

Plasma Dynamo Experiments



Cary Forest

Acknowledgements

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The Dynamo regime requires steady:

- Highly conducting and fast flowing

$$Rm = \mu_0 \sigma V L \gg 1 \quad \text{Magnetic Reynolds}$$

- Turbulent or laminar

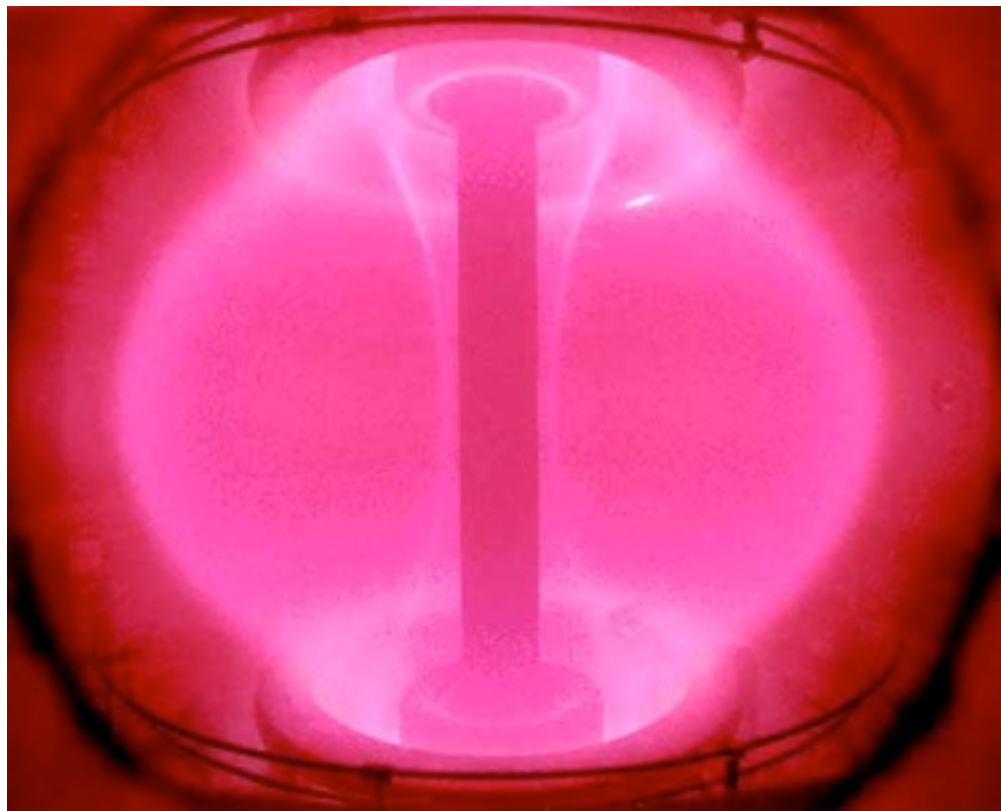
$$Re = \frac{VL}{\nu} \gg 1 \quad \text{Kinetic Reynolds}$$

- Kinetic energy dominated

$$\frac{1}{2} \rho V^2 \gg \frac{B^2}{2\mu_0} \quad \text{or} \quad V/V_A > 1 \quad \text{Alfvén Mach}$$

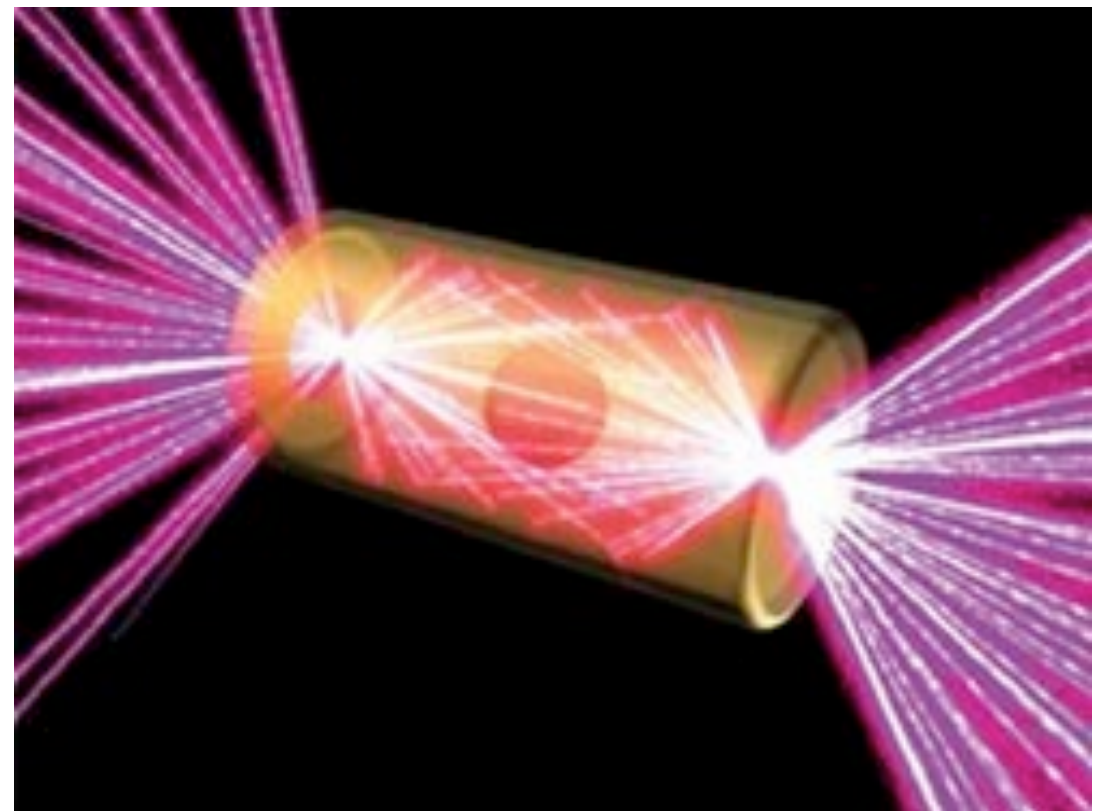
Most plasma experiments are not suitable for dynamo studies

