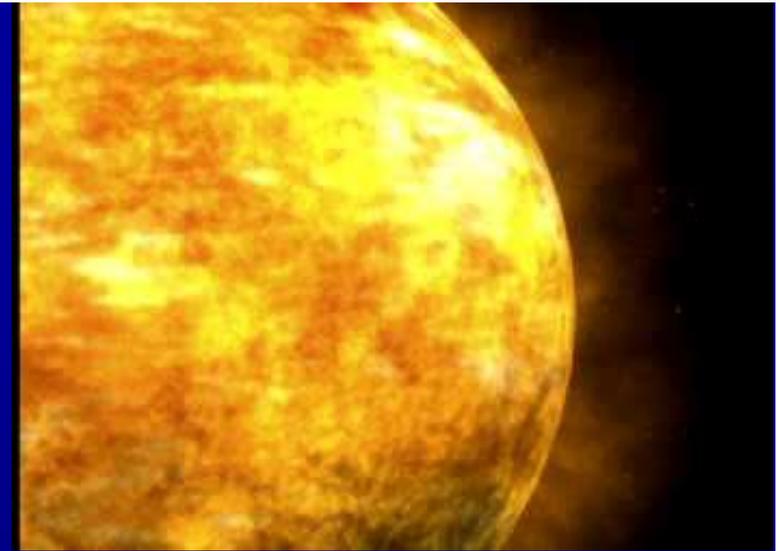


[Vaivads et al., Plasma Phys. Contr. Fus., 2009]

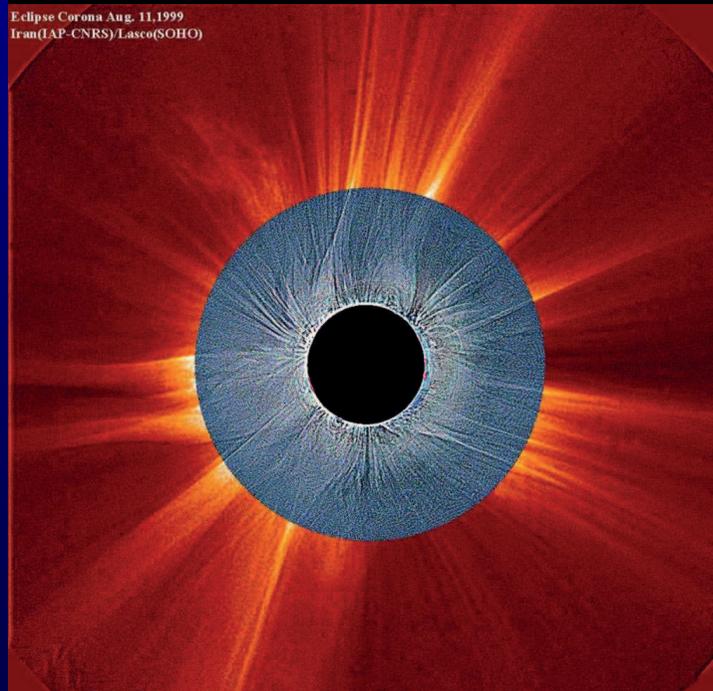
The solar wind

The solar wind plasma is generally:

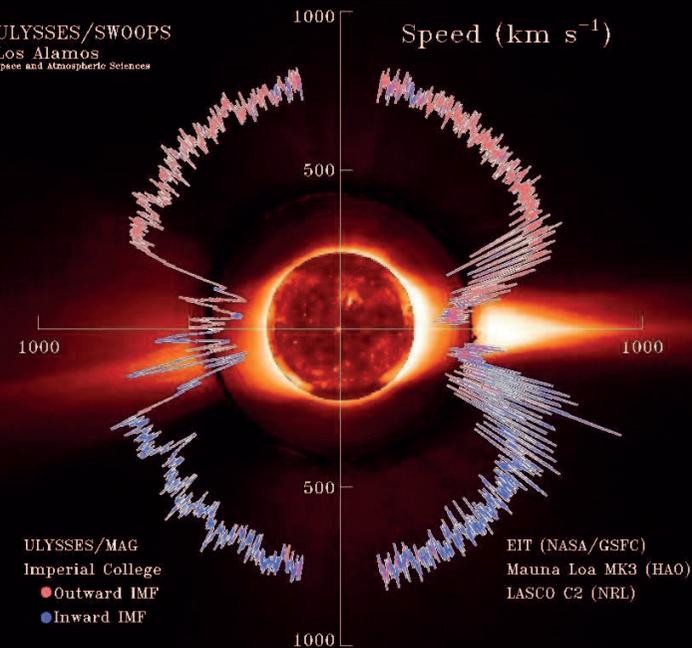
- Fully ionized (H^+ , e^-)
- Non-relativistic ($V_A \ll c$), $V \sim 350-800$ km/s
- *Collisionless*



Le champ magnétique structure la couronne



ULYSSES/SWOOPS
Los Alamos
Space and Atmospheric Sciences

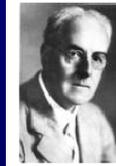
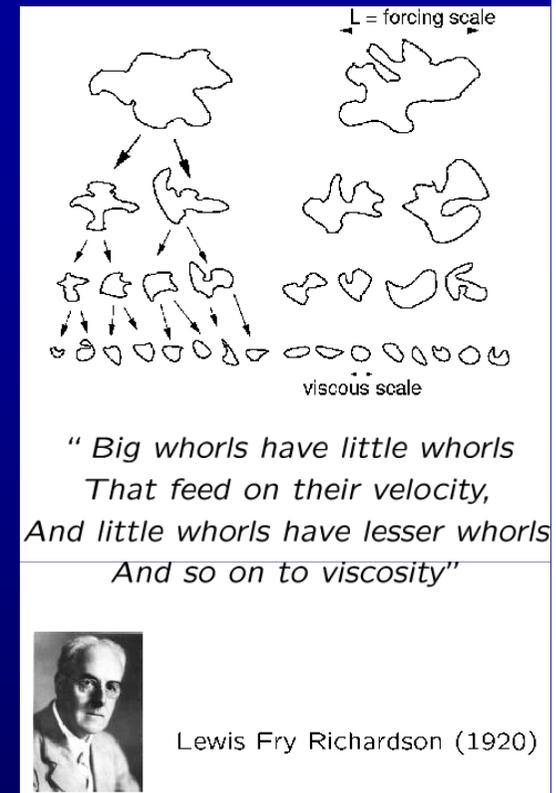
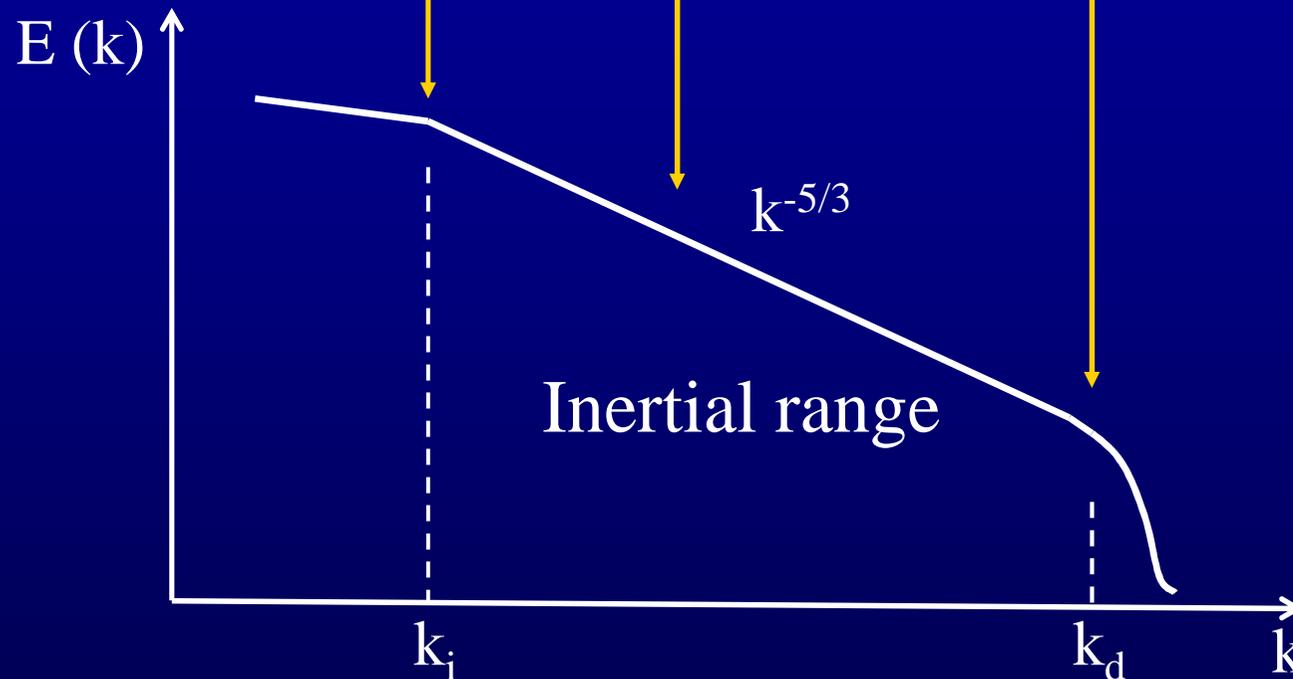


La couronne chaude crée l'héliosphère

Phenomenology of turbulence

NS equation:

$$\partial_t \mathbf{V} + \mathbf{F}_i = -\mathbf{V} \cdot \nabla \mathbf{V} - \nabla P + \nu \nabla^2 \mathbf{V}$$



Lewis Fry Richardson (1920)

Courtesy of A. Celani

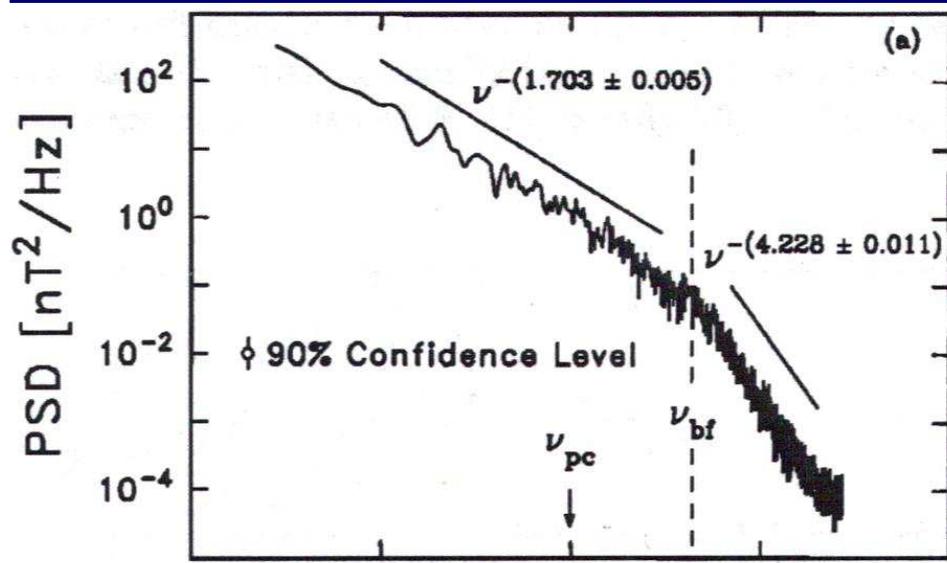
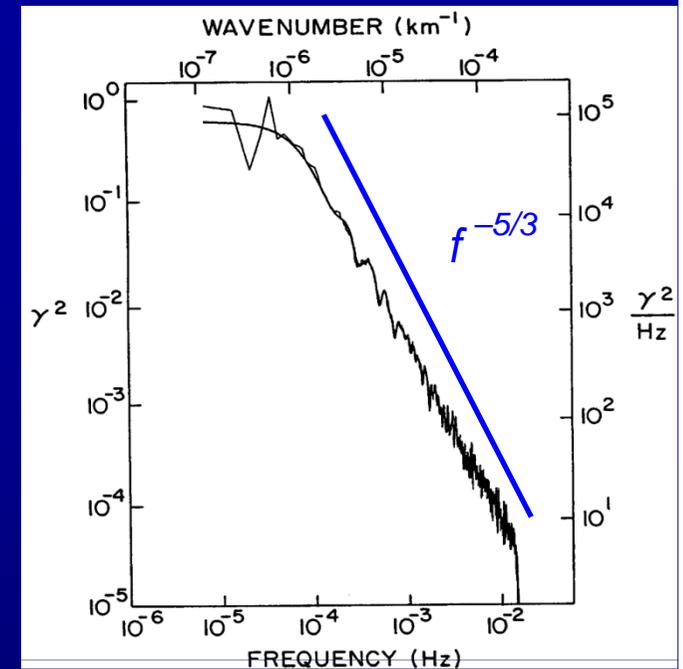
- Hydro: Scale invariance down to the dissipation scale $1/k_d$
- Collisionless Plasmas: - **Breaking of the scale invariance at $\rho_{i,e} d_{i,e}$**
- **Absence of the viscous dissipation scale $1/k_d$**

Solar wind turbulence

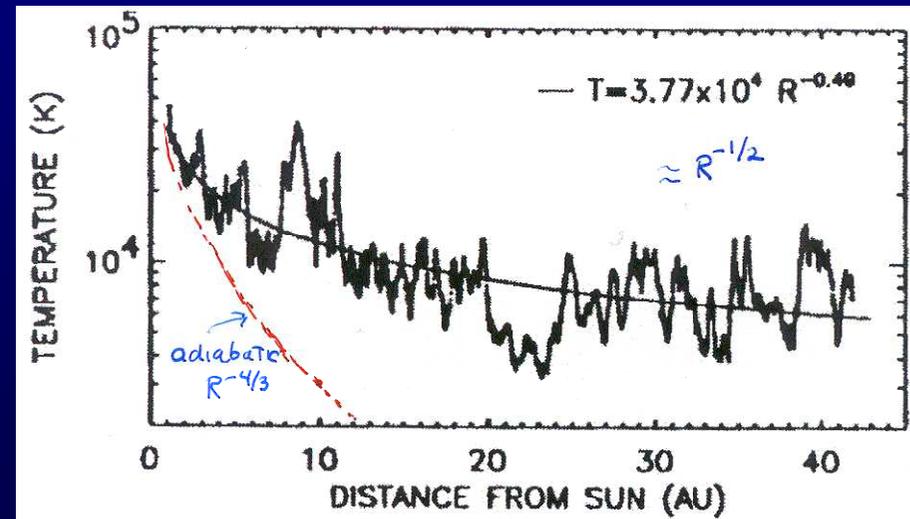
Typical power spectrum of magnetic energy at 1 AU

Does the energy cascade or dissipate below the ion scale ρ_i ?

Matthaeus & Goldstein, 82



Leamon *et al.* 98; Goldstein *et al.* JGR, 94

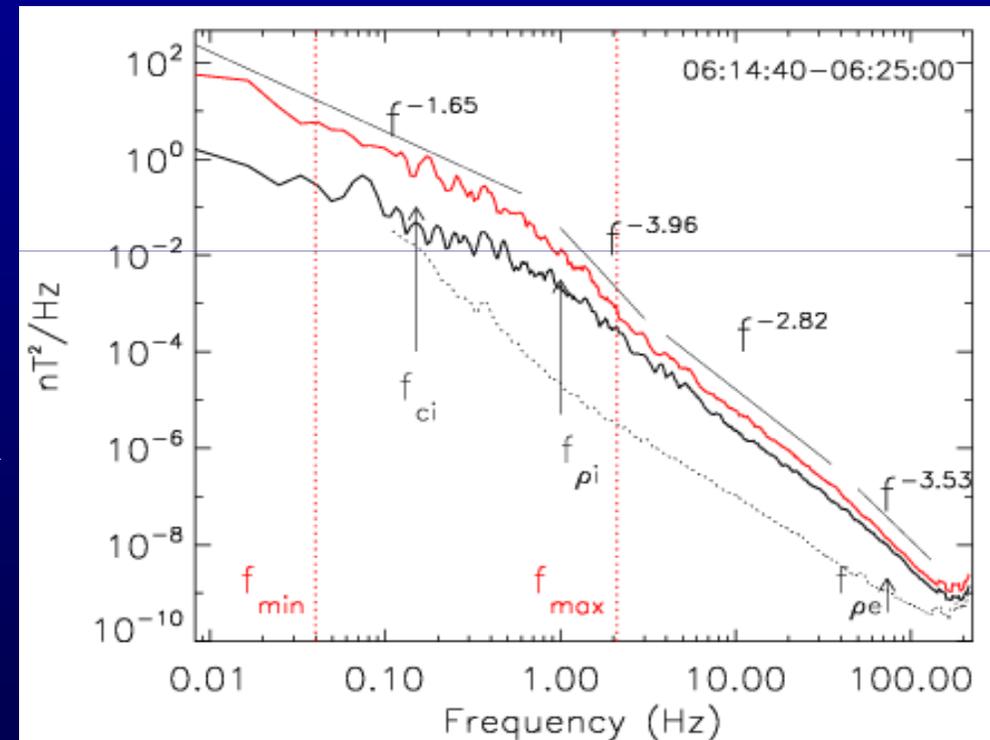


Richardson & Paularena,
GRL, 1995 (Voyager data)

How to analyse space turbulence ?

Turbulence theories generally predict **spatial spectra**: K41 ($k^{-5/3}$); IK ($k^{-3/2}$), Anisotropic MHD turbulence ($k_{\perp}^{-5/3}$), Whistler turbulence ($k^{-7/3}$), ...

Example of measured spectra in the SW



But measurements provide only **temporal spectra** (generally with different power laws at different frequencies)



How to infer *spatial spectra* from *temporal* ones measured in the spacecraft frame? $B^2 \sim \omega_{sc}^{-\alpha} \Rightarrow B^2 \sim k_{\parallel}^{-\beta} k_{\perp}^{-\gamma}$?

The Taylor frozen-in flow assumption

In the solar wind (SW) the Taylor's hypothesis can be valid
at MHD scales

High SW speeds: $V \sim 600\text{km/s} \gg V_\phi \sim V_A \sim 50\text{km/s} \Rightarrow$

$$\omega_{spacecraft} = \omega_{plasma} + \mathbf{k} \cdot \mathbf{V} \approx \mathbf{k} \cdot \mathbf{V} = k_V V$$

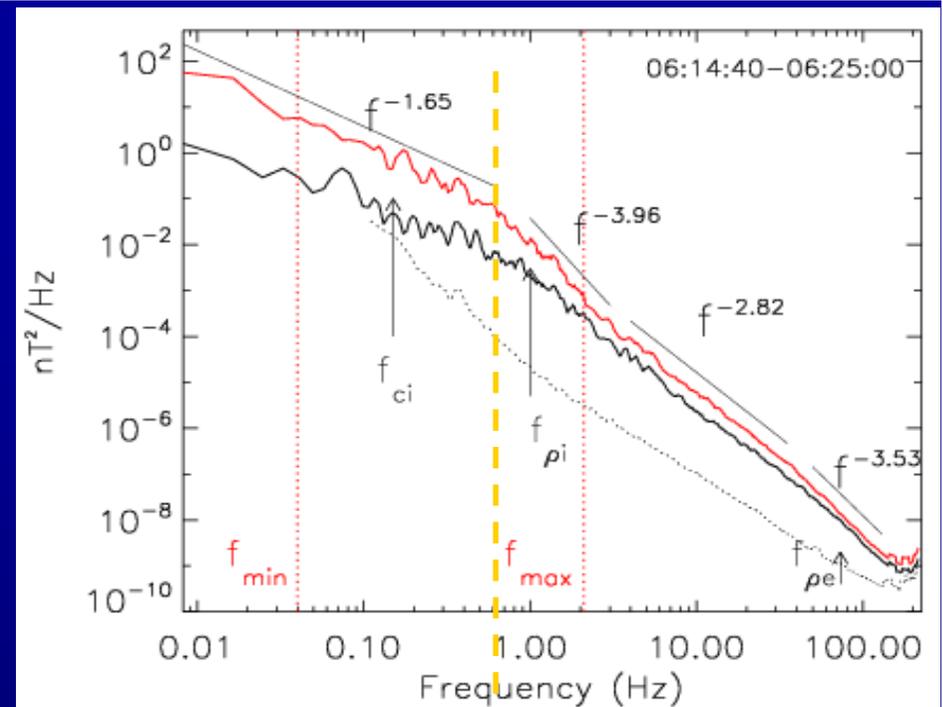
\Rightarrow Inferring the k -spectrum is possible with one spacecraft

But only along one single direction

1. At MHD scales, even if the Taylor assumption is valid, **inferring 3D k -spectra from an ω -spectrum is impossible**

2. At sub-ion and electron scales $V\phi$ can be larger than $V_{sw} \Rightarrow$ The **Taylor's hypothesis is invalid**

1 & 2 \Rightarrow ***Need to use multi-spacecraft measurements and appropriate methods to infer 3D k -spectra***



← MHD scales | Sub-ion scales →

Anisotropy and the critical balance conjecture

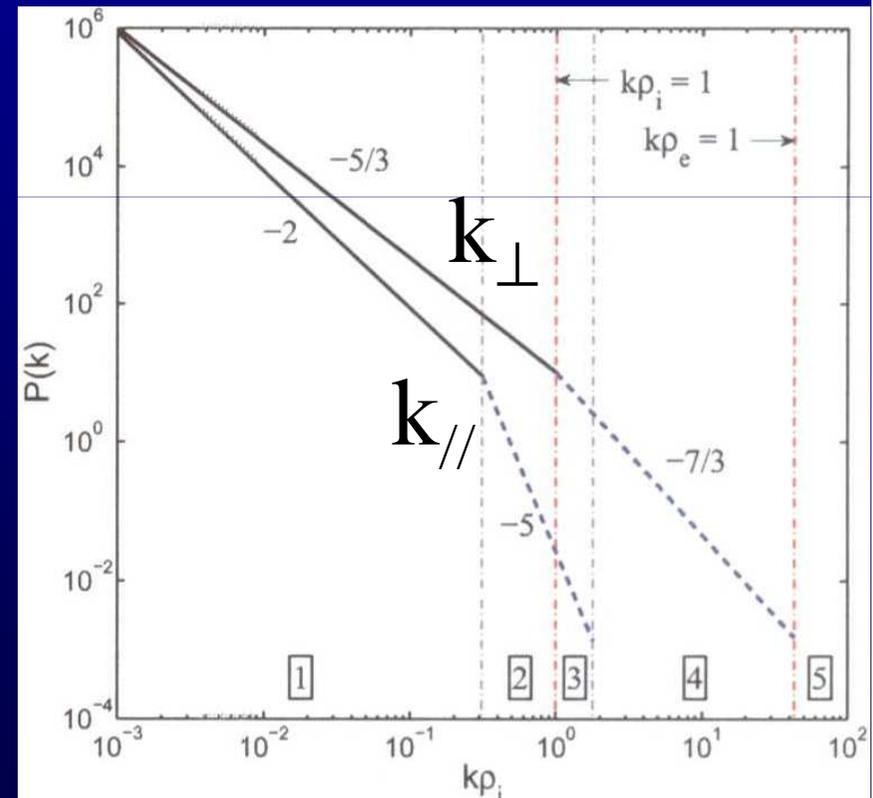
The critical balance conjecture [Goldreich & Sridhar, 1995]:

Linear (Alfvén) time \sim nonlinear (turnover) time

$$\Rightarrow \omega \sim k_{\parallel} V_A \sim k_{\perp} u_{\perp}$$

$$\Rightarrow k_{\parallel} \sim k_{\perp}^{2/3}$$

See also [Boldyrev, ApJ, 2005] and [Galtier et al., Phys. Plasmas, 2005]



[Chen et al., ApJ, 2010]

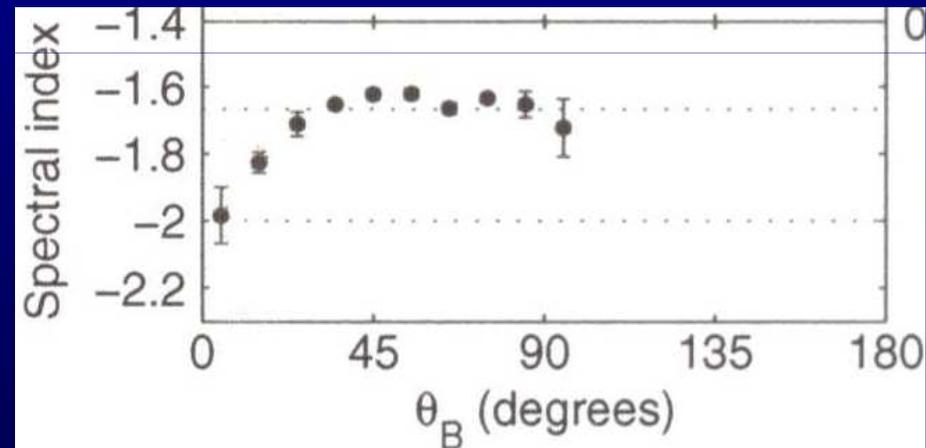
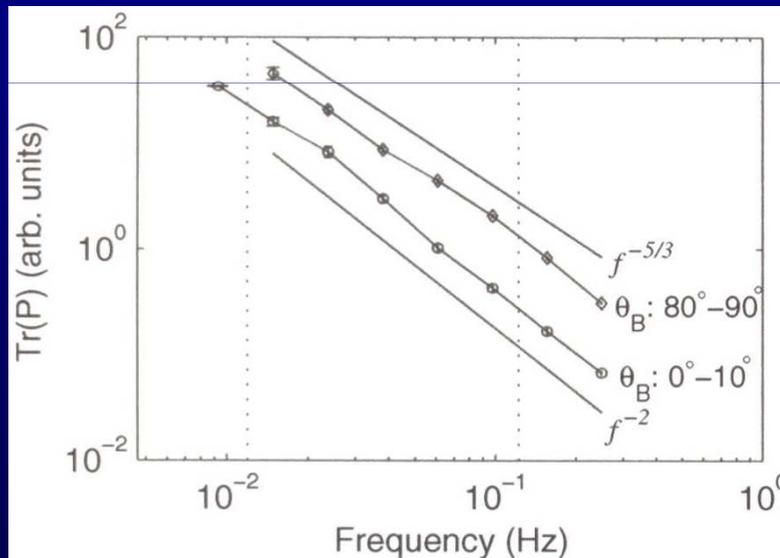
Single satellite analysis \rightarrow use of the Taylor assumption:

$$\omega_{sc} \sim \mathbf{k} \cdot \mathbf{V}_{sw} \sim k_{\parallel} V_{sw}$$

$$\mathbf{V} // \mathbf{B} \rightarrow k_{\parallel} = k_{\parallel}$$

$$\mathbf{V} \perp \mathbf{B} \rightarrow k_{\parallel} = k_{\perp}$$

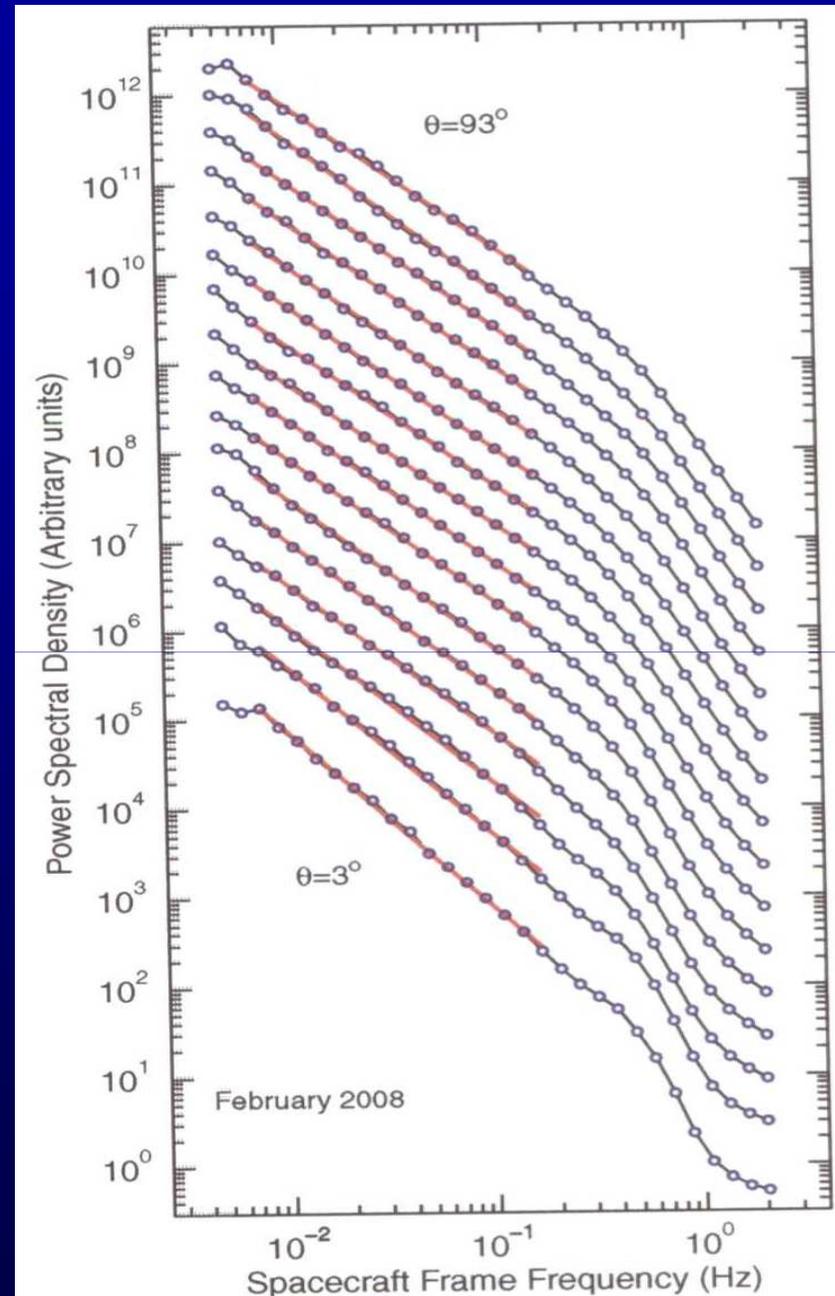
*Assumes axisymmetry
around B*



$\Theta_{BV} \rightarrow 0 \Rightarrow B^2 \sim k_{\parallel}^{-2} \Rightarrow$ *Partial* evidence of the critical balance [Horbury et al., PRL, 2008]

Results confirmed by
Podesta, ApJ, 2009

See also Chen et al.,
PRL, 2010

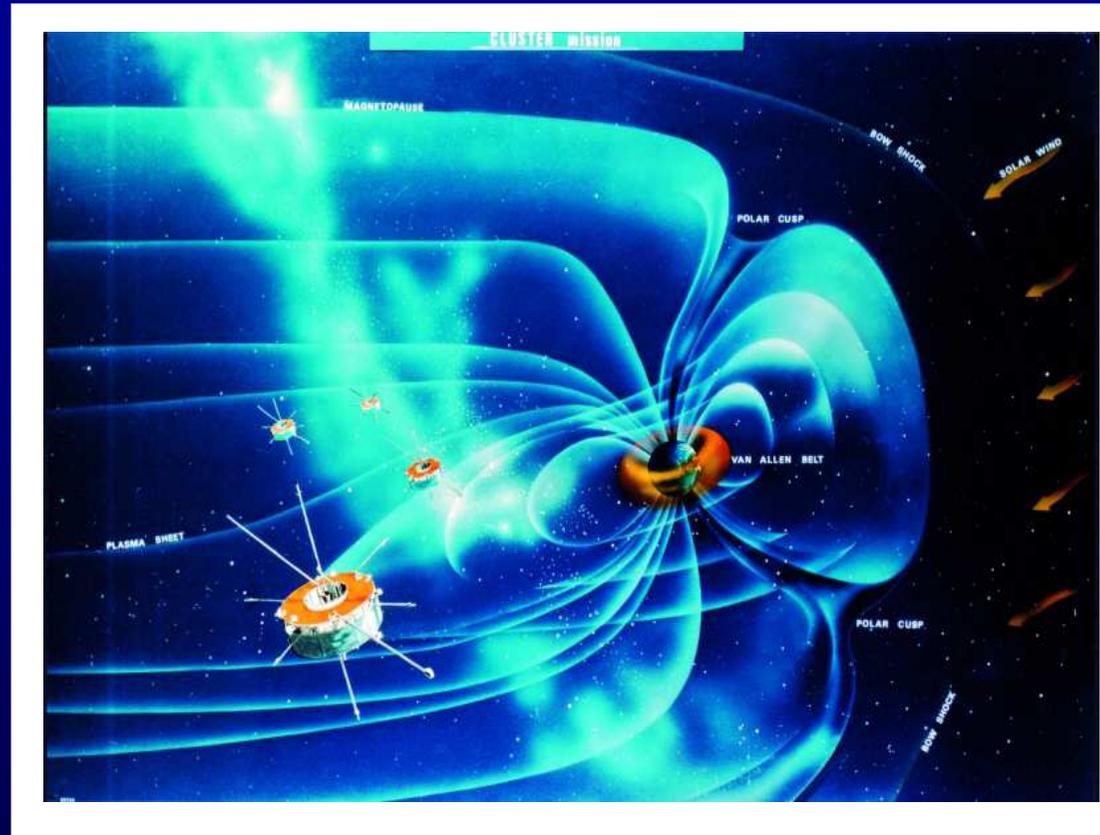


The ESA/Cluster mission

The first multispacecraft mission: 4 identical satellites

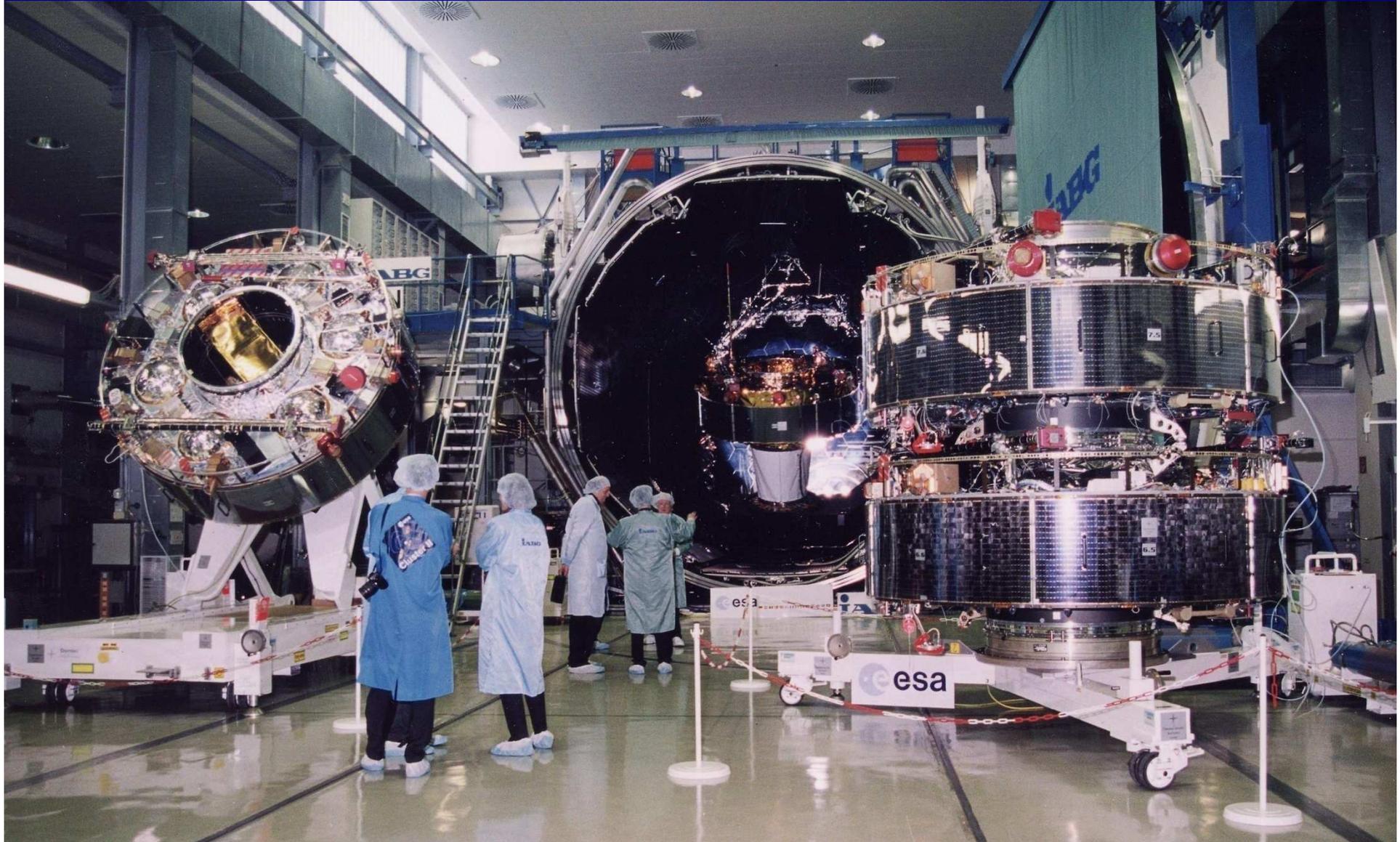
Objetives:

- **3D exploration** of the Earth magnetosphere boundaries (magnetopause, bow shock, magnetotail) & SW
- **Mesurements of 3D quantities:**
 $\mathbf{J}=\nabla\times\mathbf{B}$, ...
- **Fundamental physics:**
turbulence, reconnection, particle acceleration, ...



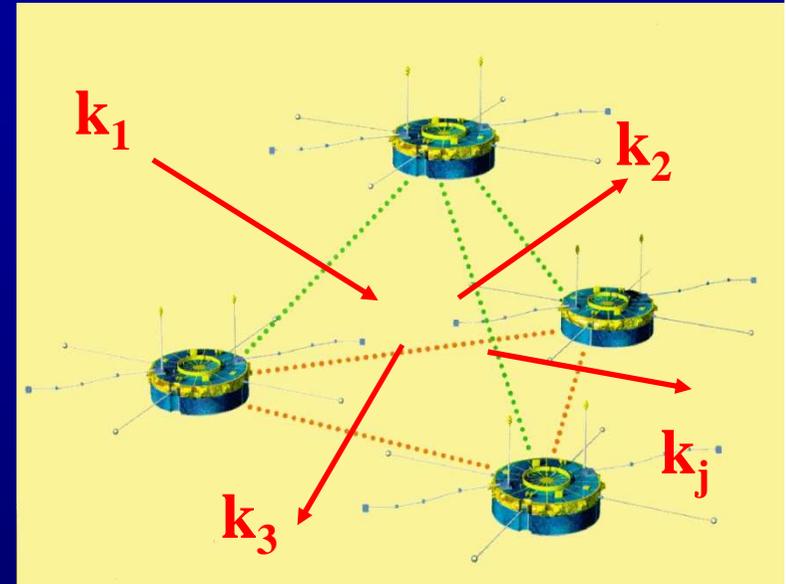
Different orbits and separations (10^2 to 10^4 km) depending on the scientific goal

The 4 satellites before launch



The k -filtering technique

Interferometric method: it provides, by using a NL filter bank approach, an optimum estimation of the 4D spectral energy density $P(\omega, \mathbf{k})$ from simultaneous multipoints measurements [Pinçon & Lefeuvre; Sahraoui et al., 03, 04, 06, 10; Narita et al., 03, 06, 09]



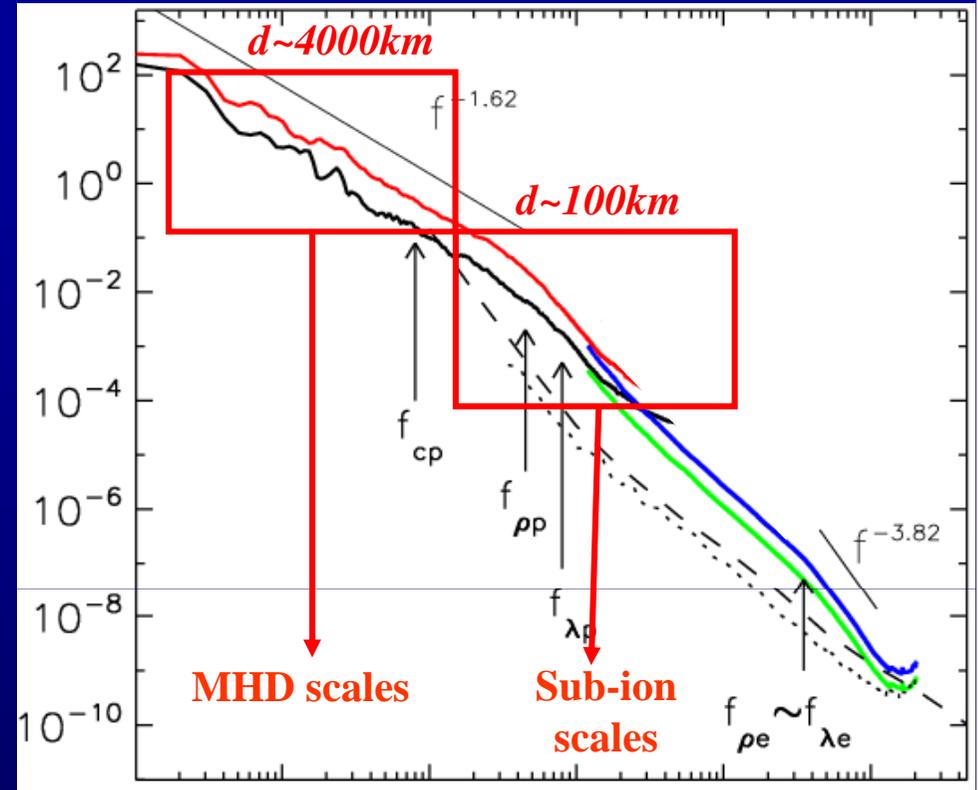
We use $P(\omega, \mathbf{k})$ to calculate

1. 3D ω - \mathbf{k} spectra \Rightarrow plasma mode identification e.g. Alfvén, whistler
2. 3D \mathbf{k} -spectra (anisotropies, scaling, ...)