Magnetically Confined Plasmas



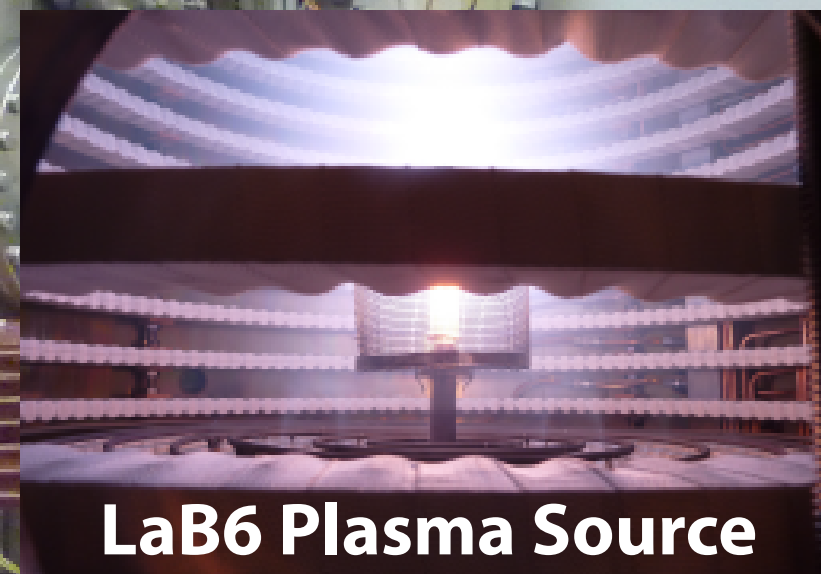
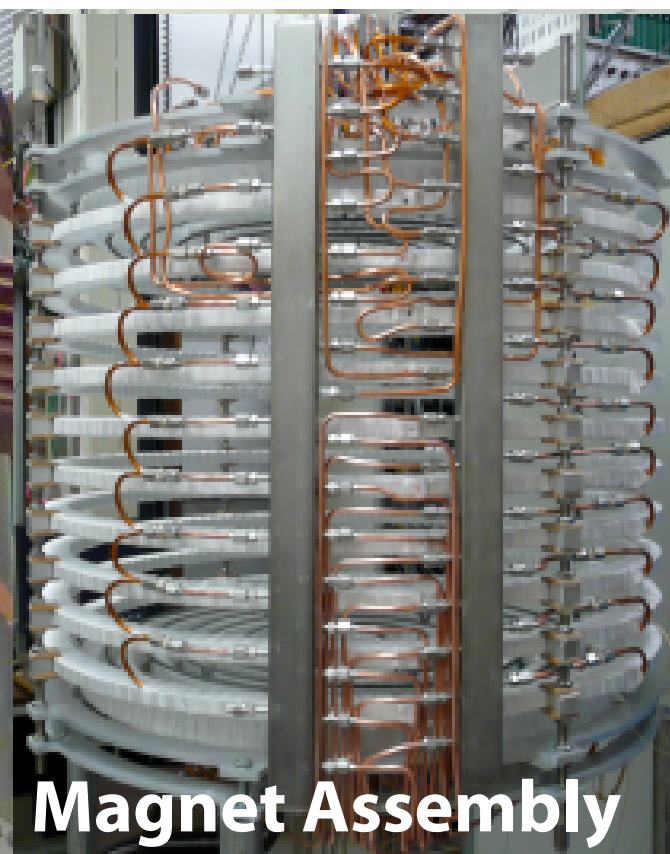
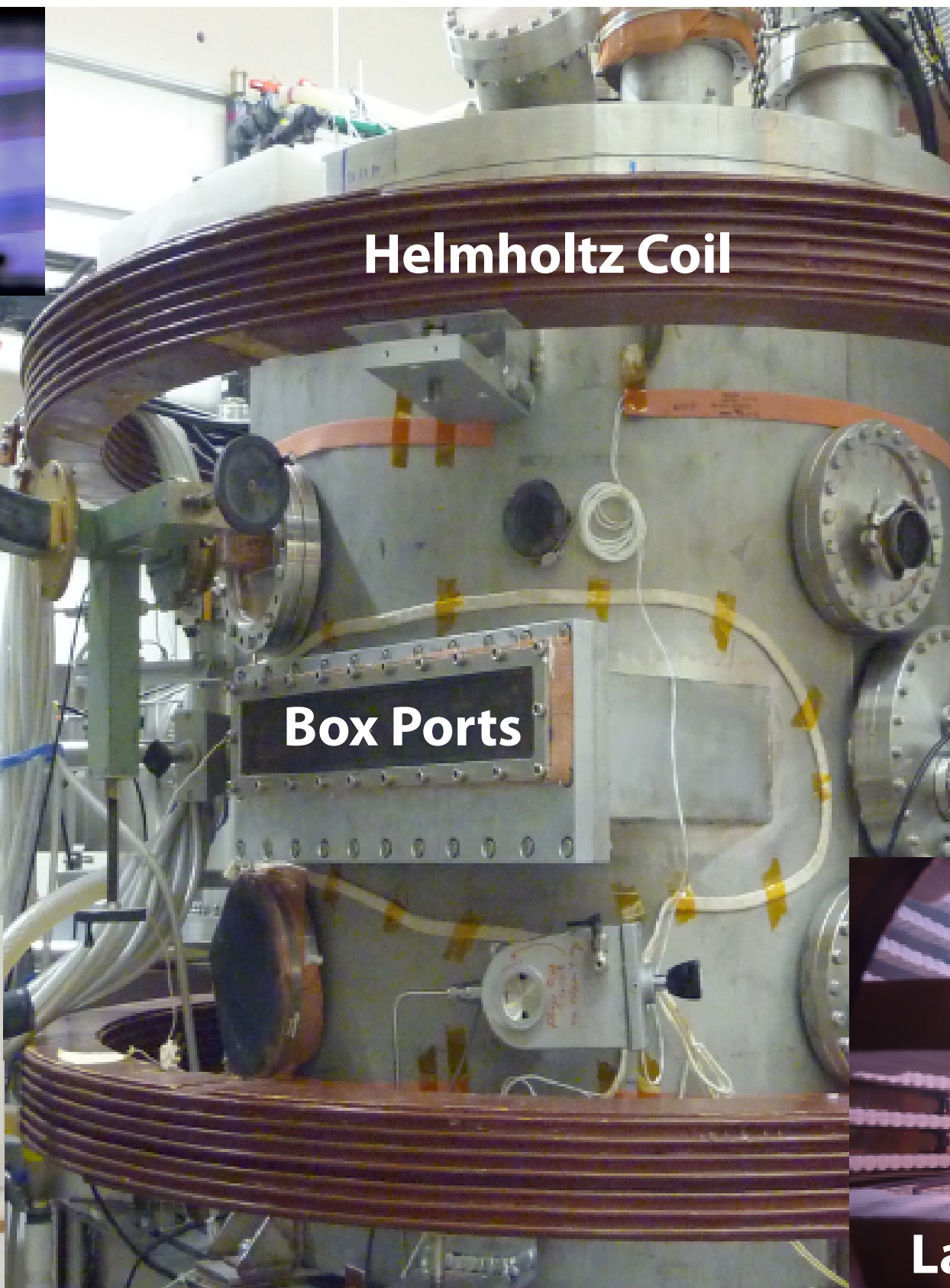
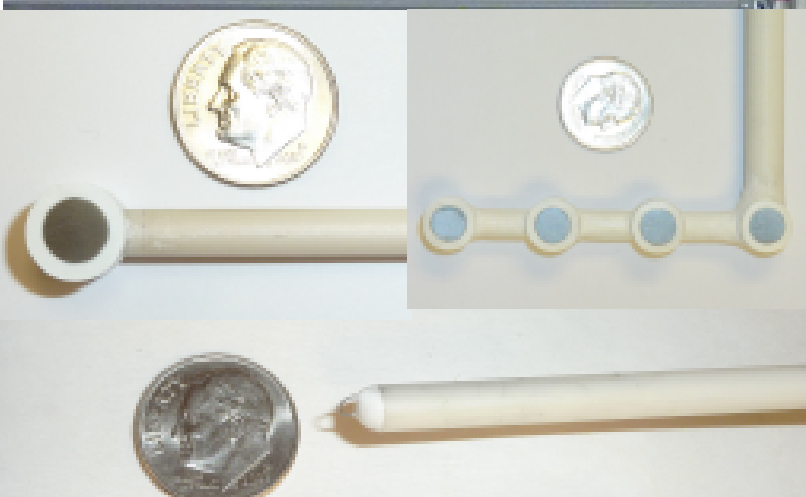
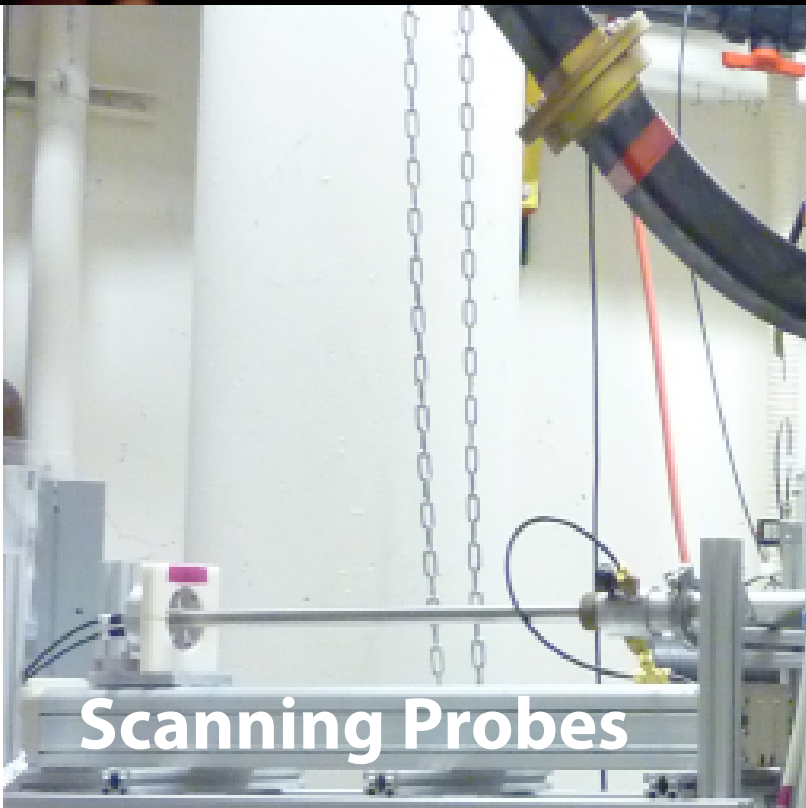
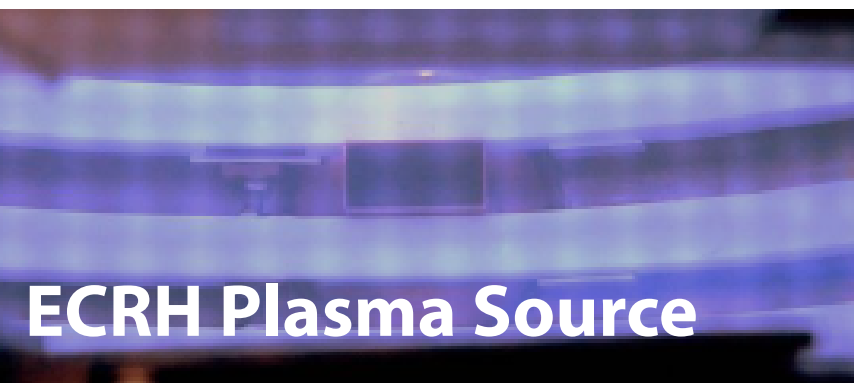
- $V = 100 \text{ km/s}$, $T=1 \text{ keV}$, $B=1 \text{ T}$
- high Rm
- strongly magnetized
- pressure and magnetically driven MHD
- $\tau \ll \mu\sigma L^2$