Measurable spatial scales

Given a spacecraft separation d
 only one decade of scales $2d < \lambda < 30d$ can be correctly determined

- $\lambda_{min} \cong 2d$, otherwise *spatial aliasing* occurs.
- $\lambda_{max} \cong 30d$, because larger scales are subject to important uncertainties



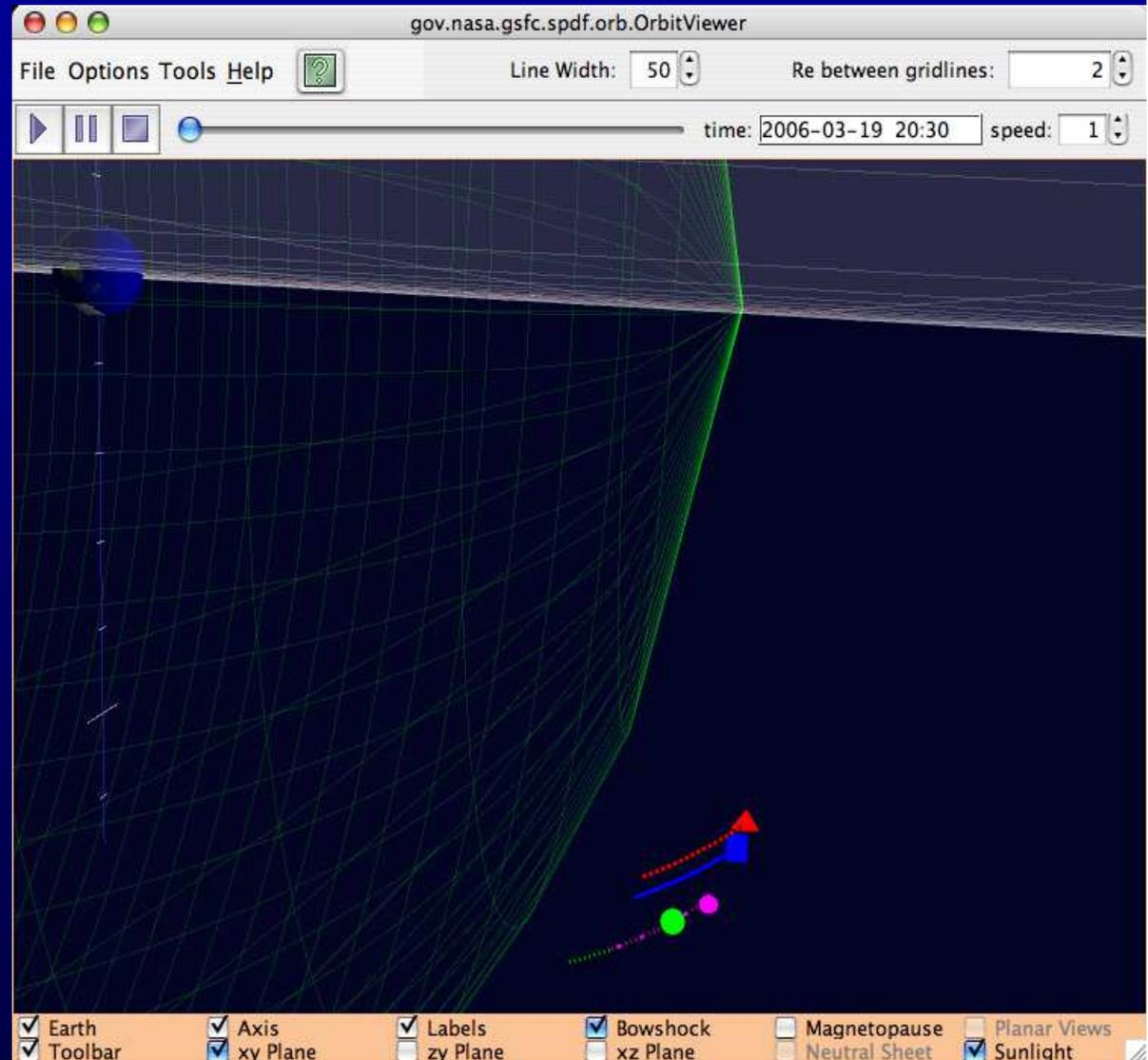
$$\omega_{\text{sat}} \sim kV \Rightarrow f_{\text{max}} \sim k_{\text{max}} V / \lambda_{\text{min}} \quad (V \sim 500 \text{ km/s})$$

- $d \sim 10^4 \text{ km} \Rightarrow$ MHD scales
- $d \sim 10^2 \text{ km} \Rightarrow$ Sub-ion scales
- $d \sim 1 \text{ km} \Rightarrow$ Electron scales (*but not accessible with Cluster: $d > 100$*)

1- MHD scale solar wind turbulence

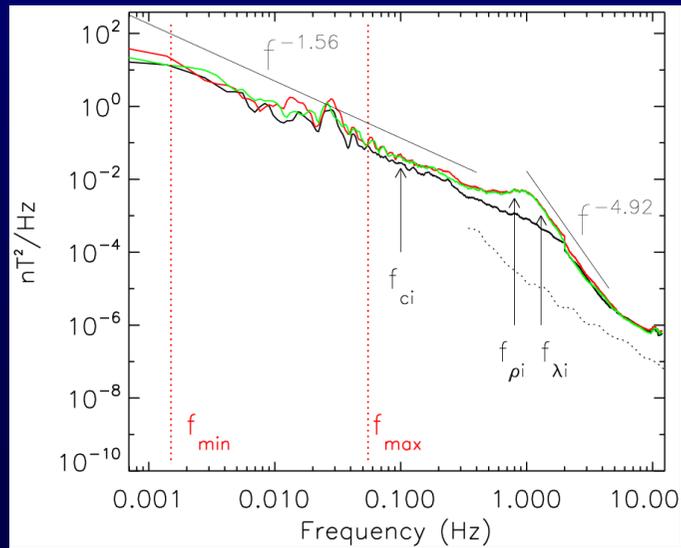
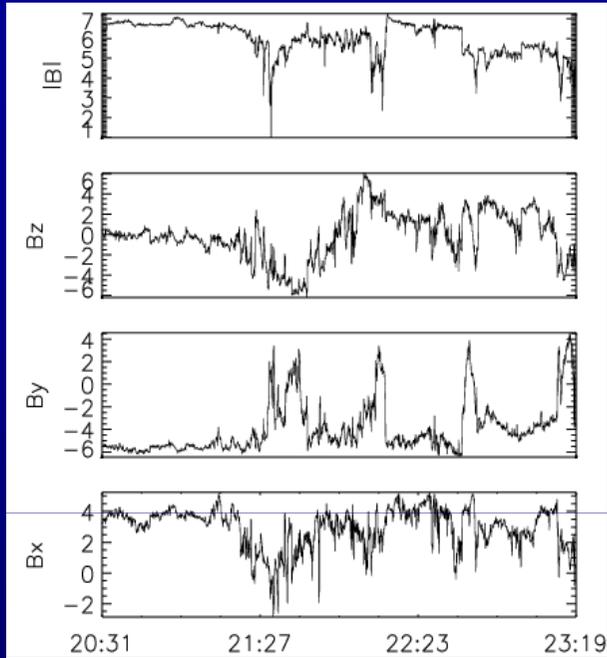
Position of the Quartet
on March 19, 2006

Satellite	Color	X	Y	Z
Cluster-1	Blue	15.038	-6.569	-9.299
Cluster-2	Red	15.139	-7.034	-8.672
Cluster-3	Green	13.979	-7.397	-10.41
Cluster-4	Magenta	14.587	-7.292	-9.987

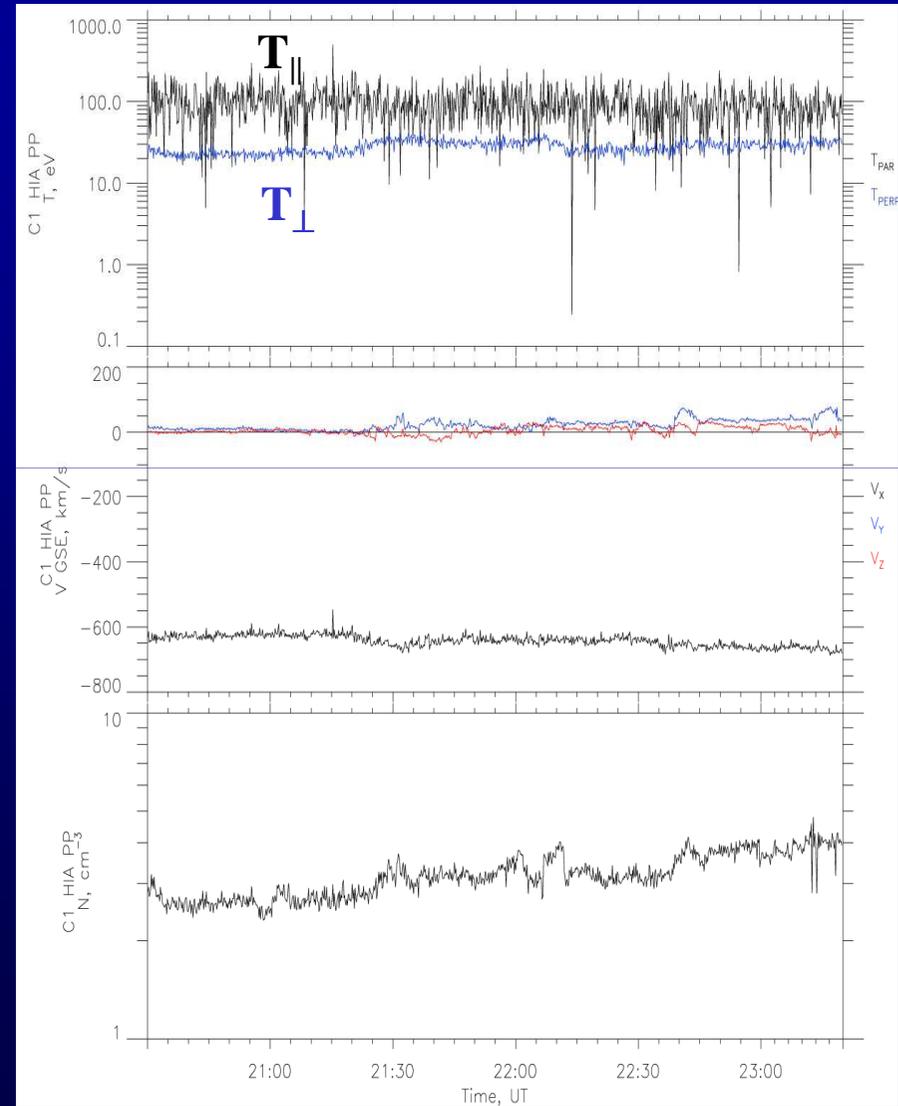


Data overview

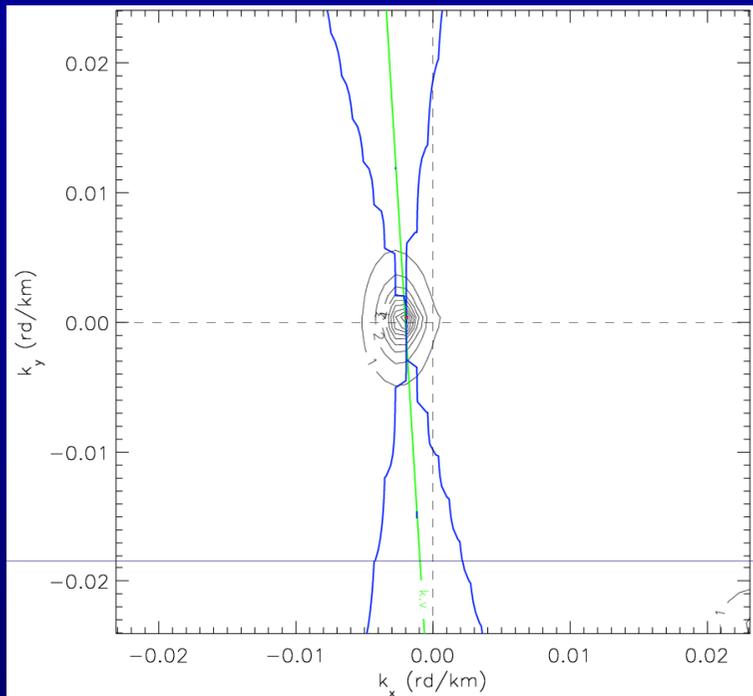
FGM data (CAA, ESA)



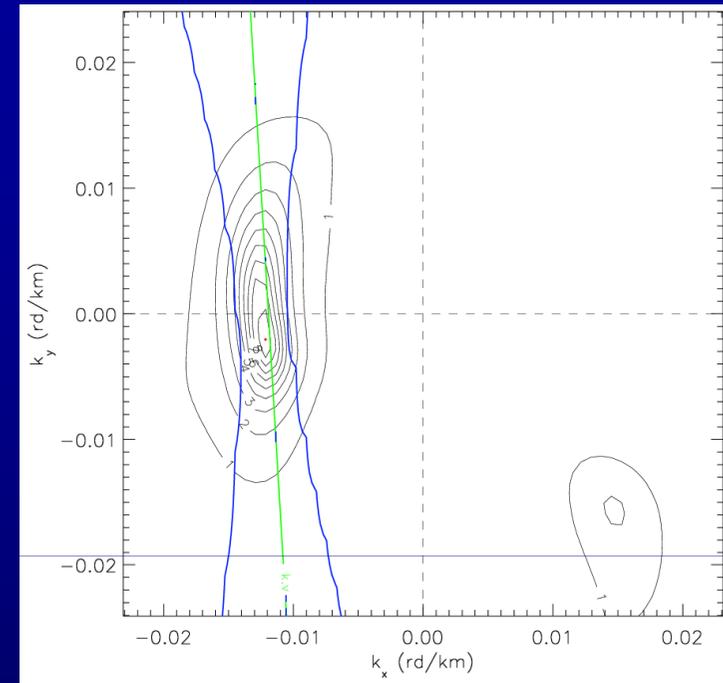
Ion plasma data from CIS (AMDA, CESR)



$$f_1 = 0.23 \text{ Hz} \sim 2f_{ci}$$



$$f_2 = 0.9 \text{ Hz} \sim 6f_{ci}$$



To compute reduced spectra we integrate over

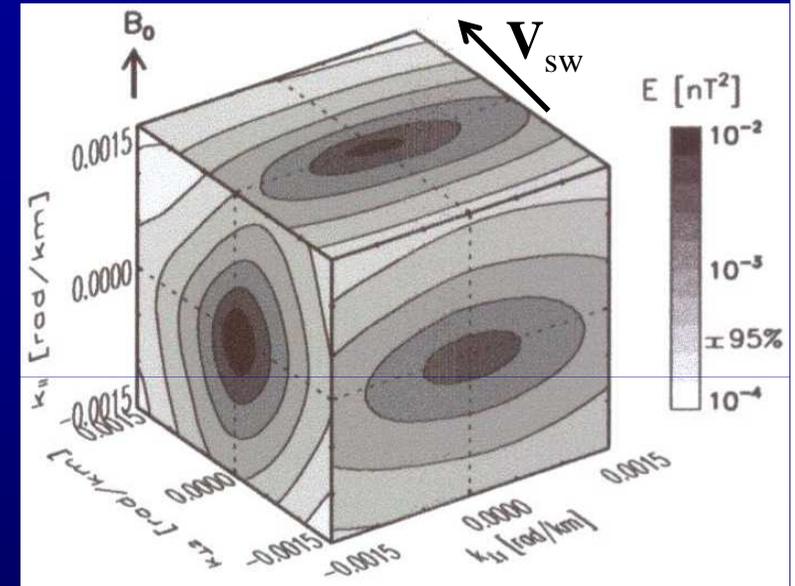
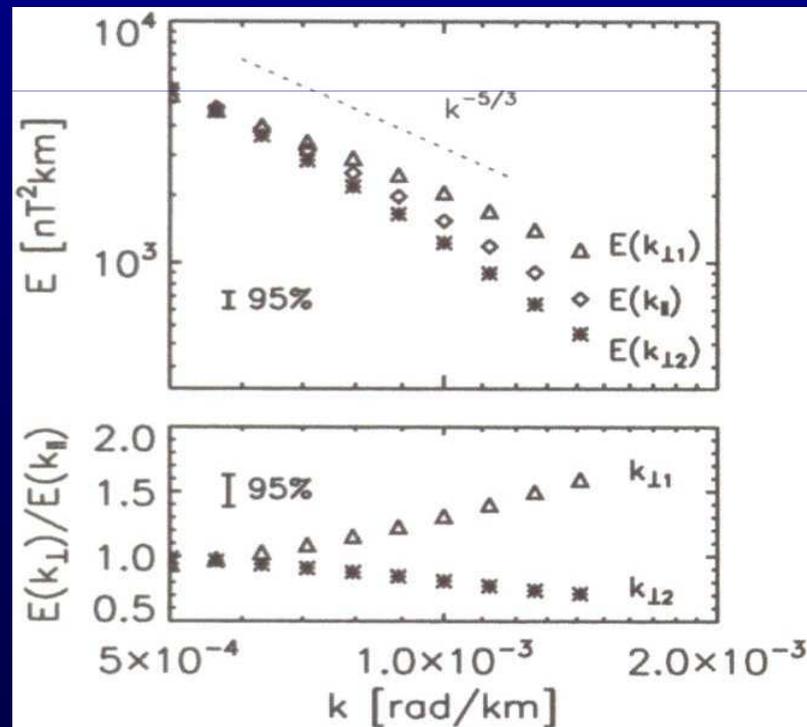
1. all frequencies f_{sc} :
$$\tilde{P}(\mathbf{k}) = \sum_{k_y, k_z} P(f_{sc}, \mathbf{k})$$

2. all $k_{i,j}$:
$$\tilde{\tilde{P}}(k_x) = \sum_{k_y, k_z} \tilde{P}(k_x, k_y, k_z)$$

Anisotropy of MHD turbulence along B_0 and V_{sw}

[Narita et al. , PRL, 2010]

Turbulence is not axisymmetric (around B) [see also Sahraoui, PRL, 2006]



The anisotropy ($\perp B$) is along $V_{sw} \rightarrow SW$ expansion effect ? [Saur & Bieber, JGR, 1999]

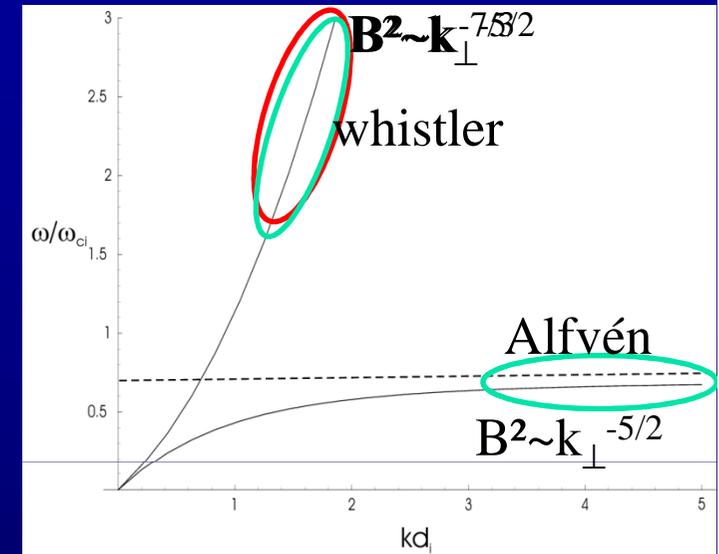
Kinetic (sub-ion scale) turbulence in the SW

I- Theoretical predictions on small scale turbulence

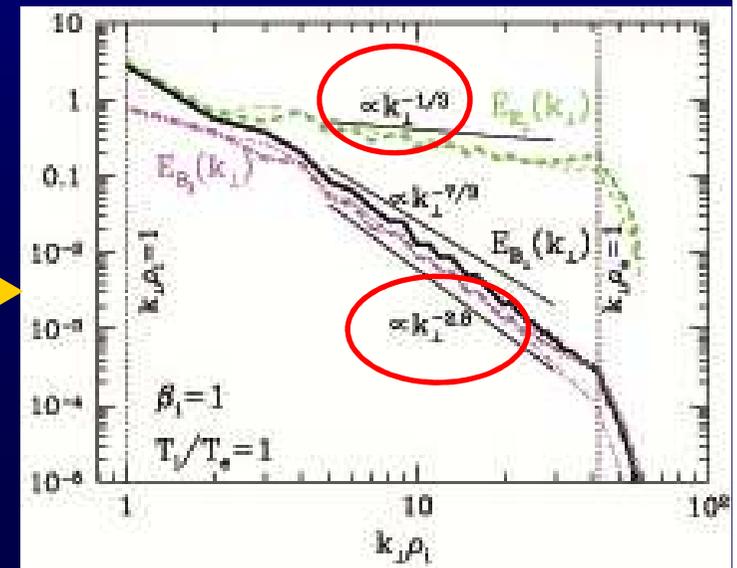
$$\mathbf{E} + \mathbf{V} \times \mathbf{B} = \frac{1}{en} \mathbf{J} \times \mathbf{B} - \frac{\nabla P_e}{en} + \dots$$