Inertially Confined Plasmas

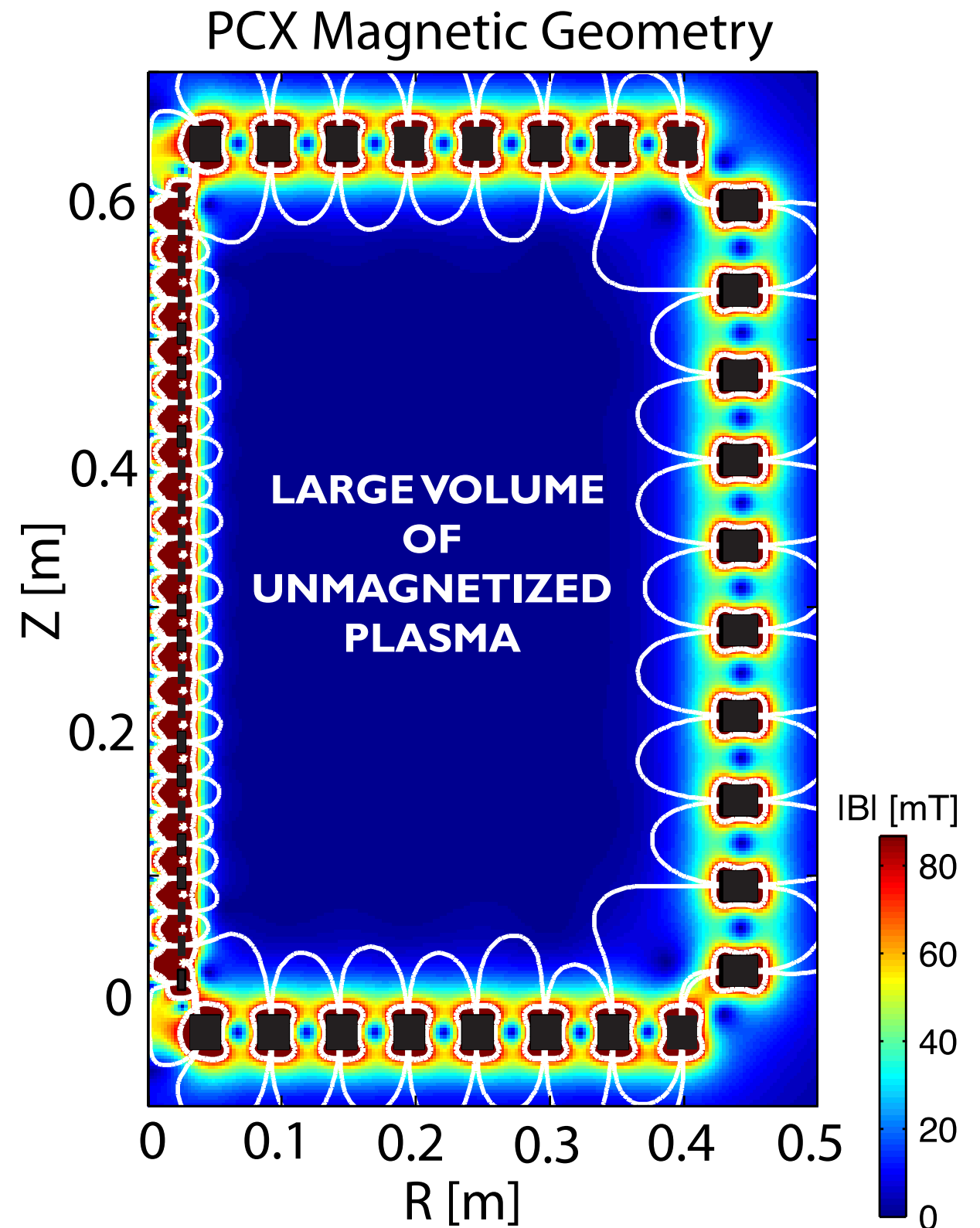


- $V \sim 50 \text{ km/s}$, $T_e \sim 100 \text{ eV}$
- modest Rm
- weakly magnetized
- $\tau \sim a/V \ll \mu_0\sigma L^2$