1. Fluid models (Hall-MHD) \rightarrow

- Whistler turbulence (E-MHD): (Biskamp *et al.*, 99, Galtier, 08)
- Weak Turbulence of Hall-MHD (Galtier, 06; Sahaoui *et al.*, 07)



2. Gyrokinetic theory: $k_{\parallel} \ll k_{\perp}$ and $\omega \ll \omega_{ci}$ \rightarrow (Schekochihin *et al.* 06; Howes *et al.*, 11)



Other numerical predictions on electron scale turbulence

2D PIC simulations gave evidence of a power law dissipation range at $k\rho_e > 1$

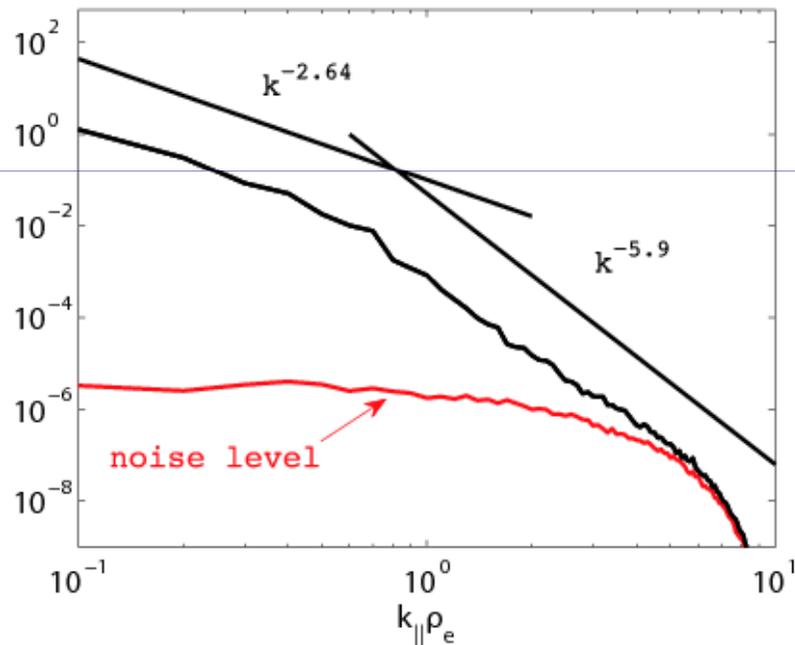


Figure 4. Spectrum of magnetic fluctuation $|\delta B|^2/|B_0|^2$ in the parallel direction $k\rho_{e\parallel}$. The noise level curve is in red. The power-law best fits are superimposed. (A color version of this figure is available in the online journal.)

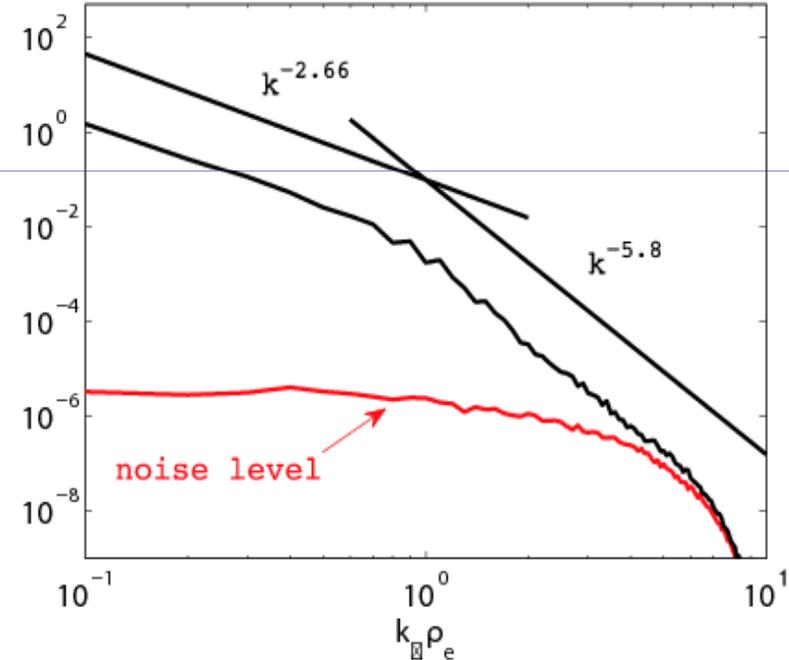
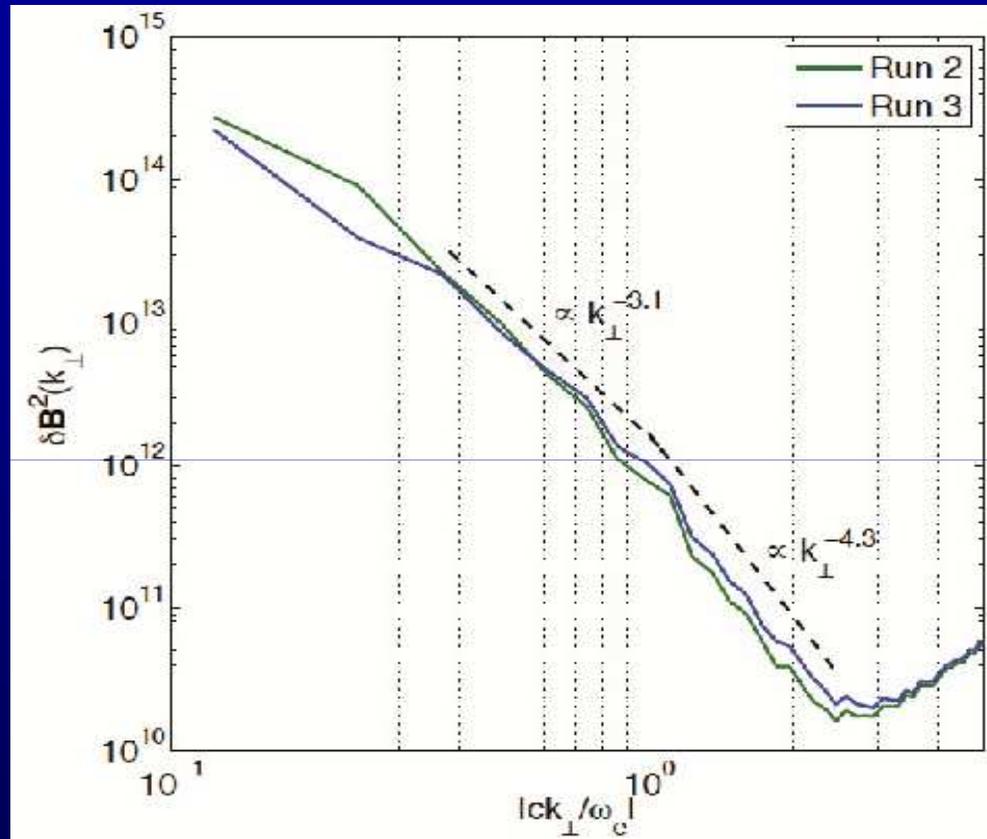


Figure 5. Spectrum of magnetic fluctuations $|\delta B|^2/|B_0|^2$ in the perpendicular direction $k\rho_{e\perp}$. The noise level curve is in red. The power-law best fits are superimposed.

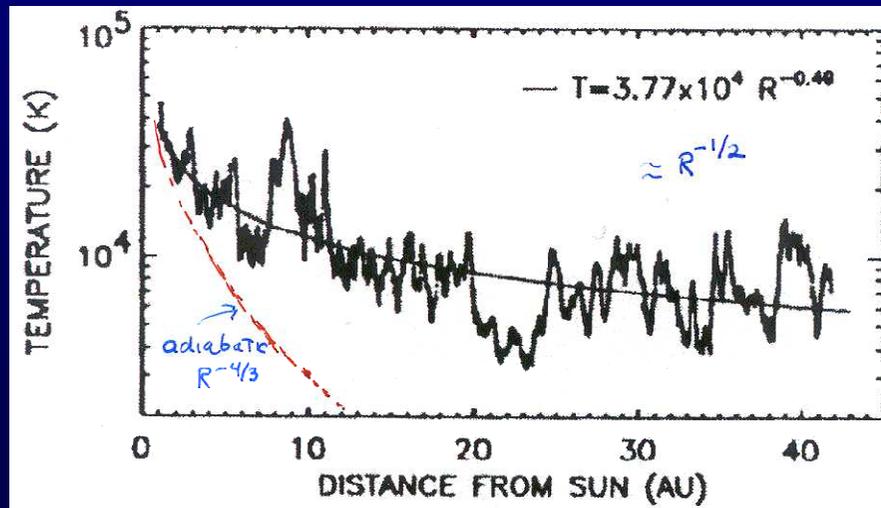
3D PIC simulations of whistler turbulence : $k^{-4.3}$ at $kd_e > 1$



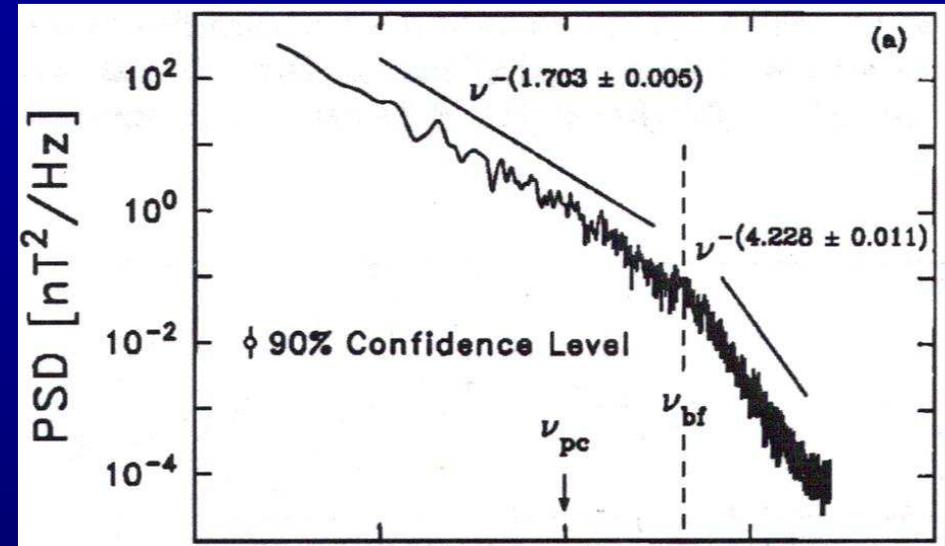
Chang & Gary, GRL 2011

2- Observations of kinetic SW turbulence

What happens to the energy at, and below, the ion scale ρ_i (not f_{ci}): a total dissipation or a new cascade?



Richardson & Paularena, GRL, 1995 (Voyager data)

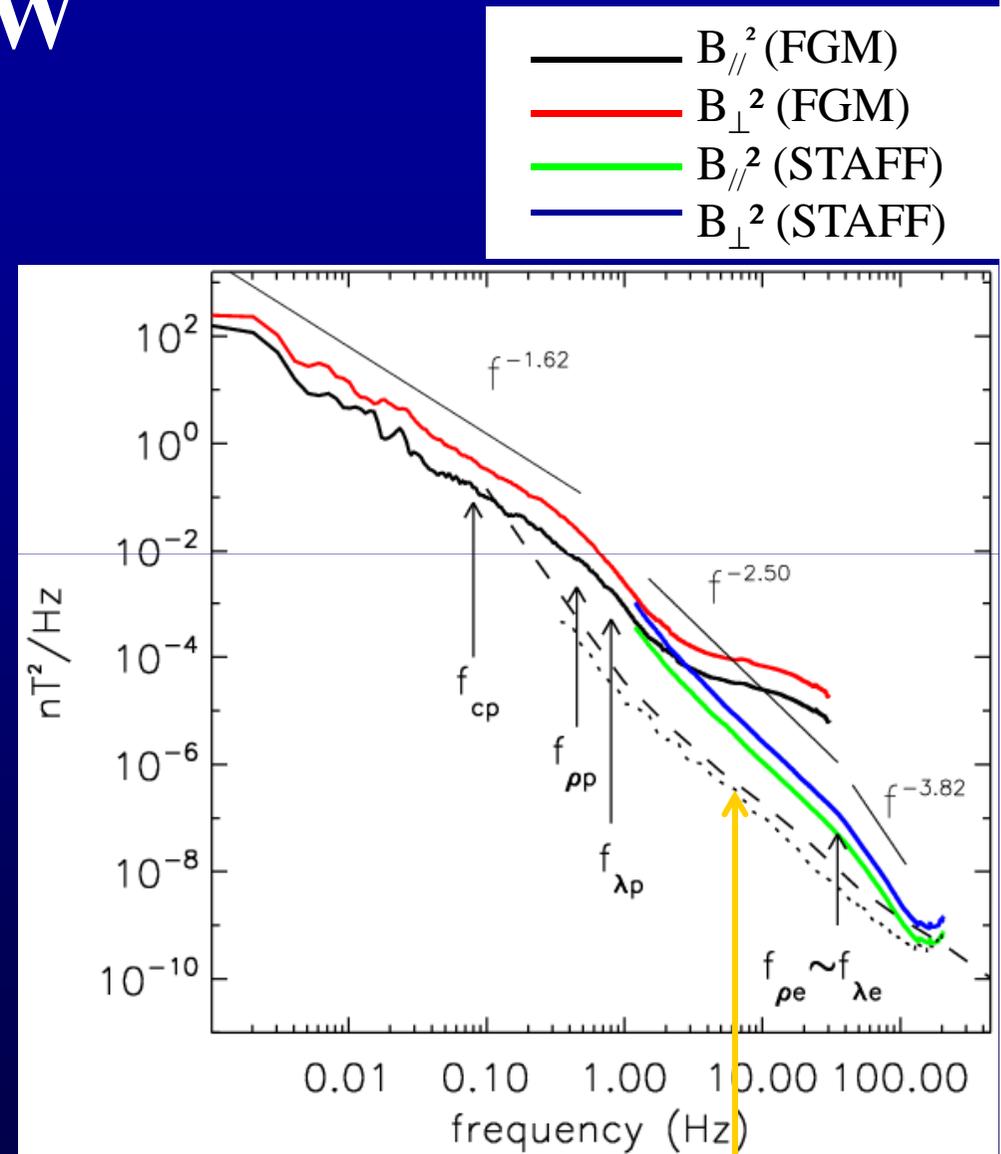


Leamon *et al.*, 98; Goldstein *et al.*, 94

First evidence of a cascade from MHD to electron scale in the SW

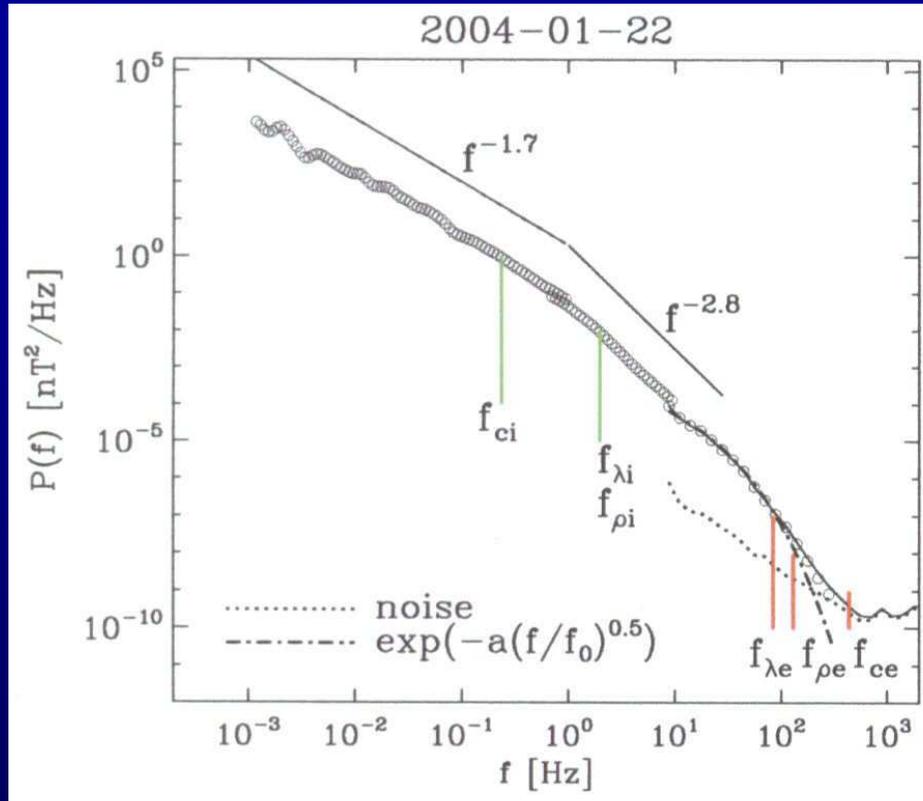
1. Two breakpoints corresponding to ρ_i and ρ_e are observed.
2. A clear evidence of a new inertial range $\sim f^{-2.5}$ below ρ_i
3. *First evidence of a dissipation range $\sim f^{-4}$ near the electron scale ρ_e*

Sahraoui et al., PRL, 2009

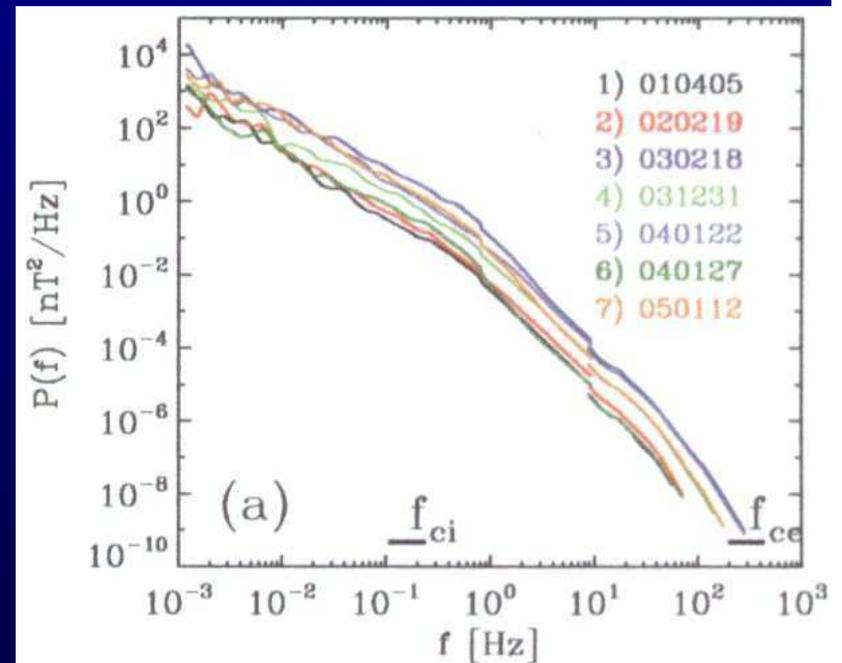


STAFF-SC sensitivity floor

[Alexandrova et al., 2009, 2013]