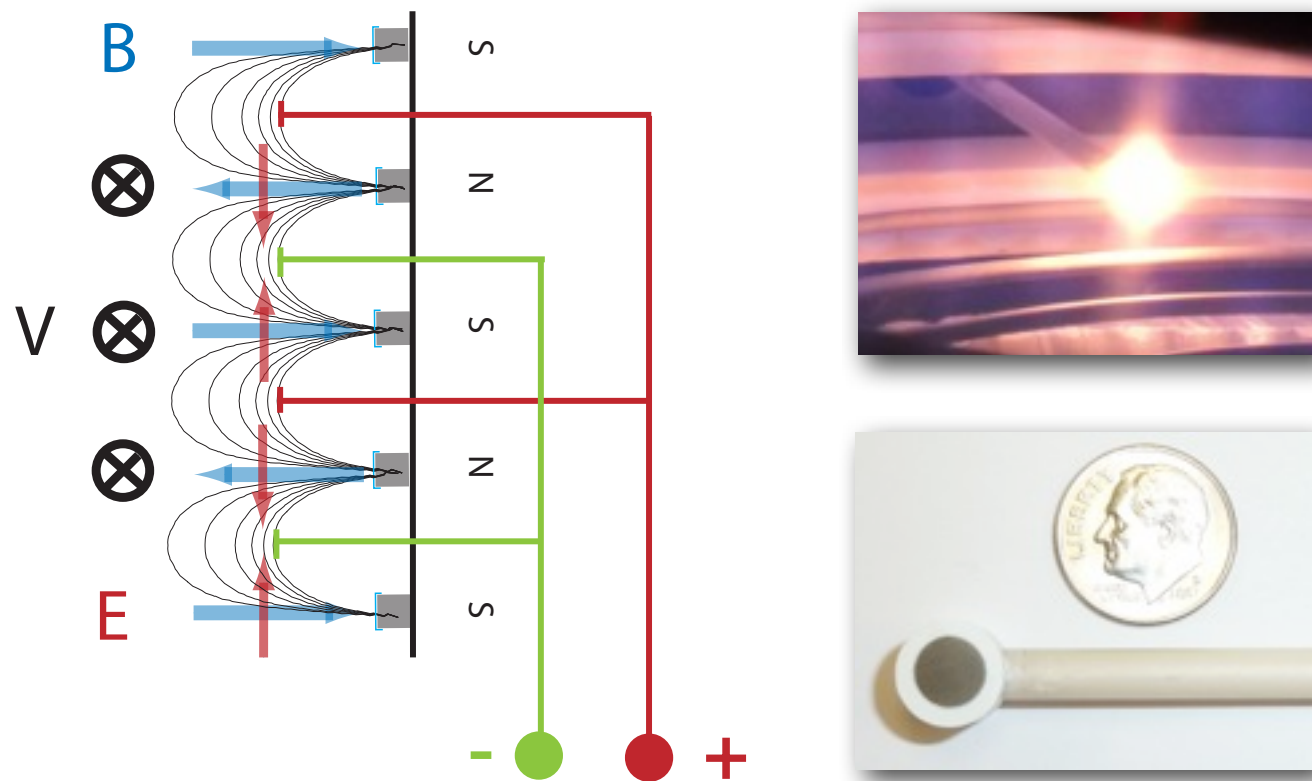
Review: Stirring Unmagnetized Plasma



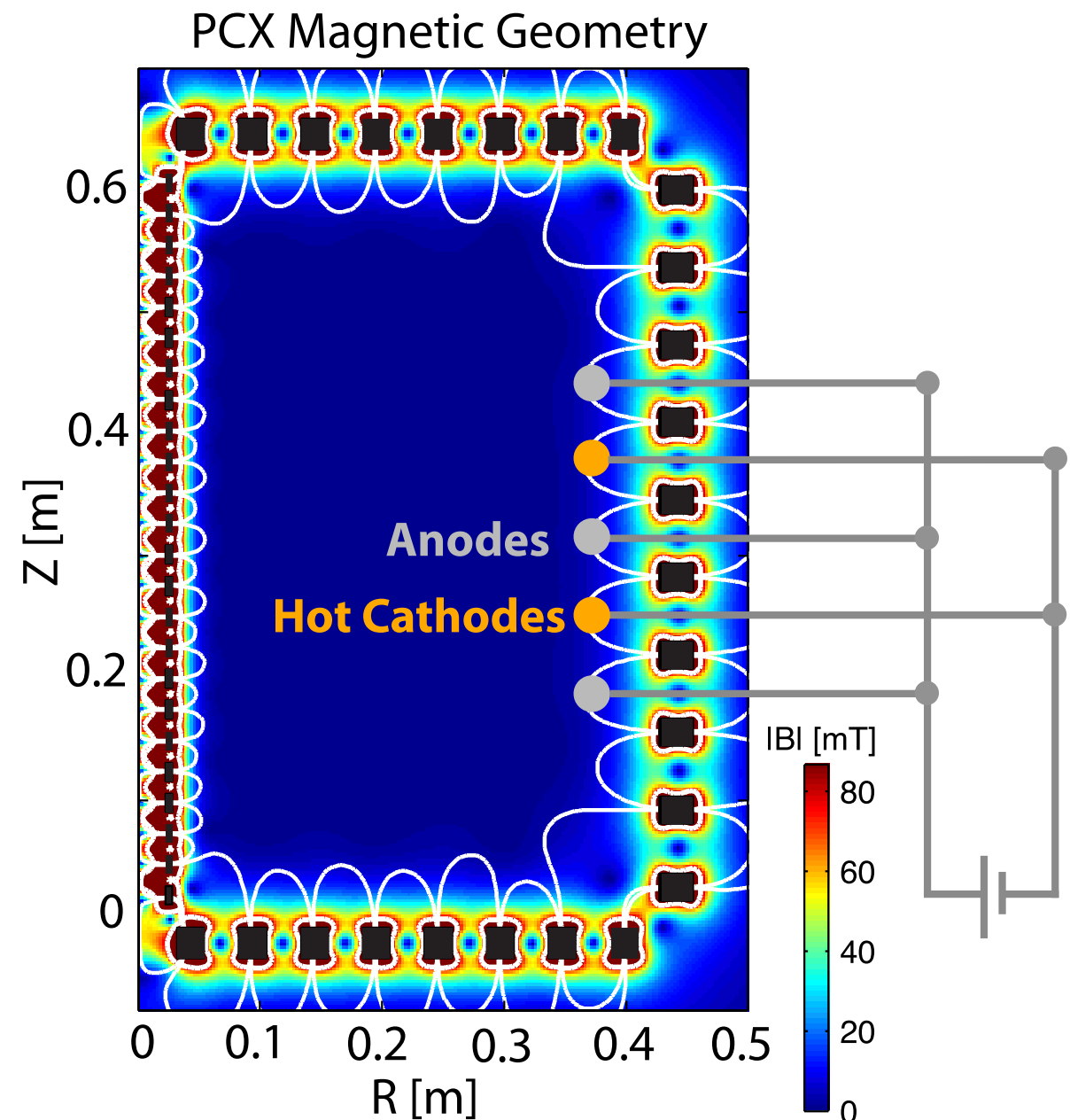
Axisymmetric Plasma Bucket Provides confinement



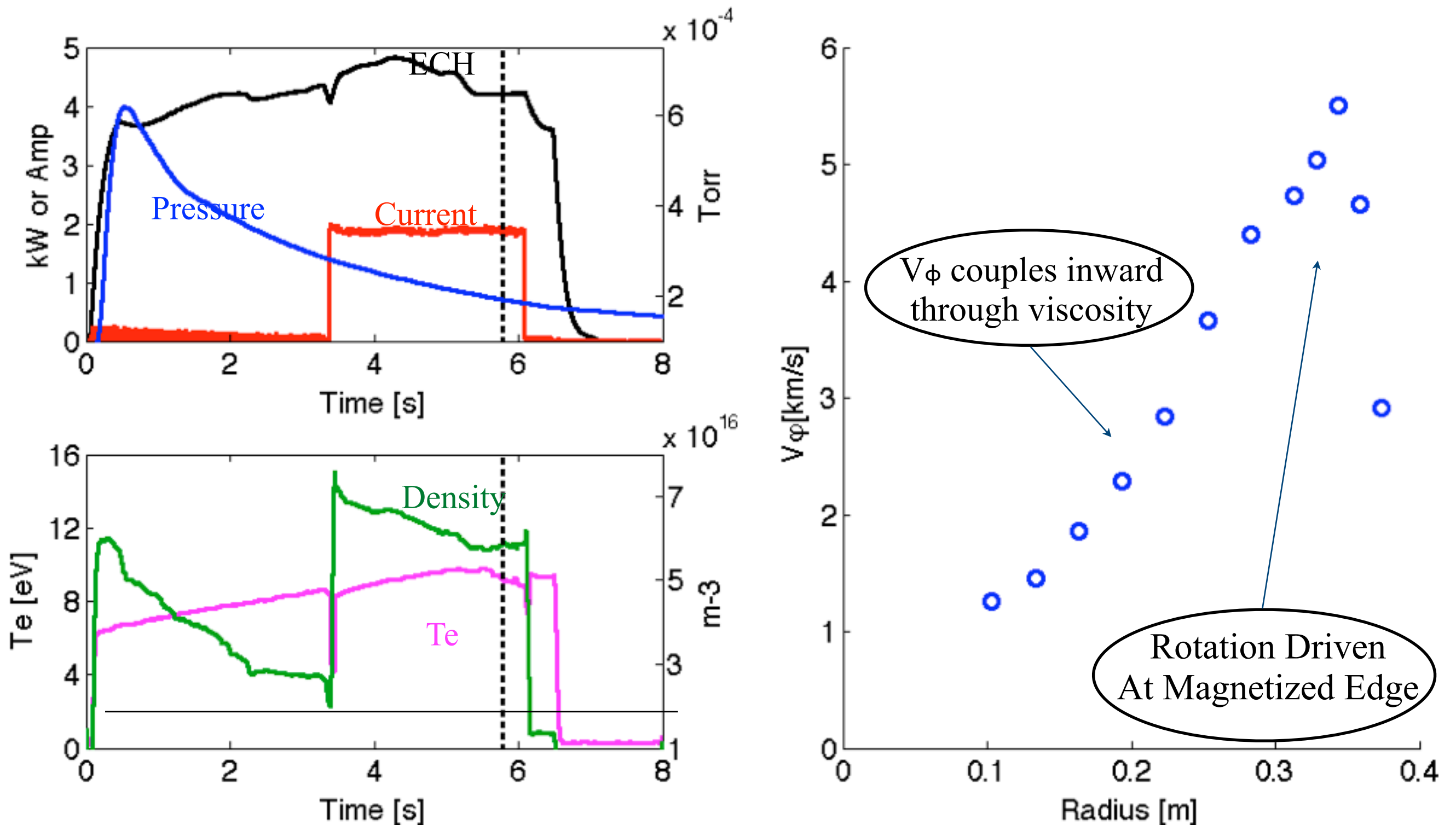
Plasma Stirring from boundary using electrodes



- Toroidally localized cathodes and anodes are differentially biased to create $J \times B$ torque
- Velocity couples inward to the unmagnetized region through viscosity
- Rotation (measured with Mach probe) is axisymmetric



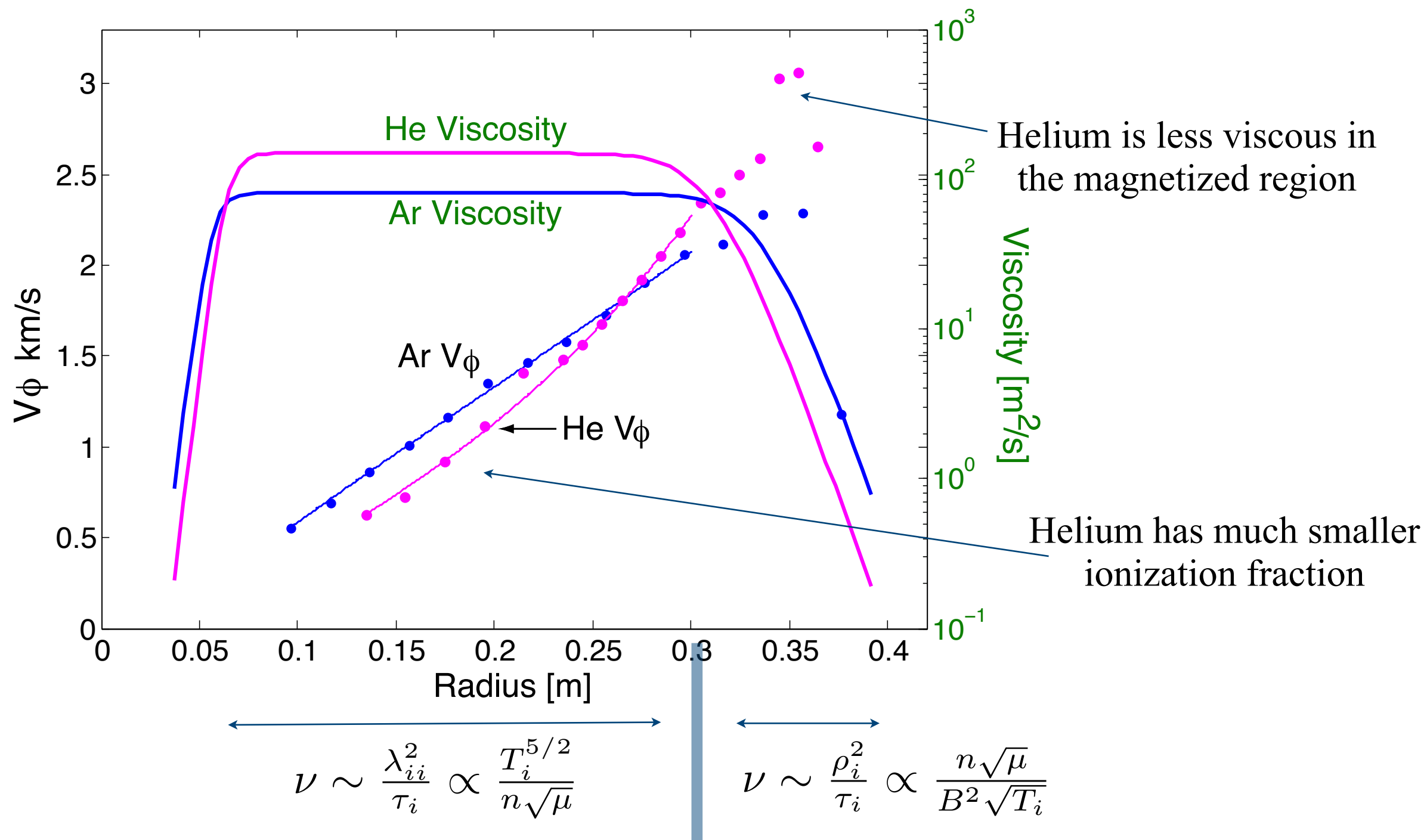
High speed flows in unmagnetized core!



$Pm \sim 10$ $Rm \sim 50$

Plasma Viscosity Depends on Magnetization

Case	$n \times 10^{11}$ (cm^{-3})	f_{ion} %	Te (eV)	Ti _{fit} (eV)	V _{max} (km/s)	Rm	Re	Pm
Ar400VT2	.77	27	7	0.6	2.2	11	8	1.3
He400VT1	.45	0.6	11	0.4	2.4	22	5	4.5

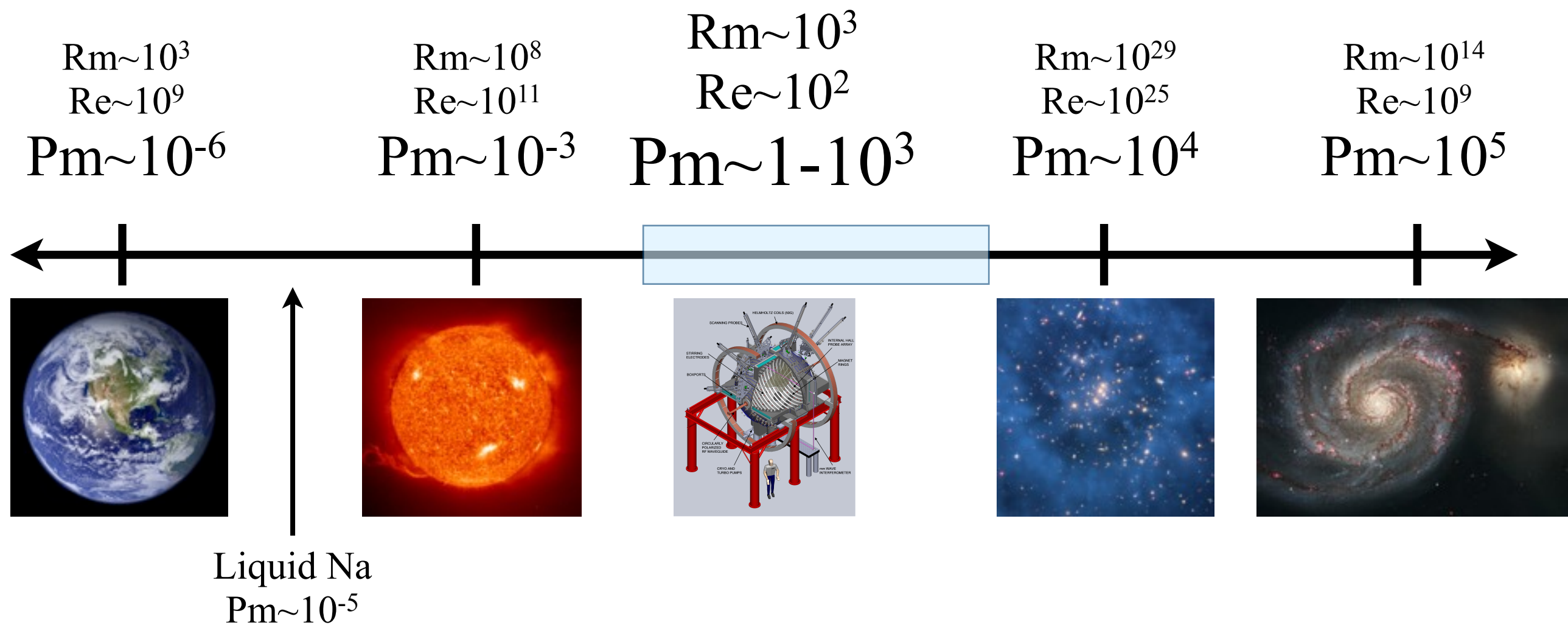


Next Step: Plasma Dynamo Experiments

- $R_m > 1000$ $R_m = 50 T_{e,10\text{eV}}^{3/2} V_{\text{km/s}} L_m$
- Independent $Re = 8 \mu^{1/2} n_{10^{18}\text{m}^{-3}} V_{\text{km/s}} L_m / T_{i,\text{eV}}^{5/2}$
- Rapidly Rotating
- Compressibility, stratification, buoyancy
- Plasma Effects beyond MHD: neutrals, kinetic effects, Hall MHD