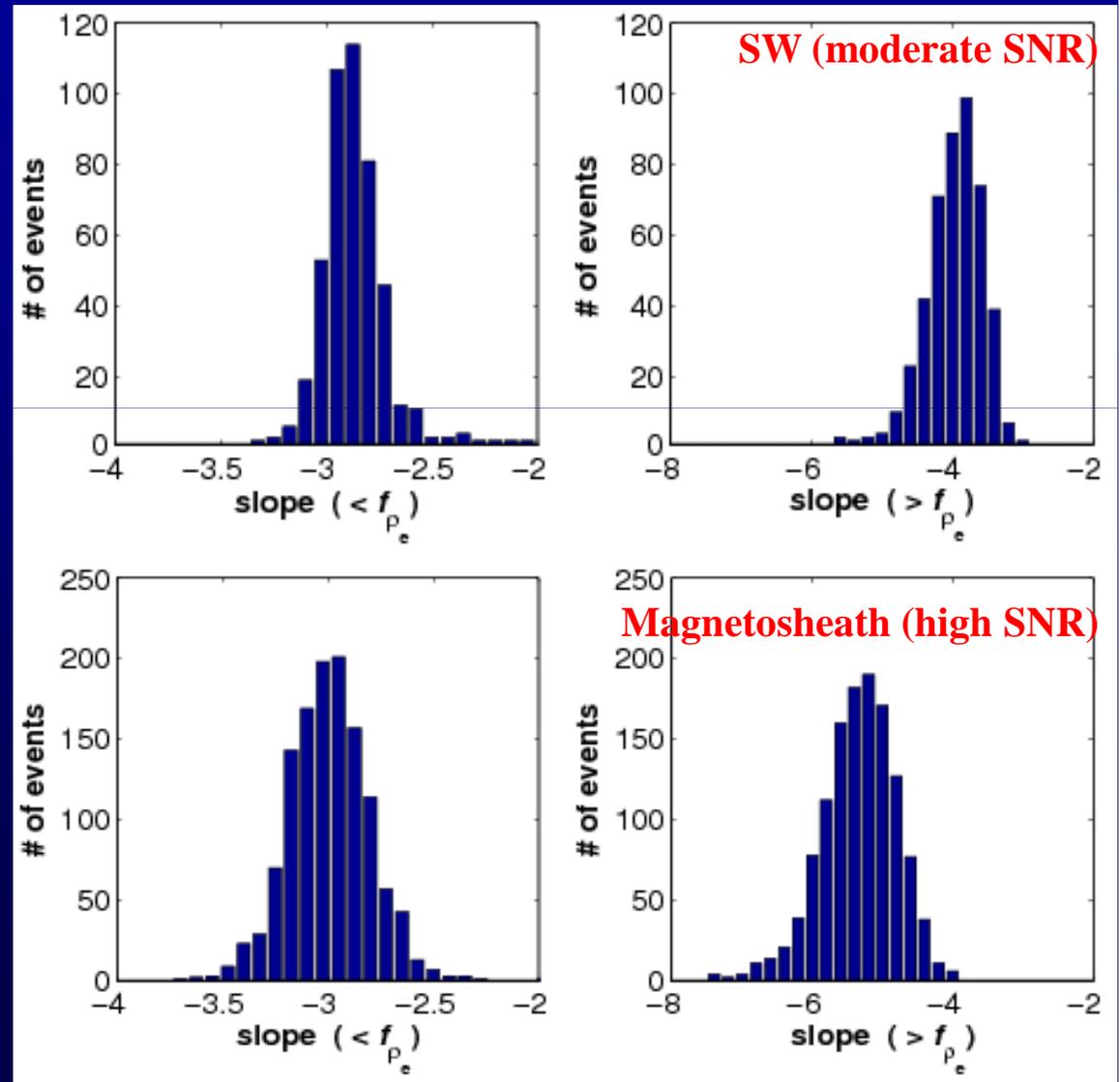
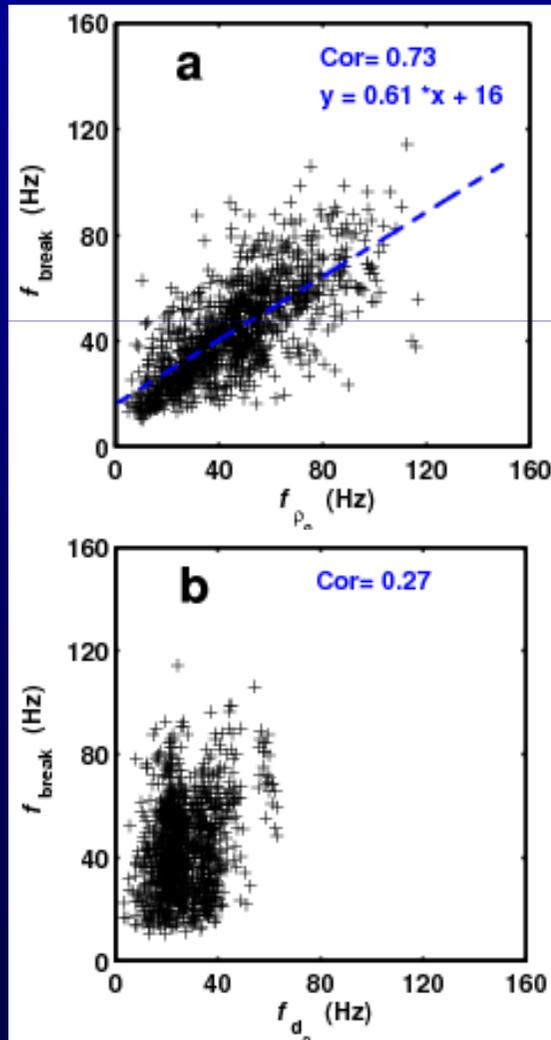
Similar observations from STAFF-SA data, but the spectra were fit by an exponential model



The largest survey of Cluster/STAFF-SC data

A better correlation
with ρ_e than with d_e

[Sahraoui+, 2013; Huang+, 2013]

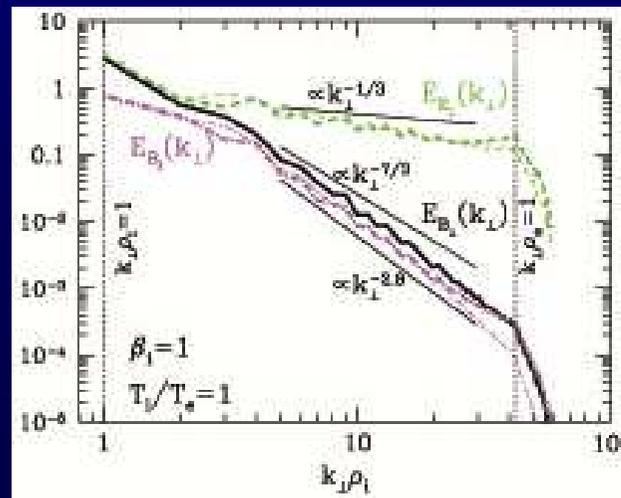
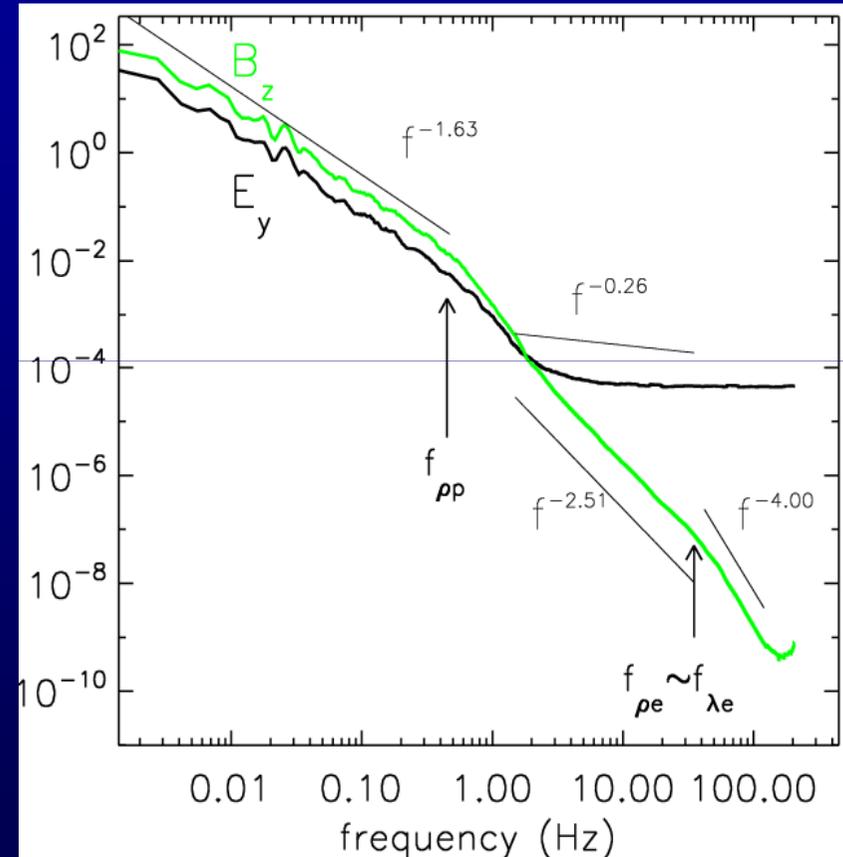


Whistler or KAW turbulence?

1. Large (MHD) scales ($L > \rho_i$): **strong correlation of E_y and B_z** in agreement with $\mathbf{E} = -\mathbf{V} \times \mathbf{B}$
2. Small scales ($L < \rho_i$): **steepening of B^2 and enhancement of E^2** (however, strong noise in E_y for $f > 5\text{Hz}$)

⇒ **Good agreement with GK theory of Kinetic Alfvén Wave turbulence**

FGM, STAFF-SC and EFW data



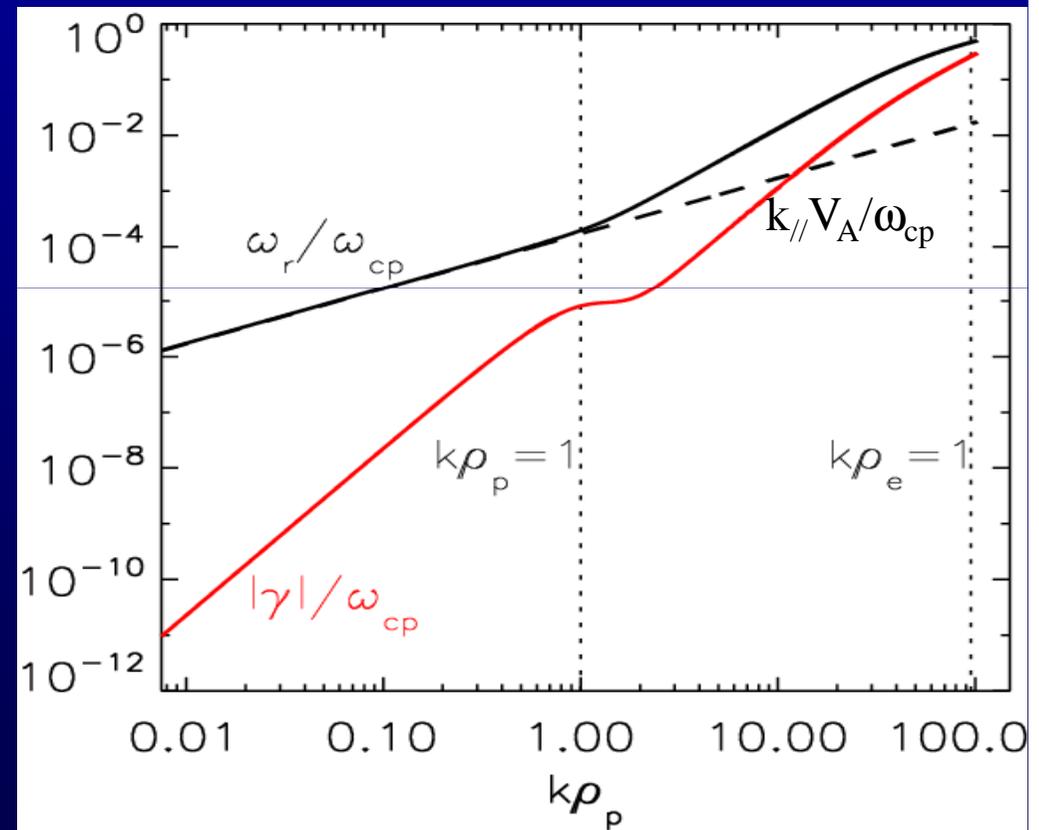
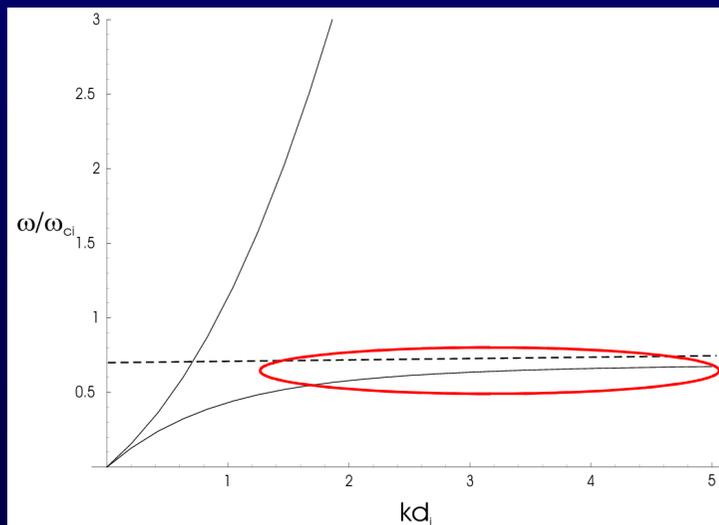
Howes *et al.*
PRL, 11

See also Bale *et al.*,
PRL, 2005

Theoretical interpretation : KAW turbulence

Linear Maxwell-Vlasov solutions: $\Theta_{\text{KB}} \sim 90^\circ$, $\beta_i \sim 2.5$, $T_i/T_e \sim 4$

The **Kinetic Alfvén Wave** solution extends **down to $k\rho_e \sim 1$** with $\omega_r < \omega_{ci}$
 [See also Podesta, ApJ, 2010]



$$\omega_r = k_{\parallel} V_A k_{\perp} \rho_i / \sqrt{\beta_i + 2 / (1 + T_i / T_e)}$$

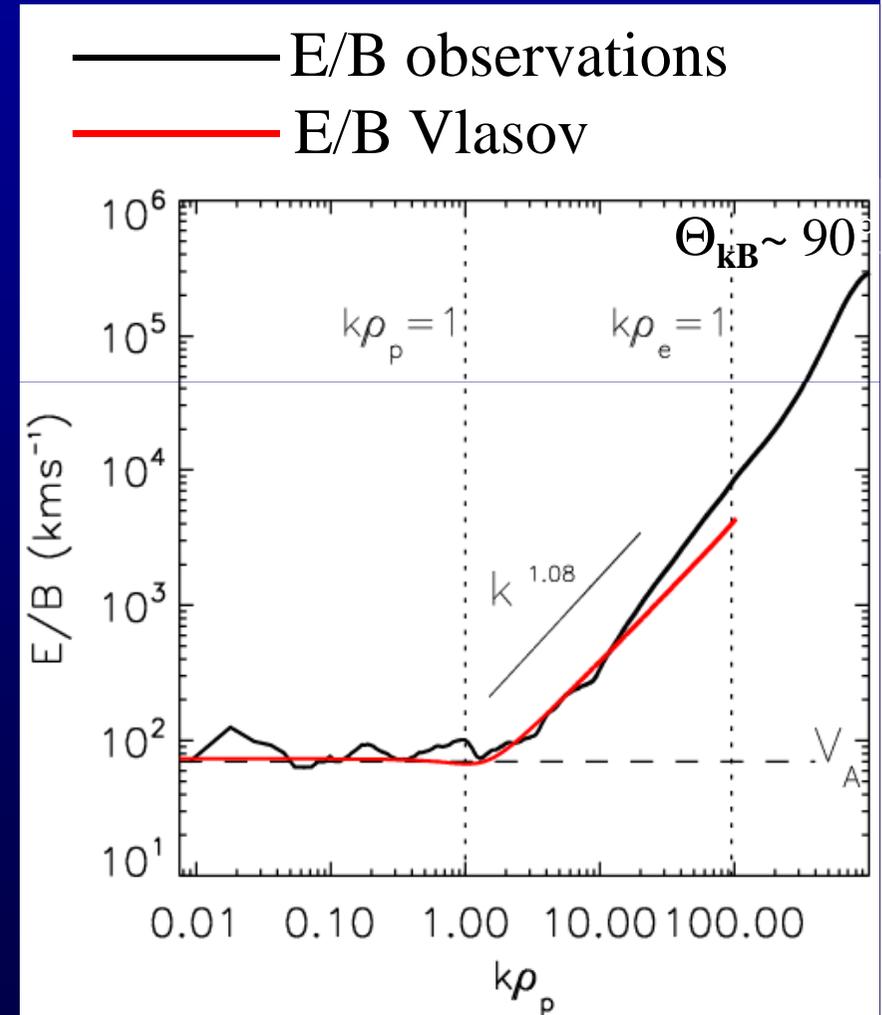
E/B : KAW theory vs observations

$$\omega_r = k_{//} V_A k_{\perp} \rho_i / \sqrt{\beta_i + 2 / (1 + T_i / T_e)}$$

➤ Lorentz transform: $\mathbf{E}_{\text{sat}} = \mathbf{E}_{\text{plas}} + \mathbf{V} \times \mathbf{B}$

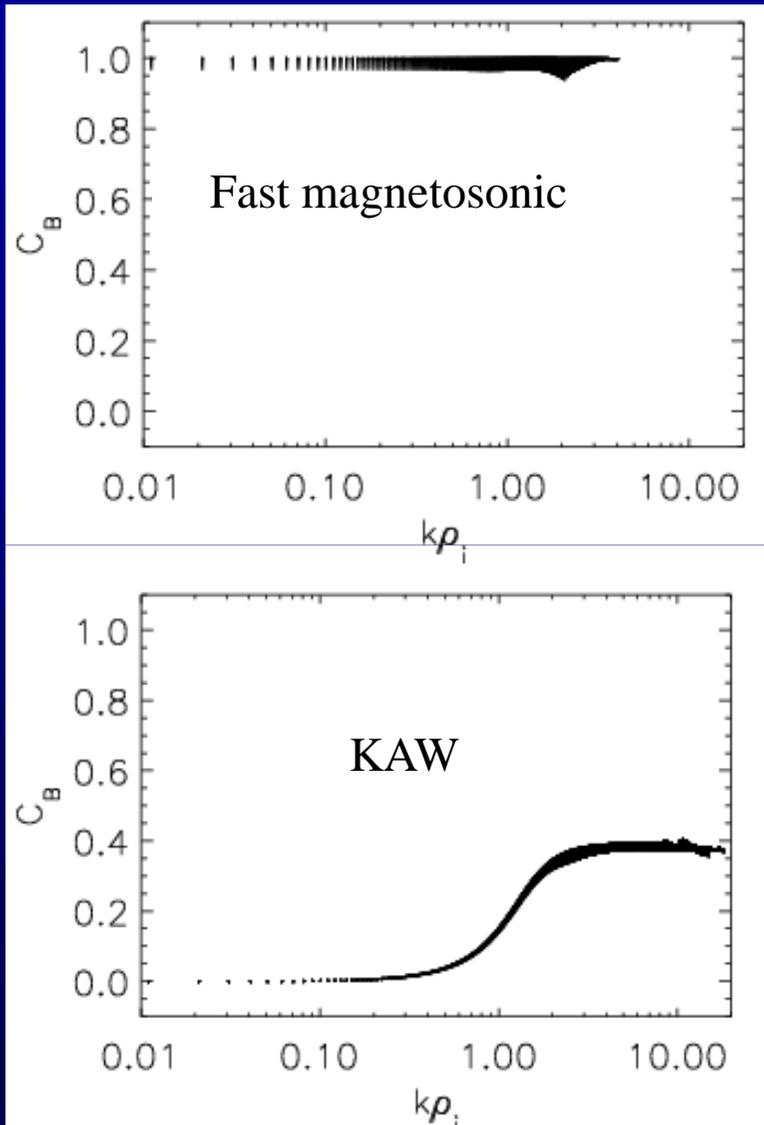
➤ Taylor hypothesis to transform the spectra from f (Hz) to $k\rho_i$

1. Large scale ($k\rho_i < 1$): $\delta E / \delta B \sim V_A$
2. Small scale ($k\rho_i > 1$): $\delta E / \delta B \sim k^{1.1} \Rightarrow$
in agreement with GK theory of KAW turbulence $\delta E^2 \sim k_{\perp}^{-1/3}$ & $\delta B^2 \sim k_{\perp}^{-7/3} \Rightarrow \delta E / \delta B \sim k$
3. The departure from linear scaling ($k\rho_i \gtrsim 10$) is due to noise in Ey data