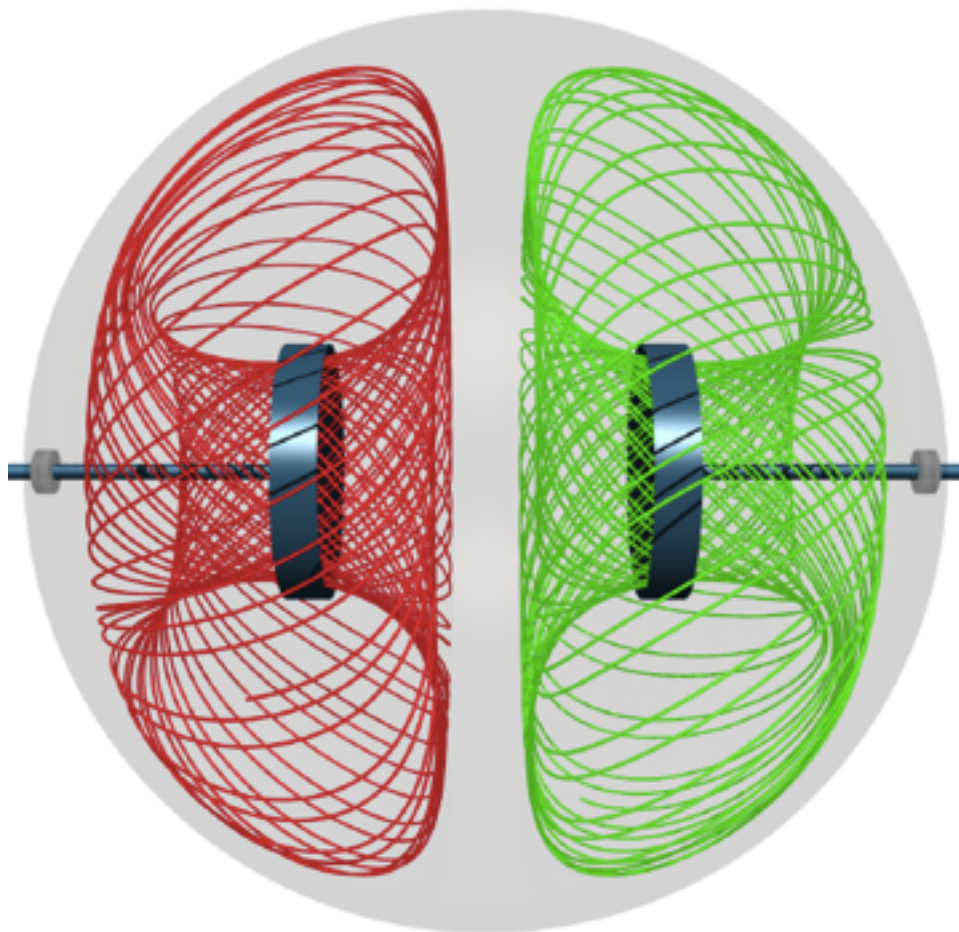
→ Study confinement and stirring in an unmagnetized plasma

Why Plasma Dynamos? complements LM experiments, computationally accessible

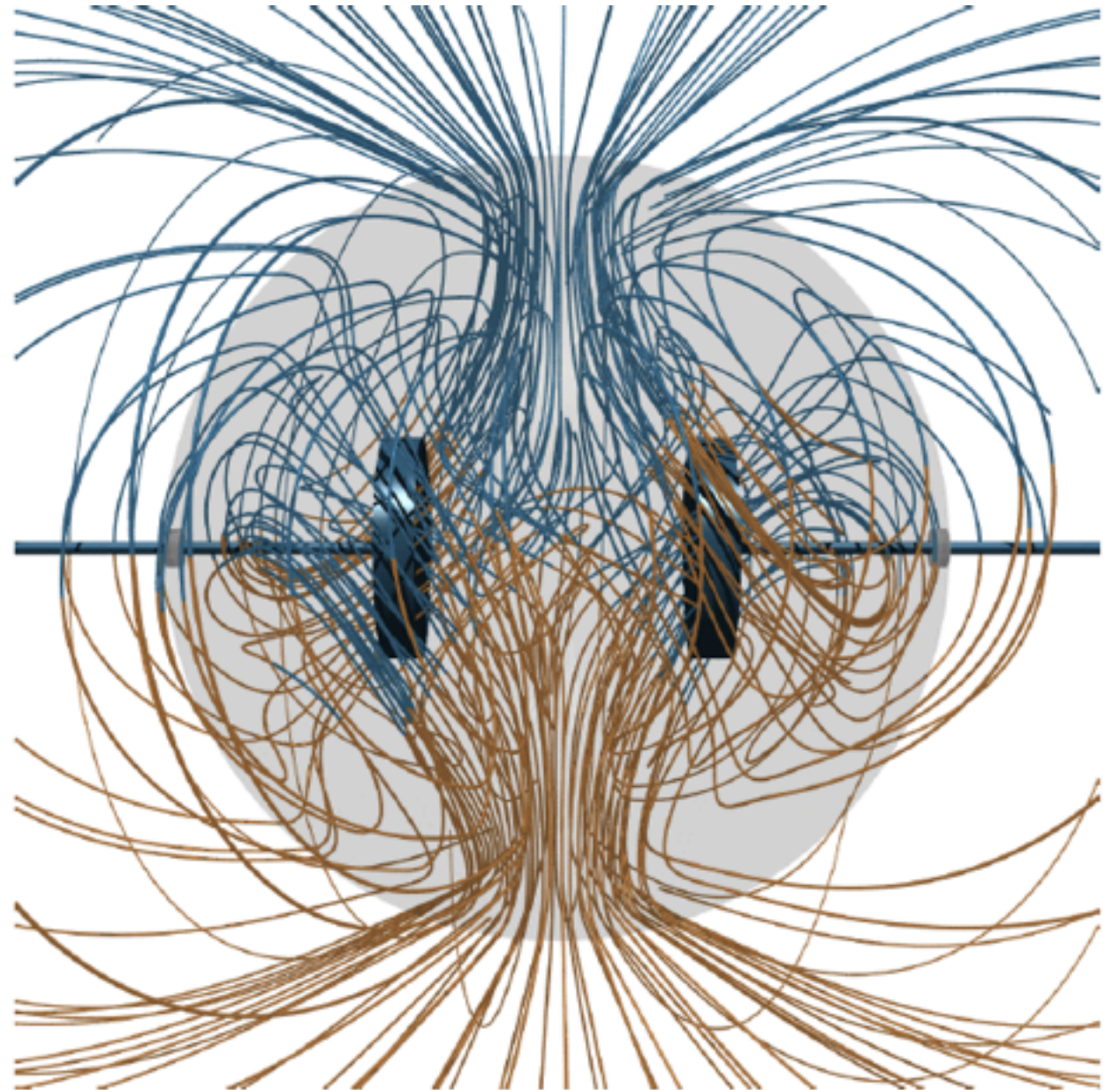


Two vortex dynamo (minimum speed set by $Rm > 50$)