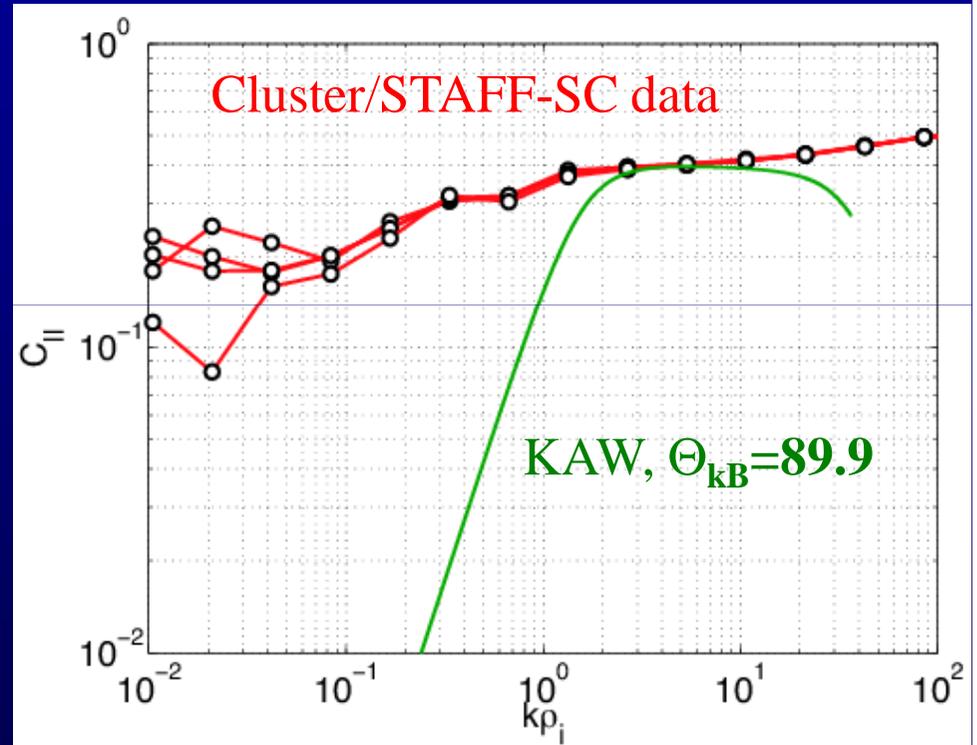
Sahraoui et al., PRL, 2009

Magnetic compressibility



[Sahraoui+, ApJ, 2012]

Additional evidence of KAW
at $k\rho_i \gtrsim 1$



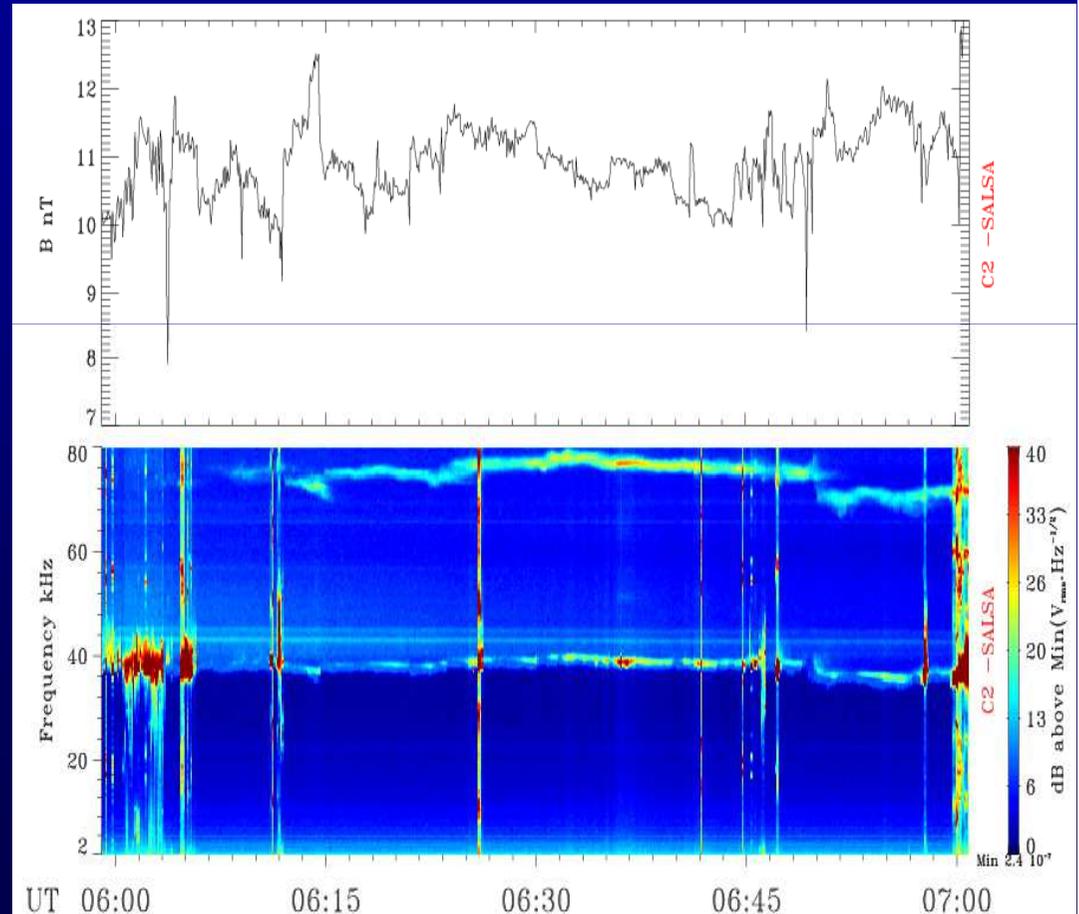
[Kiyani+, ApJ, 2012; Podesta+, 2012]

3D k -spectra at sub-proton scales of SW turbulence

Conditions required:

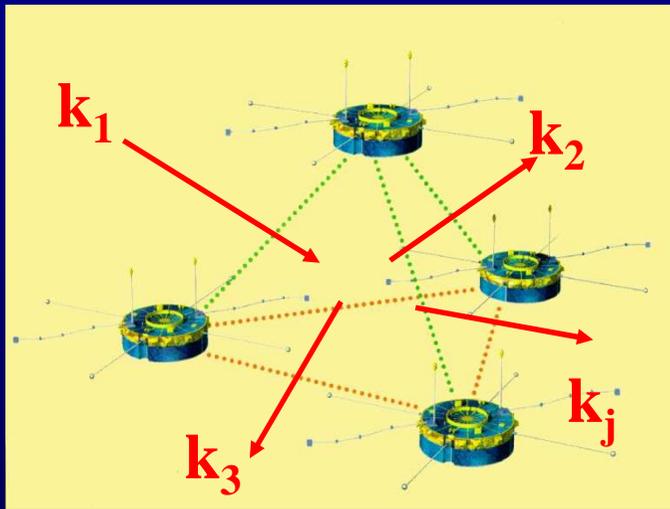
1. Quiet SW: NO electron foreshock effects
2. Shorter Cluster separations (~ 100 km) to analyze sub-proton scales
3. Regular tetrahedron to infer actual 3D k -spectra [Sahraoui et al., JGR, 2010]
4. High SNR of the STAFF data to analyse HF (>10 Hz) SW turbulence.

20040110, 06h05-06h55

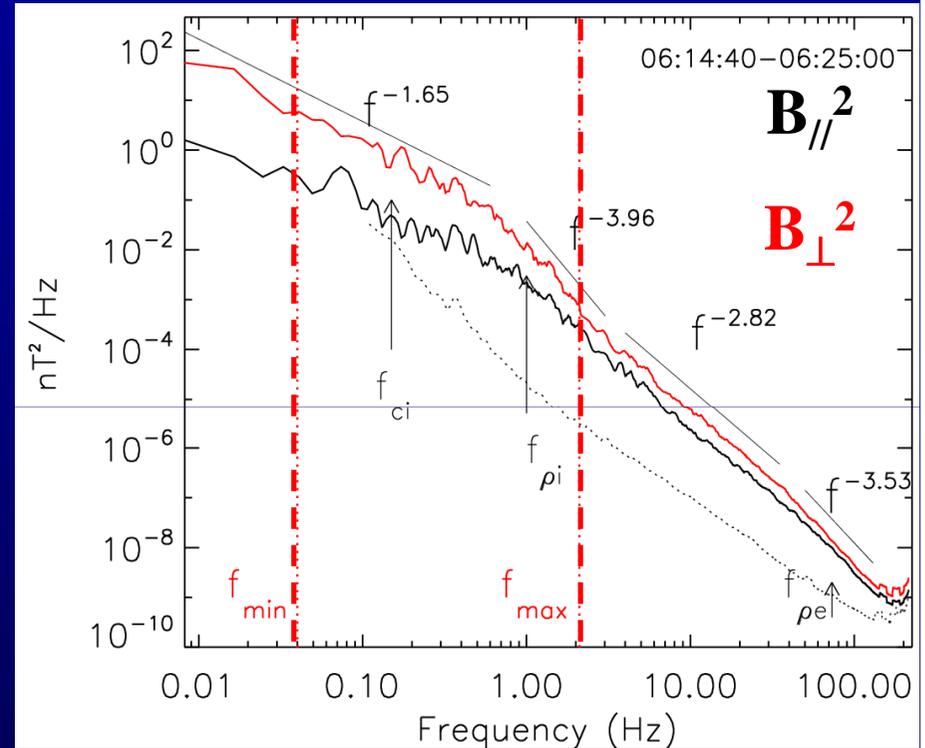


3D k -spectra at sub-proton scales

We use the k -filtering technique to estimate the 4D spectral energy density $P(\omega, \mathbf{k})$

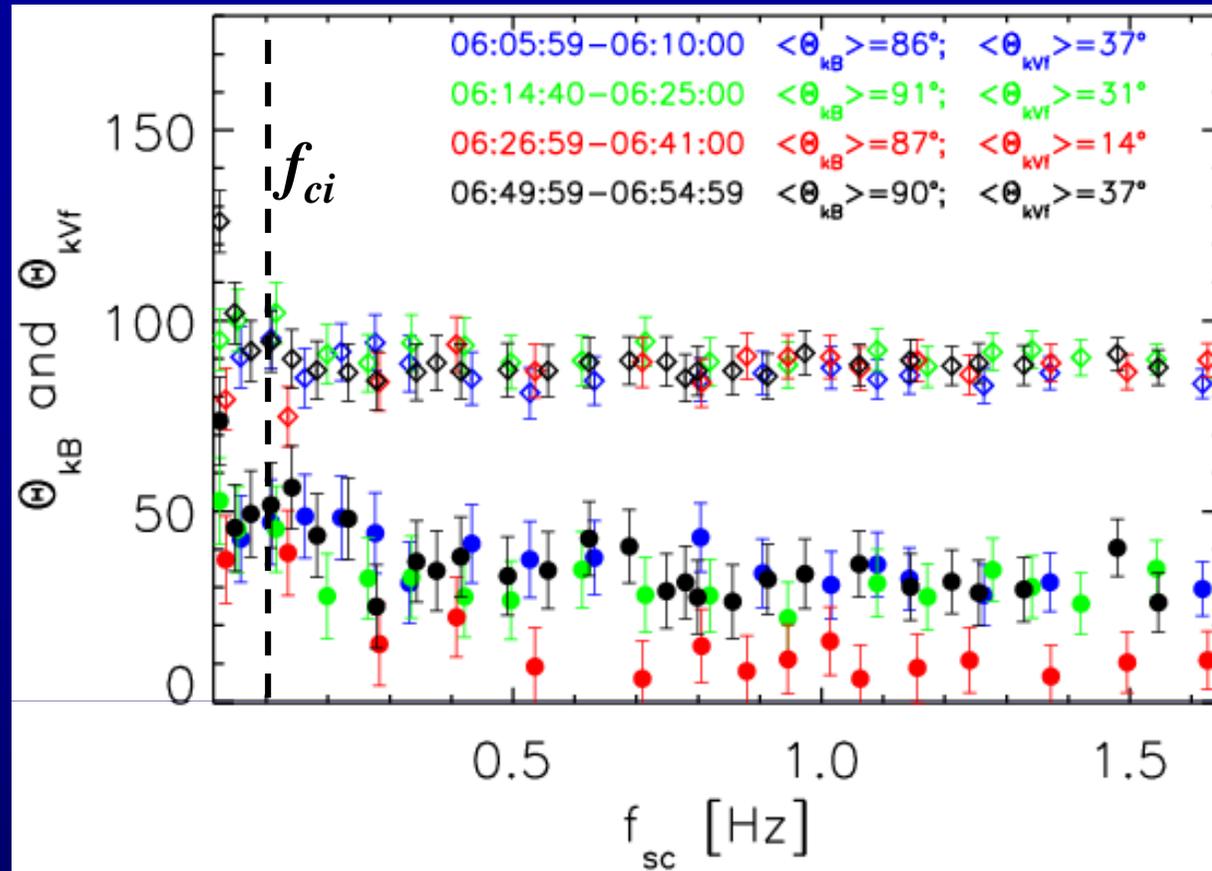


20040110 (d~200km)



We use $P(\omega, \mathbf{k})$ to calculate

1. 3D ω - k spectra
2. 3D k -spectra (anisotropies, scaling, ...)



Turbulence is

- $\perp B_0$ but non axisymmetric
- Quasi-stationary ($\omega_{plas} \sim 0$ although $\omega_{sat} \sim 20 \omega_{ci}$)

Comparison with the Vlasov theory

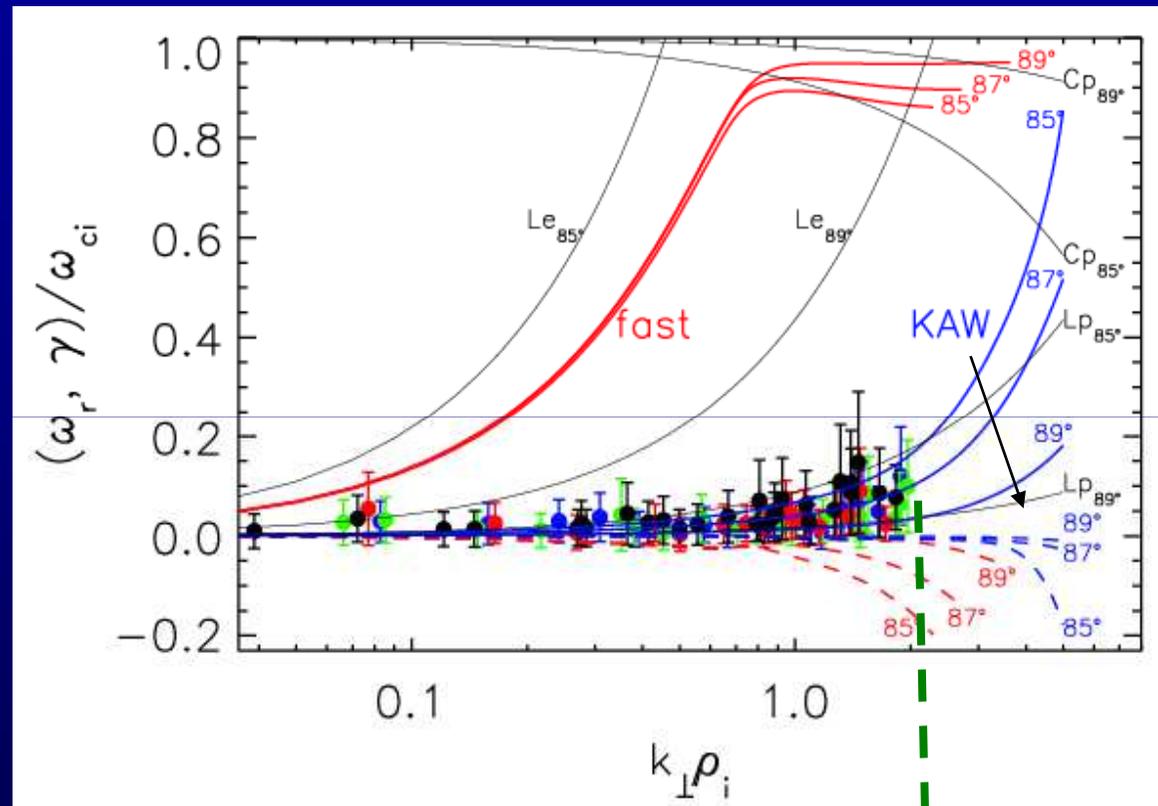
Turbulence cascades following the **Kinetic Alfvén mode (KAW)** as proposed in Sahraoui et al., PRL, 2009

→ **Rules out the cyclotron heating**

→ **Heating by p-Landau and e-Landau resonances**

[Sahraoui et al., PRL, 2010]

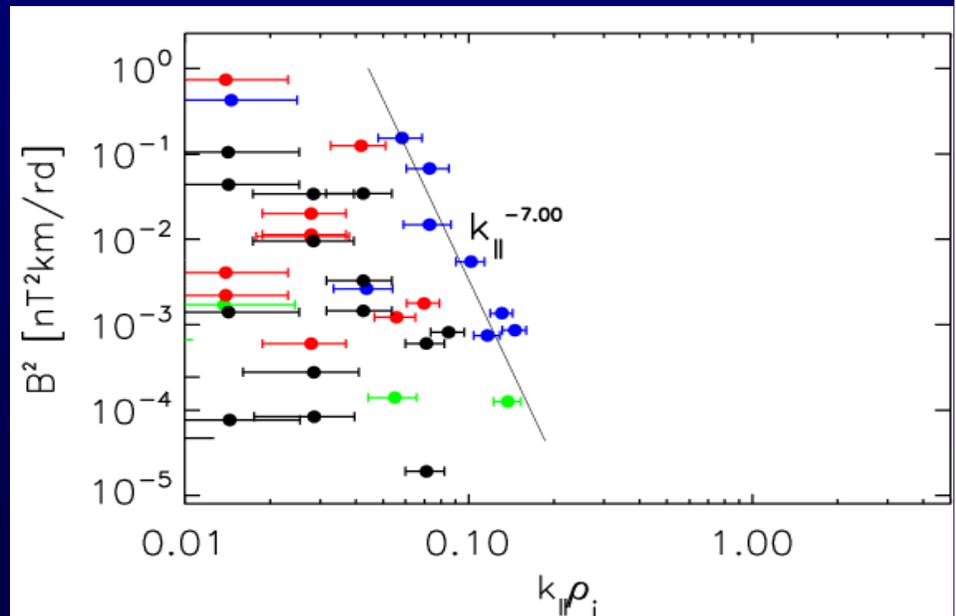
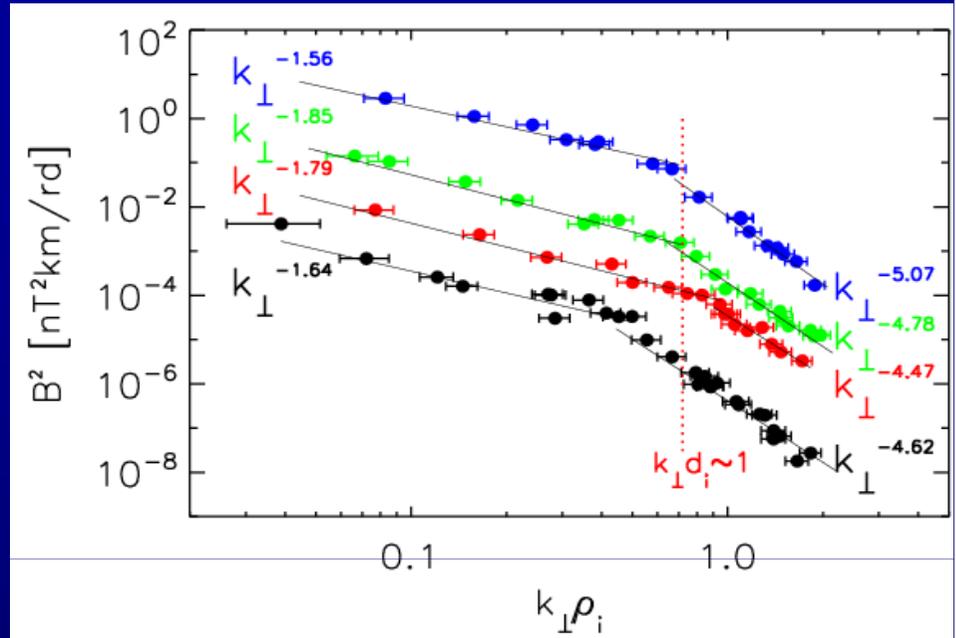
$$\beta_i \sim 2 \quad T_i/T_e = 3 \quad 85^\circ < \Theta_{kB} < 89^\circ$$



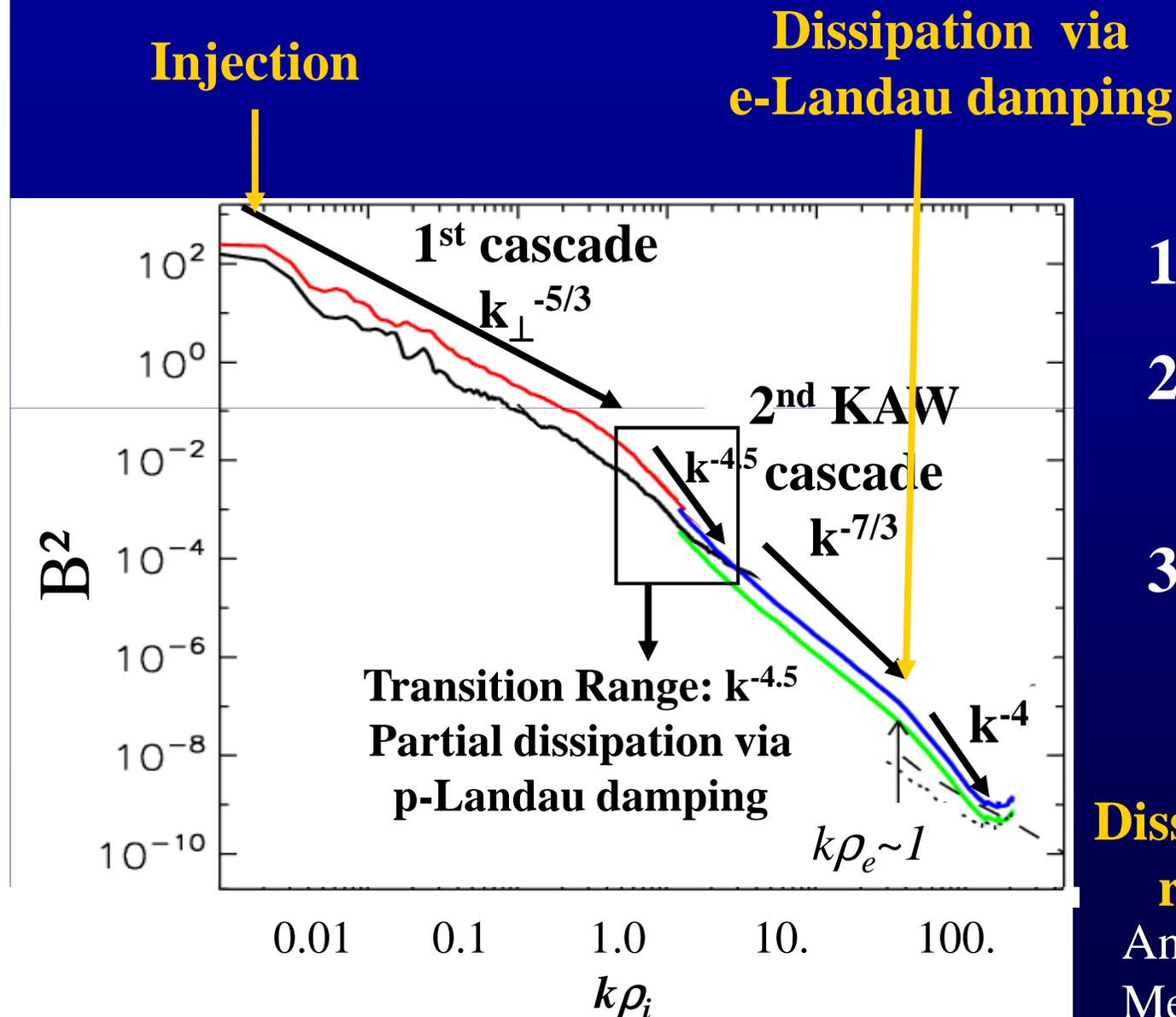
Limitation due to the
Cluster separation
($d \sim 200 \text{ km}$)

3D k -spectra at sub-ion scales

1. First *direct* evidence of the breakpoint near the proton gyroscale in k -space (*no additional assumption, e.g. Taylor hypothesis, is used*)
2. Strong steepening of the spectra below $\rho_i \rightarrow$ ***Transition Range*** to dispersive/electron cascade



Journey of the energy cascade through scales



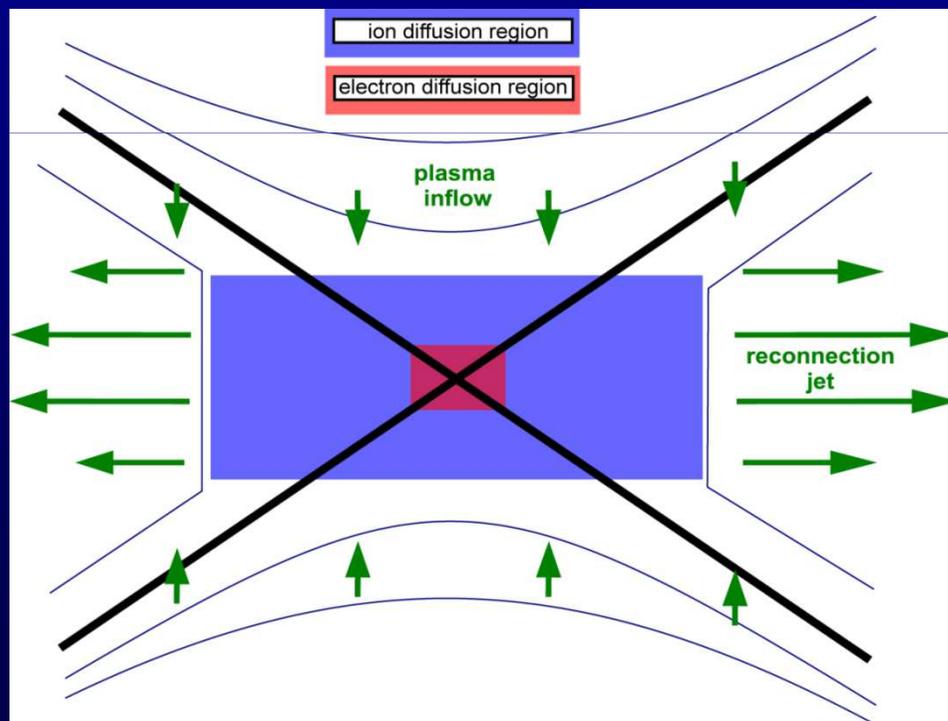
1. Turbulence
2. e-Acceleration & Heating
3. Reconnection

Dissipation range

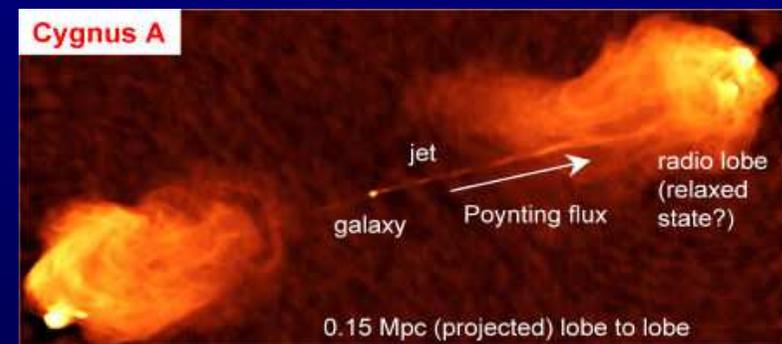
Another interpretation in Meyrand & Galtier, 2010

Dissipation through reconnection/current sheets

Large scale laminar current sheet: reconnection can occur and the can be heated or accelerated (e..g. jets)

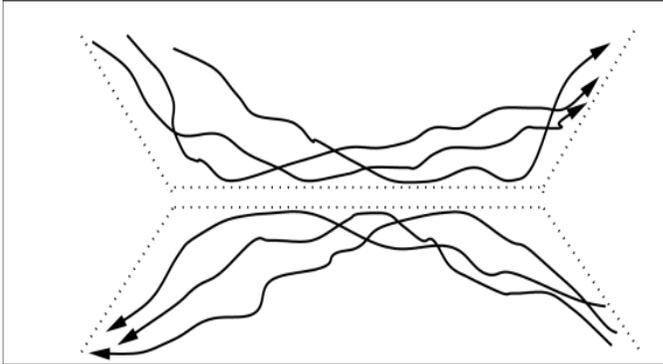


[Zhong+, Nature Physics, 2010]

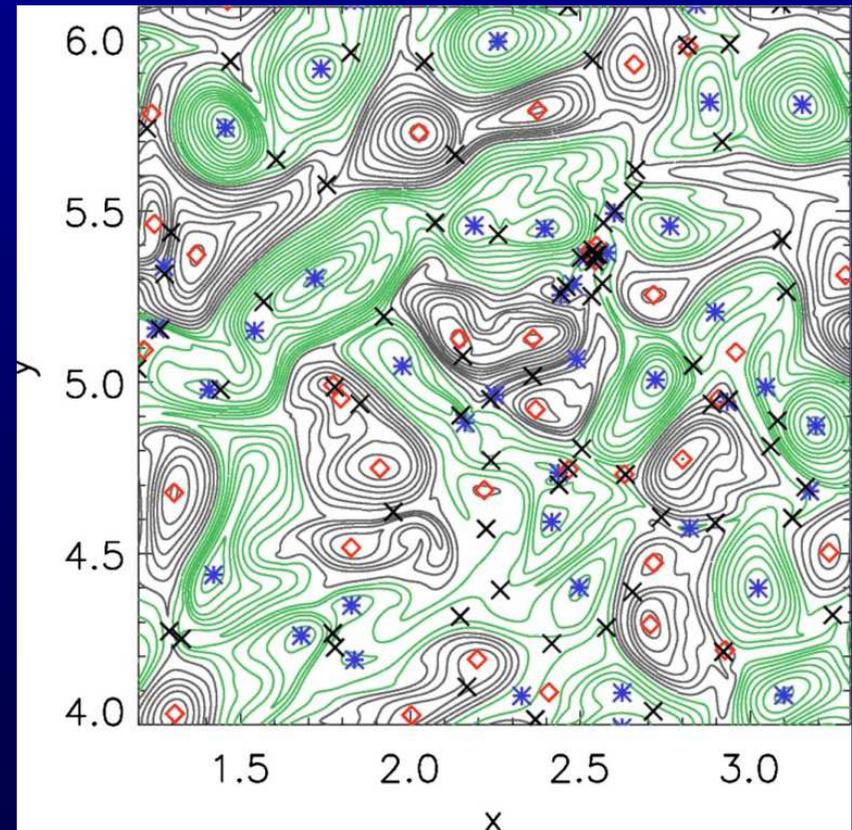
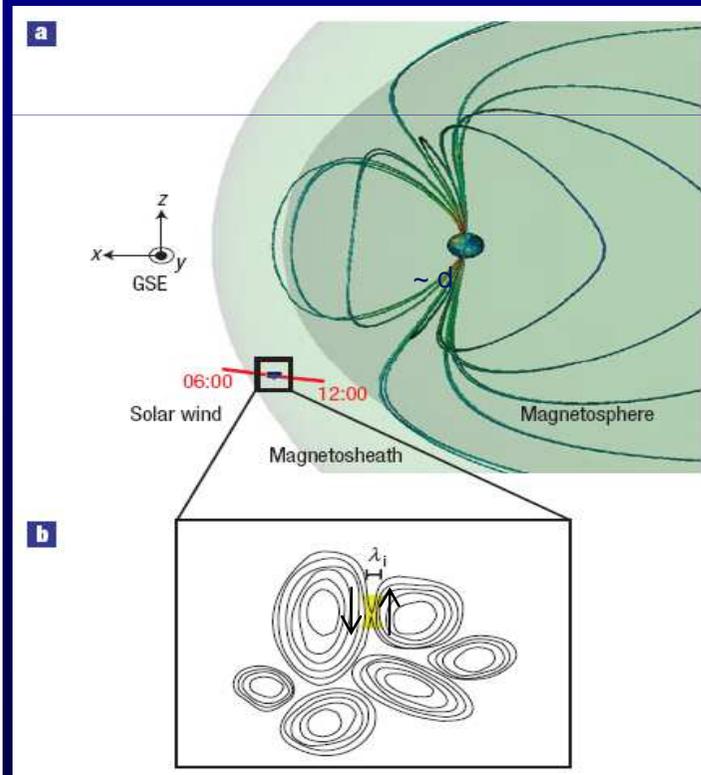


Turbulent current sheets

[Lazarian & Vishniac, 1999]



2D Hall-MHD simulation of turbulence: evidence of a large number of reconnecting regions

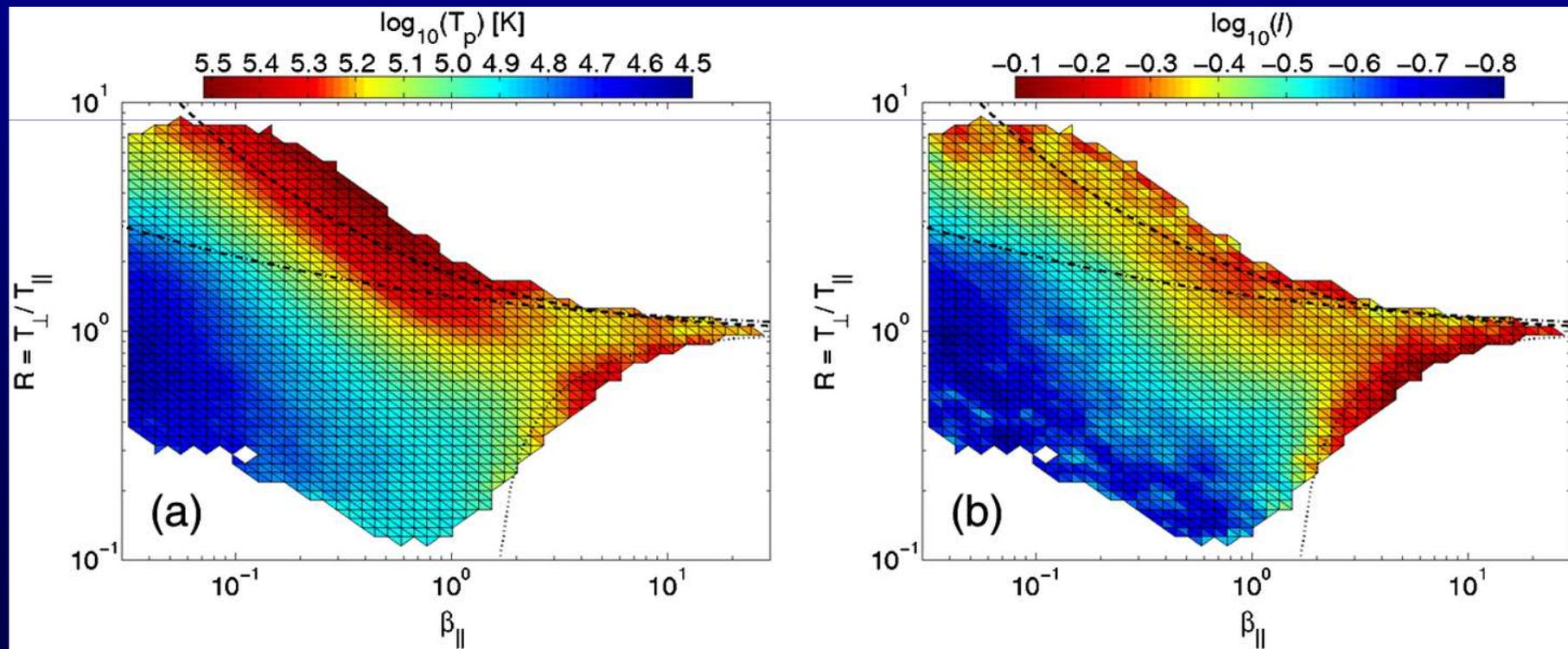


[e.g., Retinò+, Nature Physics, 2007]

Dissipation by wave-particle interaction or via reconnection?

Good correlation between enhanced T_p and threshold of linear kinetic instabilities

Good correlation between enhanced high shear B angles and the threshold of linear instabilities !!



Statistical approach to small scale SW turbulence

Which statistical description applies to sub-proton
scale SW turbulence:

1. *Weak or strong turbulence?*
2. If strong, then is it *self-similar/monofractal or
intermittent/multifractal?*

1. Strong vs Weak Turbulence:

Often it has been argued that small scale/high frequency turbulence in the solar wind is a weak turbulence because $|\delta\mathbf{B}|/B \ll 1$

This is wrong !

Because *only* the ratio nonlinear/linear times (or terms) for *each physical* system can indicate how weak or strong is the turbulence

Let us consider the example of Incompressible MHD

Incompressible Alfvénic Turbulence

$$\partial_t z^\pm \mp v_A \cdot \nabla z^\pm + z^\mp \cdot \nabla z^\pm = -\nabla p$$

Linear term: $k_{\parallel} v_A z^{\pm}$

Nonlinear term: $k_{\perp} u_{\perp} z^{\pm}$

Ratio of nonlinear to linear terms:

$$\chi = \frac{k_{\perp} u_{\perp}}{k_{\parallel} v_A}$$

$\chi \ll 1 \Rightarrow$ Weak turbulence with $k_{\parallel} v_A \gg k_{\perp} u_{\perp}$

$\chi \sim 1 \Rightarrow$ Strong turbulence with $k_{\parallel} v_A \sim k_{\perp} u_{\perp}$ (or $\omega \sim \omega_{\text{NL}} \Rightarrow$
Critical balance conjecture)

For anisotropy $k_{\perp} \gg k_{\parallel}$ we have **STRONG** turbulence

($\chi \sim 1$) even when $\frac{u_{\perp}}{v_A} \sim \frac{\delta B}{B_0} \ll 1$

⇒ **One has to give up using mere criteria, e.g. $|\delta B|/B \ll 1$, to discriminate within the data between weak/strong turbulence theories**

Other alternatives?

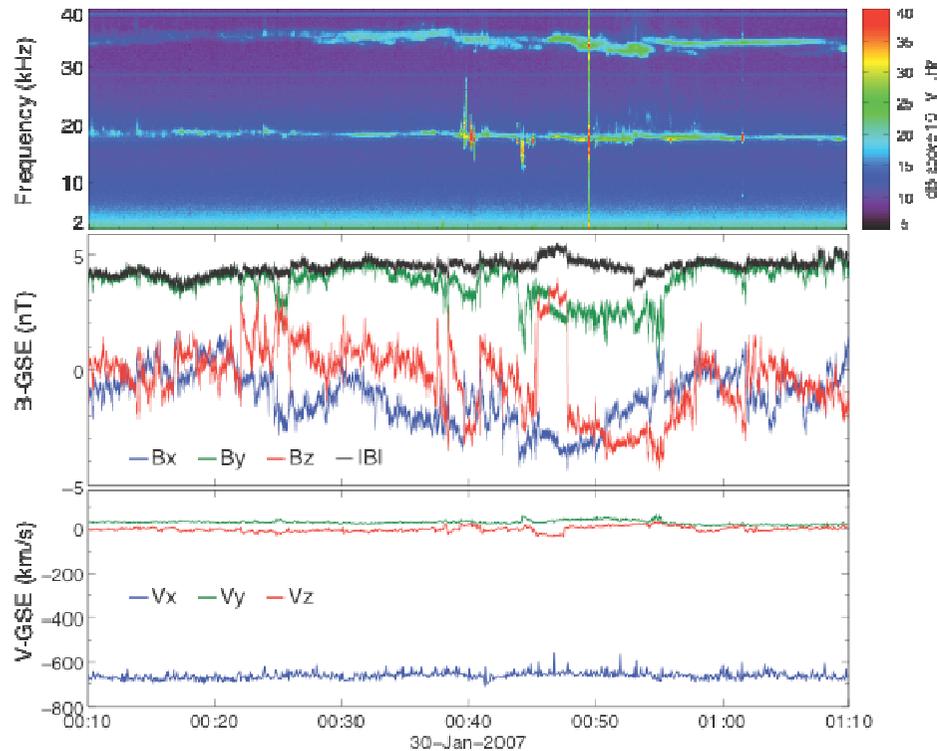
1. Estimation of the linear/nonlinear times of the turbulence from the data
→ But it is difficult because this generally requires **to know accurately** the nature of the turbulence and its spatial scales (\parallel and \perp)
2. Estimating phase coherence **directly** from the measured Fourier phases of the turbulence from the data using, e.g., **Surrogate data** [Hada et al., 2003; Sahraoui, PRE, 2008; Sahraoui & Fauvarque, in prep.]