Flow

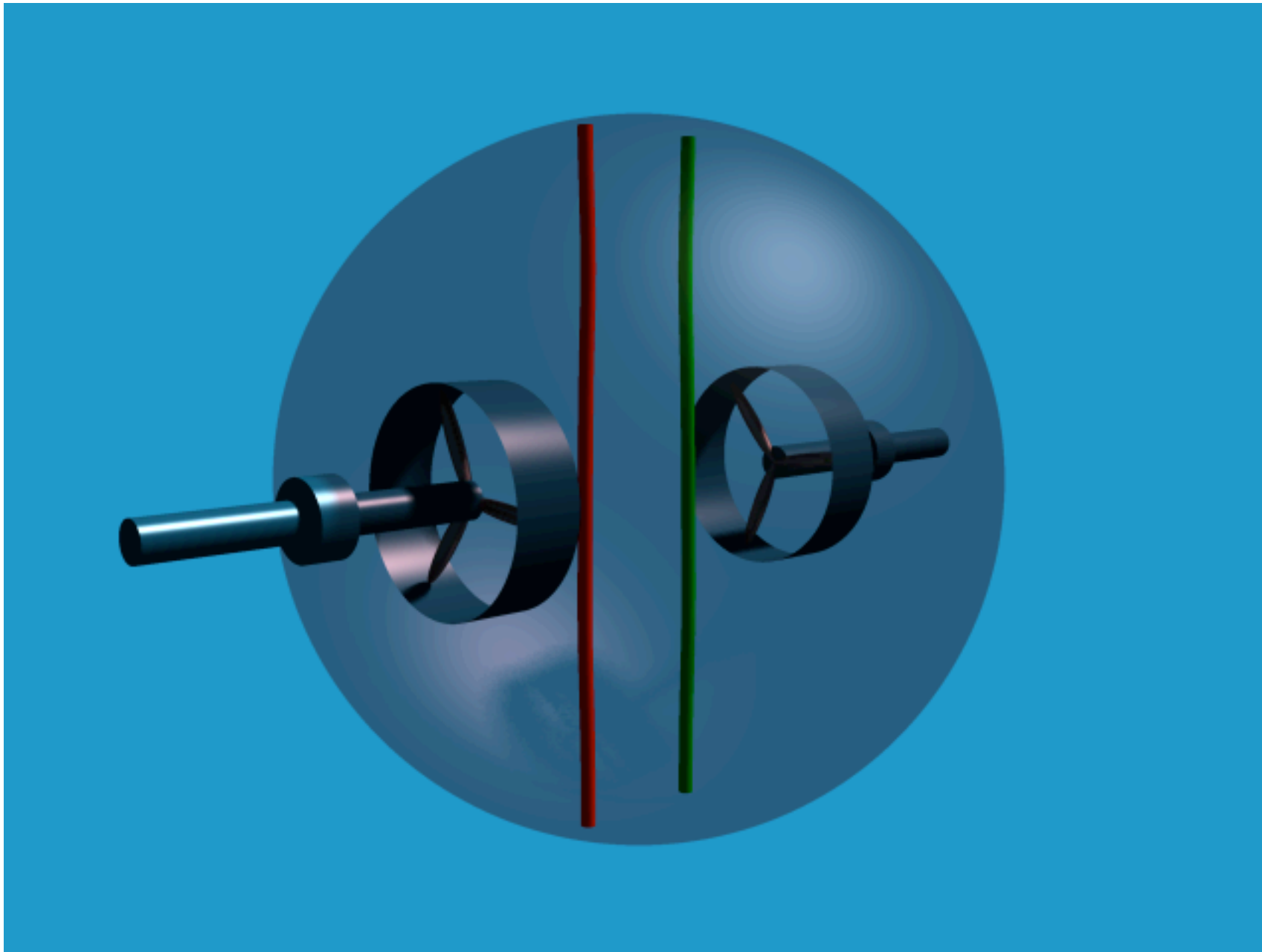


Magnetic Field

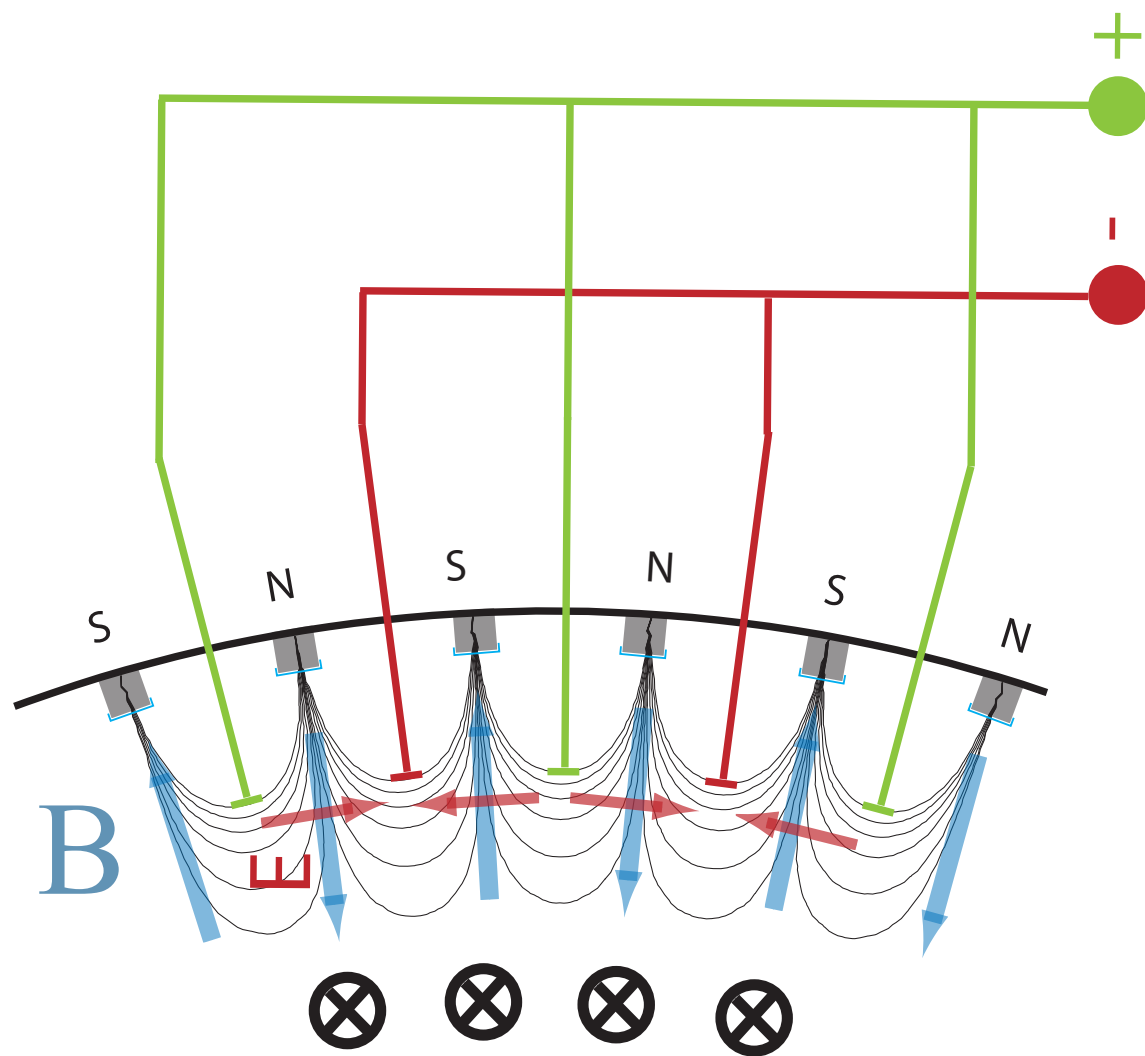


For sodium: $Rm = 12 \text{ a}_{[m]} V_{[m/s]}$

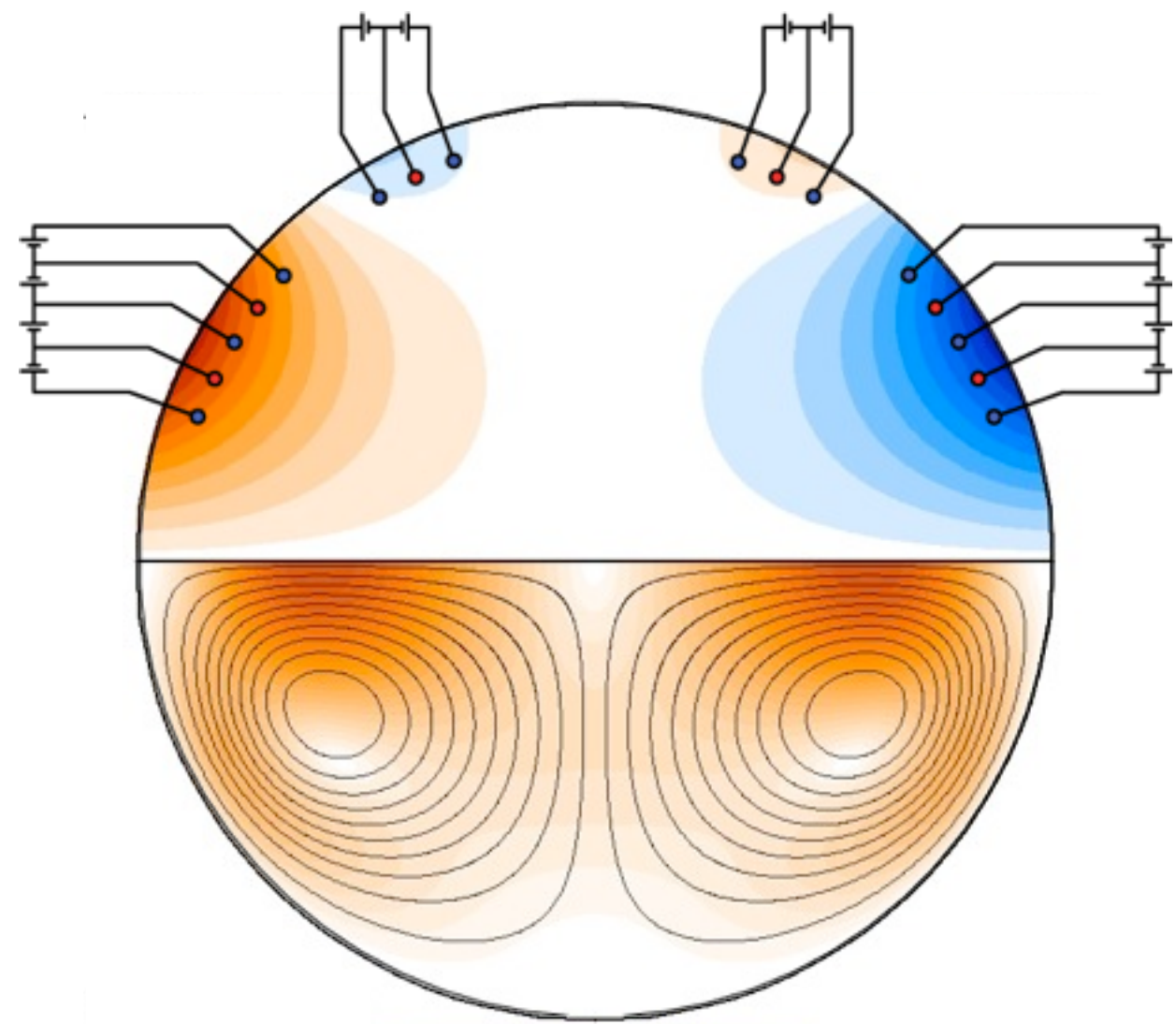
Dynamo is of the stretch-twist-fold type: field line stretching and reinforcement leads to dynamo