2. Monfractality vs multifractality in the dispersive range:

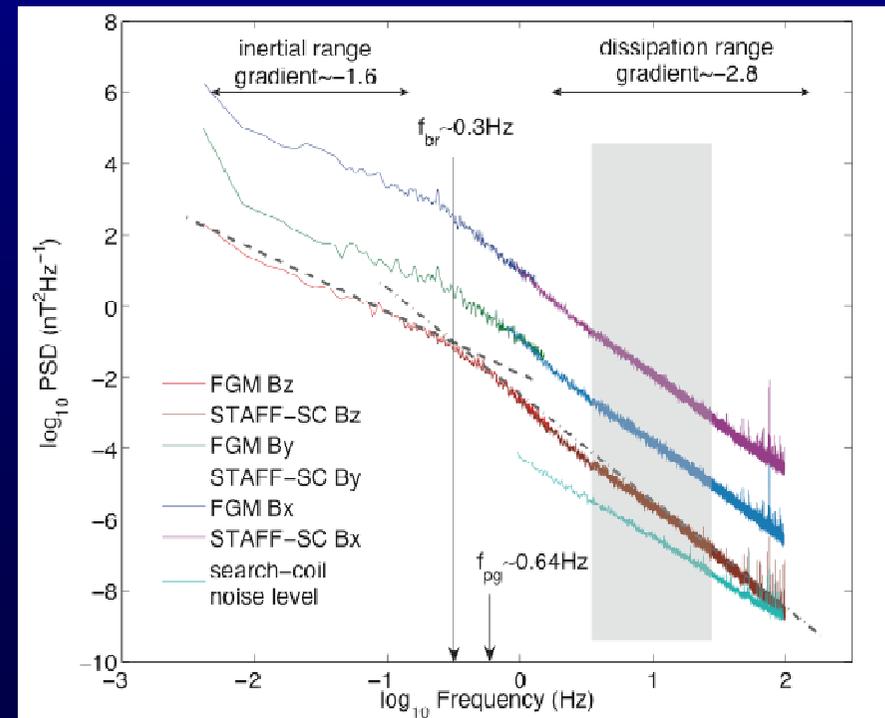
$$n_e \sim 4 \text{ cm}^{-3} \quad \text{ion } \beta \sim 2$$

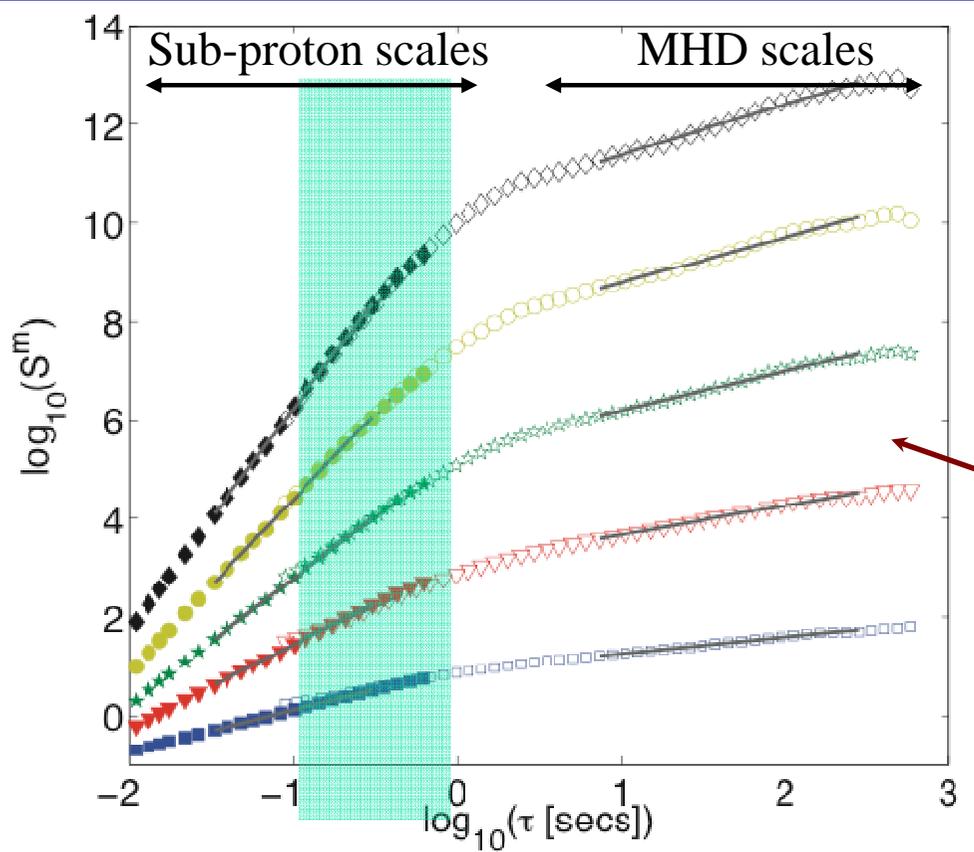
$$V_A \sim 50 \text{ km s}^{-1}$$

$$T_i \sim 103 \text{ eV} \quad |B| \sim 4 \text{ nT}$$



[Kiyani et al., PRL, 2009]





Structures functions:

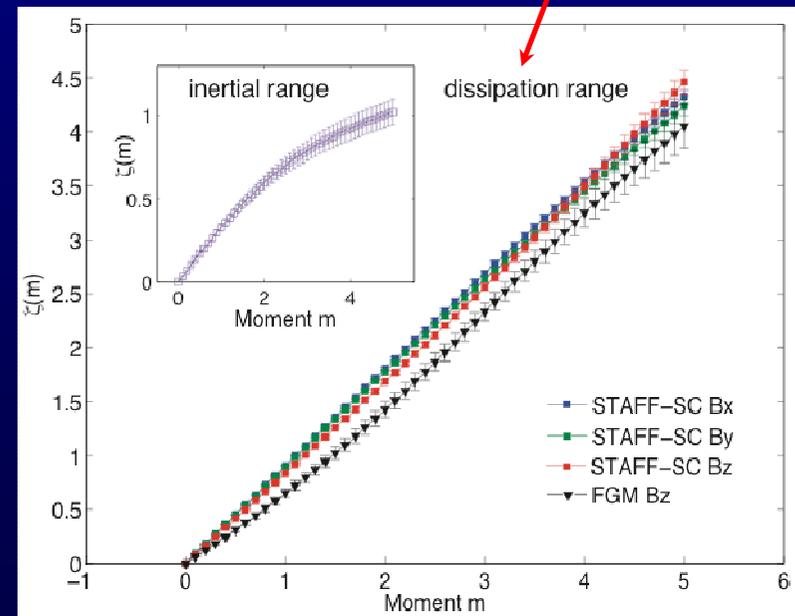
$$S^m(\tau) = \sum_t |B(t+\tau) - B(t)|^m$$

Scaling:

$$S^m(\tau) = S^m(1) \tau^{\zeta(m)}$$

Evidence of **monofractality (self-similarity)** at sub-proton scales, while MHD-scales are **multifractal (intermittent)**

[See also Alexandrova et al., ApJ, 2008]



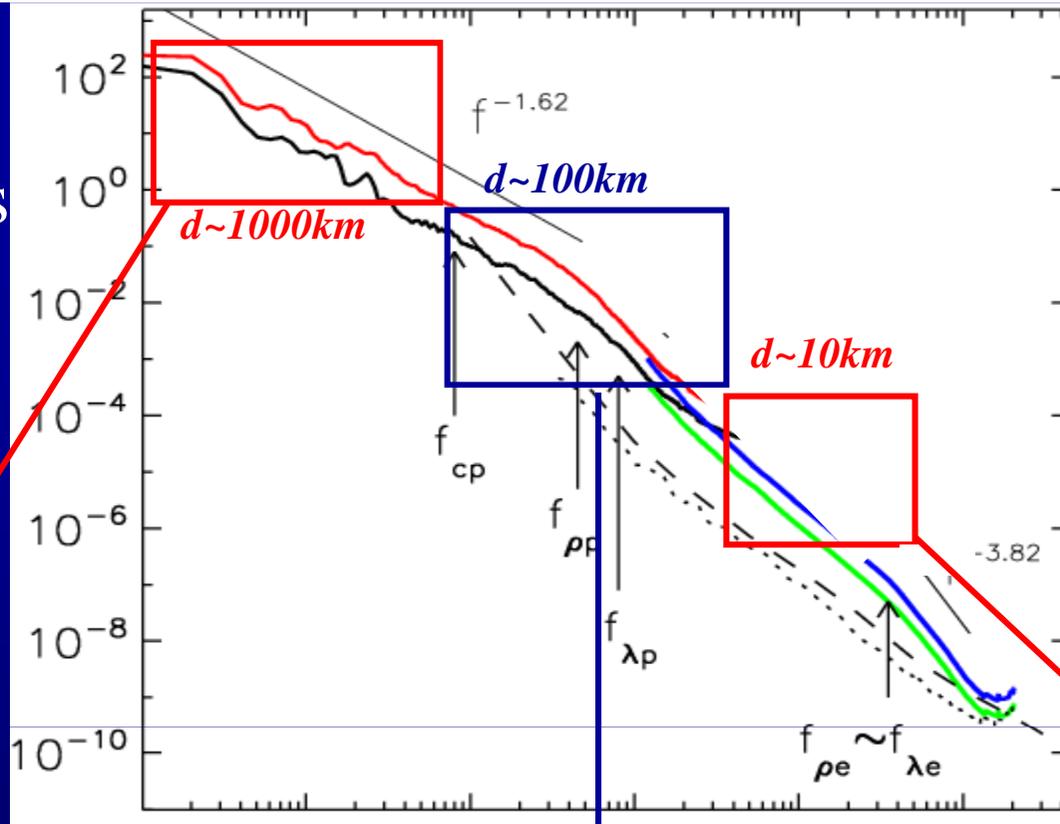
Conclusions

The Cluster data helps understanding crucial problems of astrophysical turbulence:

- Its nature and anisotropies in *k-space* at MHD and sub-ion scales
- Its cascade and dissipation down to the electron gyroscale $\rho_e \Rightarrow$ *electron heating and/or acceleration by turbulence*
- Strong evidences of KAW turbulence ($\omega \ll \omega_{ci}$, $k_{//} \ll k_{\perp}$) \Rightarrow *Heating by e-p-Landau dampings (no cyclotron heating)*
- Importance of kinetic physics in SW turbulence
- Turbulence & dissipation are at the heart of the future space missions: NASA/MMS (2014), ESA/SO (2017), NASA/SPP (2019), TOR (????)

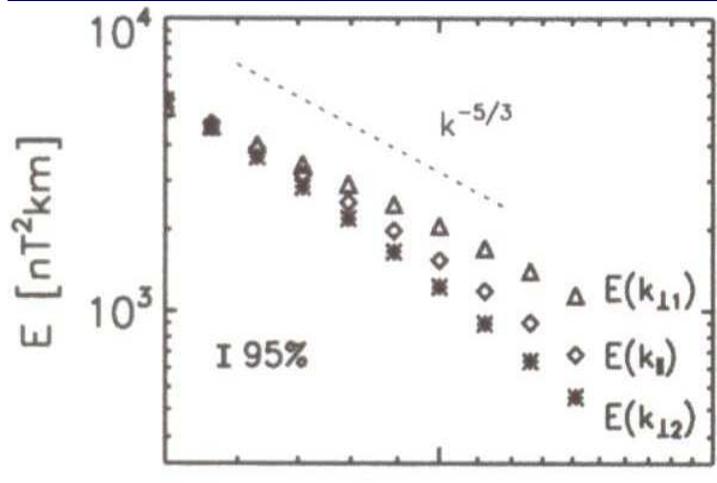
⇒ Need of **multi-scale** measurements *with appropriate spacecraft separations*

Narita *et al.* PRL, 2010

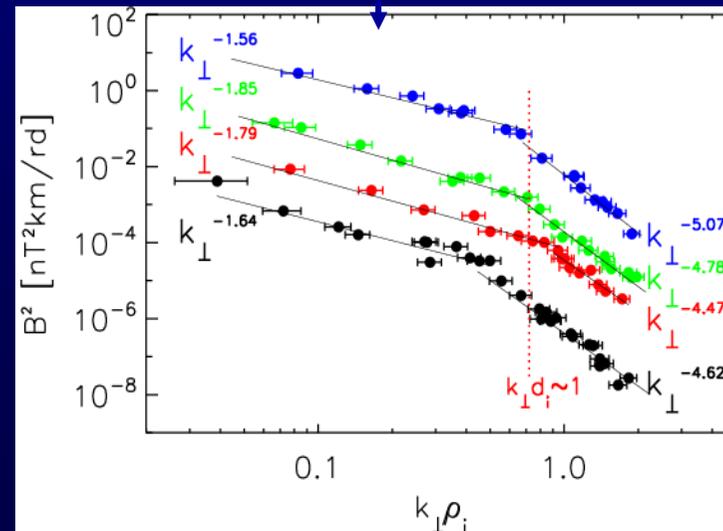


MMS

2014



Sahraoui *et al.* PRL, 2010



La turbulence et les futures missions spatiales

MAGNETOSPHERIC MULTISCALE

A SOLAR-TERRESTRIAL PROBE

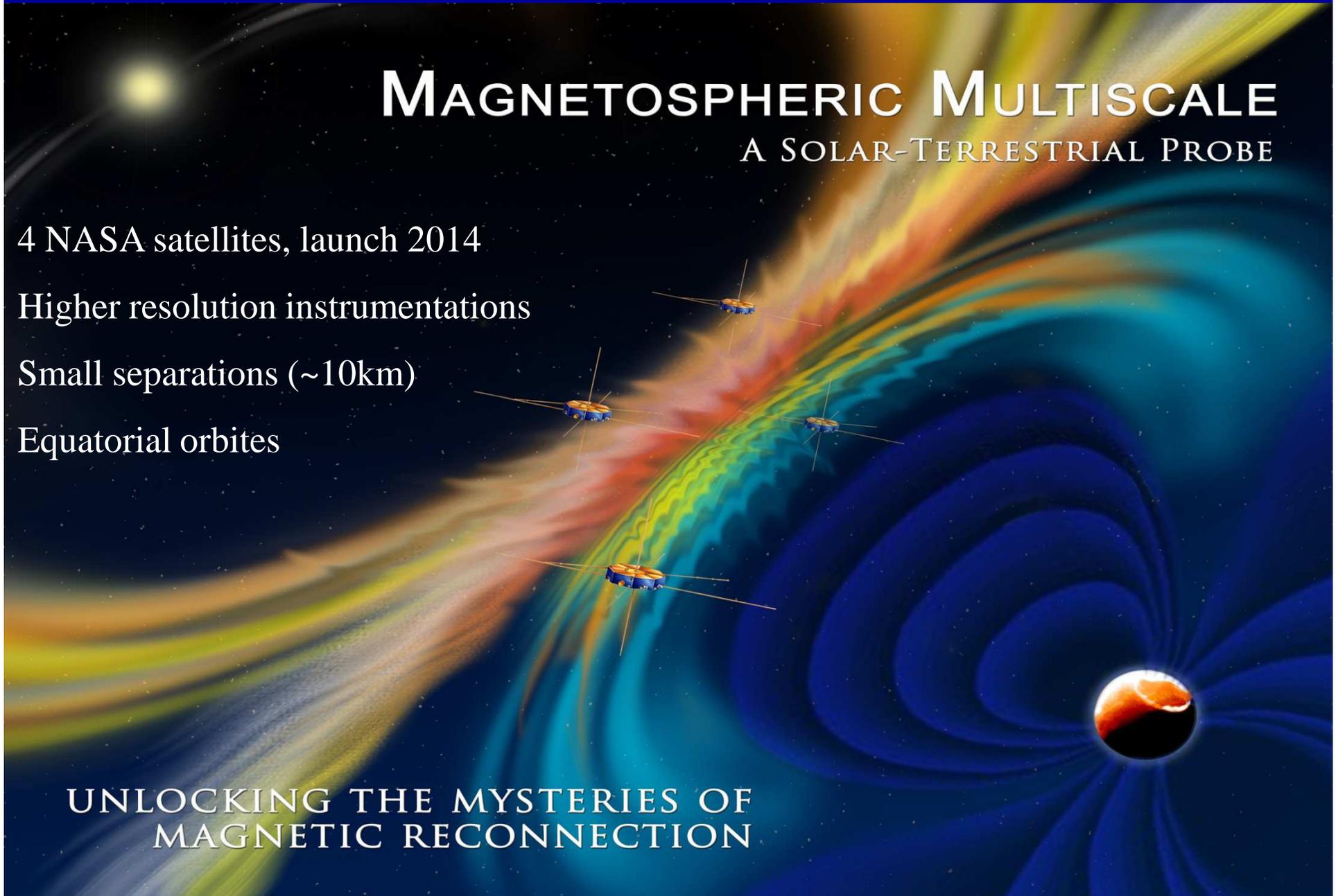
4 NASA satellites, launch 2014

Higher resolution instrumentations

Small separations (~10km)

Equatorial orbits

UNLOCKING THE MYSTERIES OF
MAGNETIC RECONNECTION





Solar Orbiter

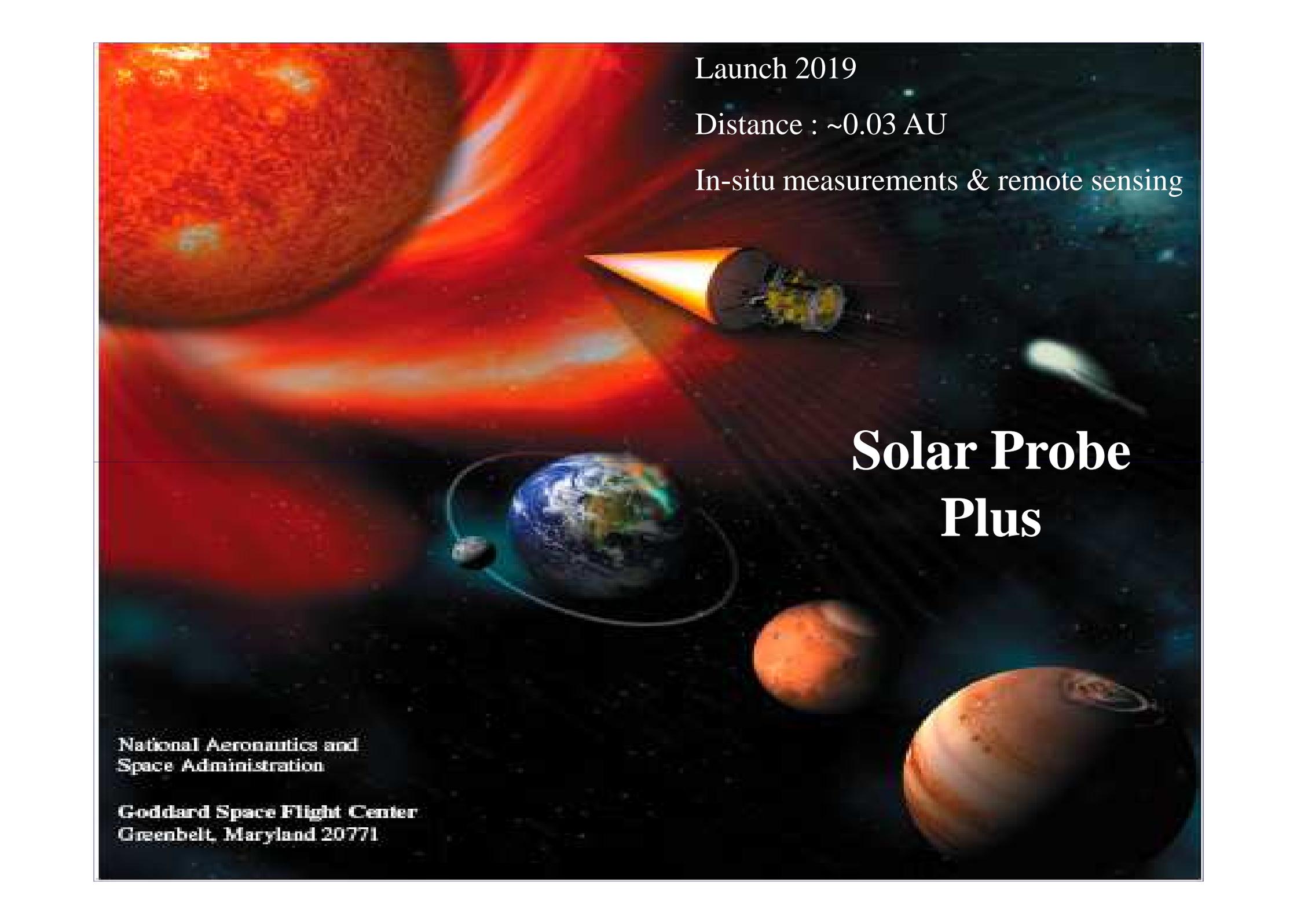
Exploring the Sun-Heliosphere Connection

Launch 2017

Distance : 0.28 AU

In-situ measurements & remote sensing



A detailed illustration of the Solar Probe Plus mission. The Sun is a large, glowing orange-red sphere in the upper left. A spacecraft with a large heat shield is shown in a highly elliptical orbit around the Sun. The Earth and Moon are shown in the lower left, with the Earth's blue and white atmosphere and the Moon's grey surface. Mars is shown in the lower right. The background is a dark space filled with stars and a nebula. Text is overlaid on the top right, and the mission name is in the center right. Agency information is in the bottom left.

Launch 2019

Distance : ~0.03 AU

In-situ measurements & remote sensing

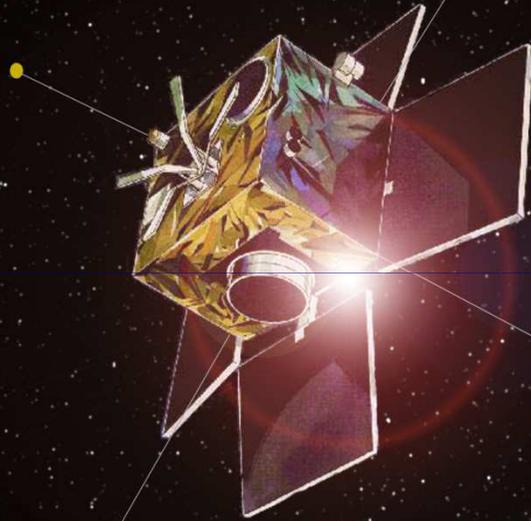
Solar Probe Plus

National Aeronautics and
Space Administration

Goddard Space Flight Center
Greenbelt, Maryland 20771

TOR

Solving energy dissipation problem at kinetic scales in the solar wind



Mission proposed to ESA (2012)
Currently under review with SNBS
Recently proposed to CNES