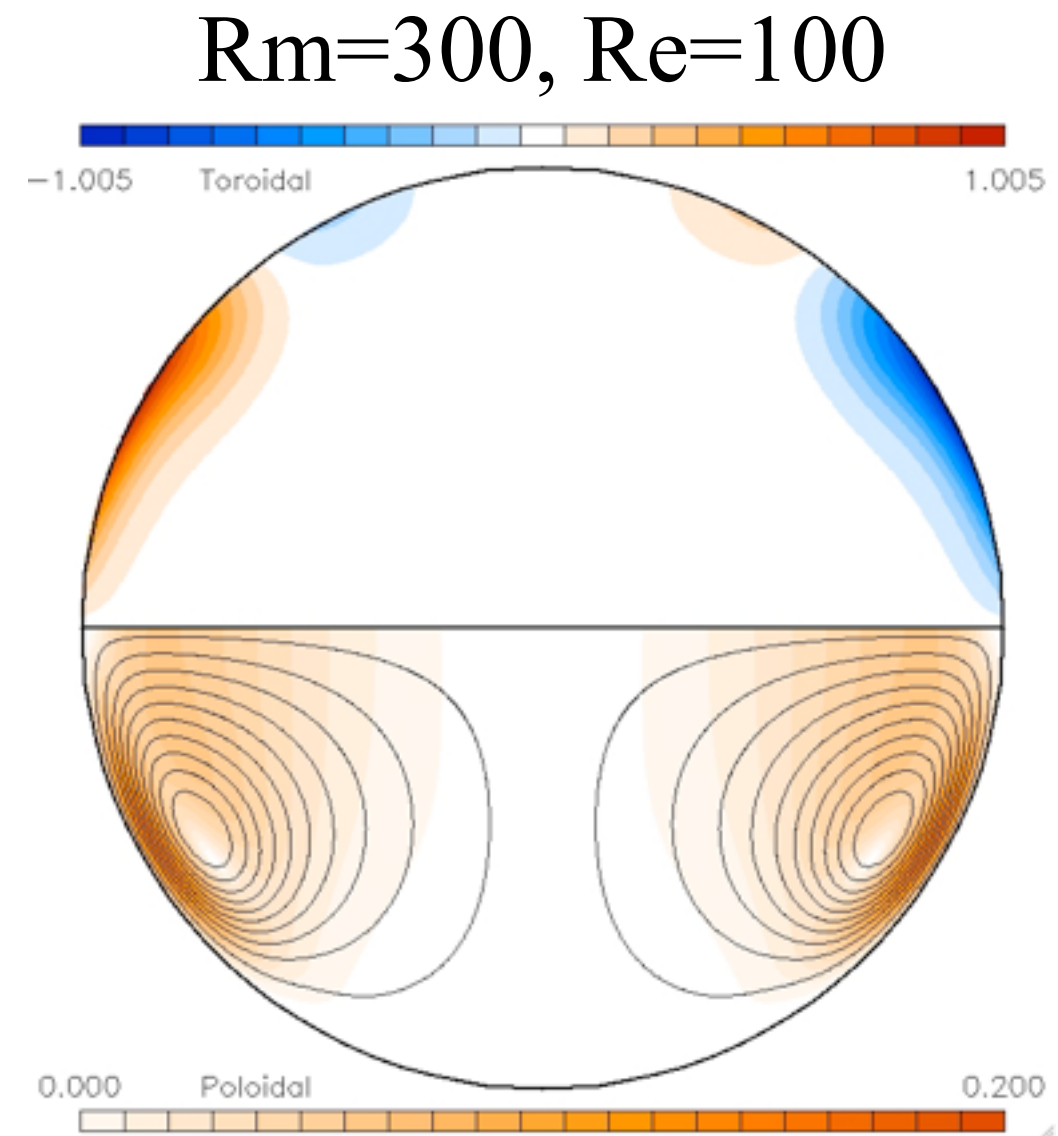
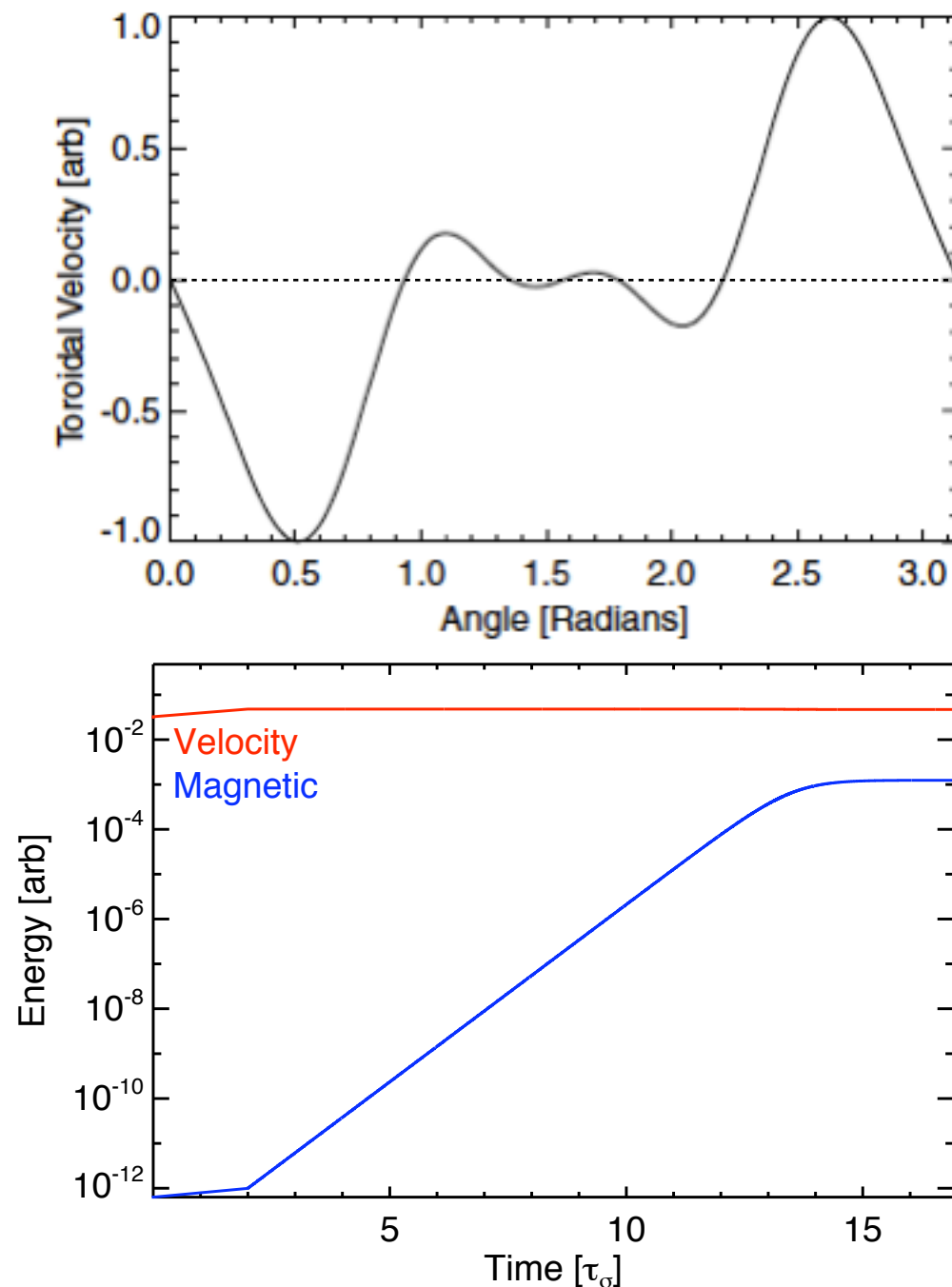
Electrostatic biasing controls edge rotation



Arbitrary
 $V_\phi (r = a, \theta)$



MHD Computation predicts laminar plasma dynamo



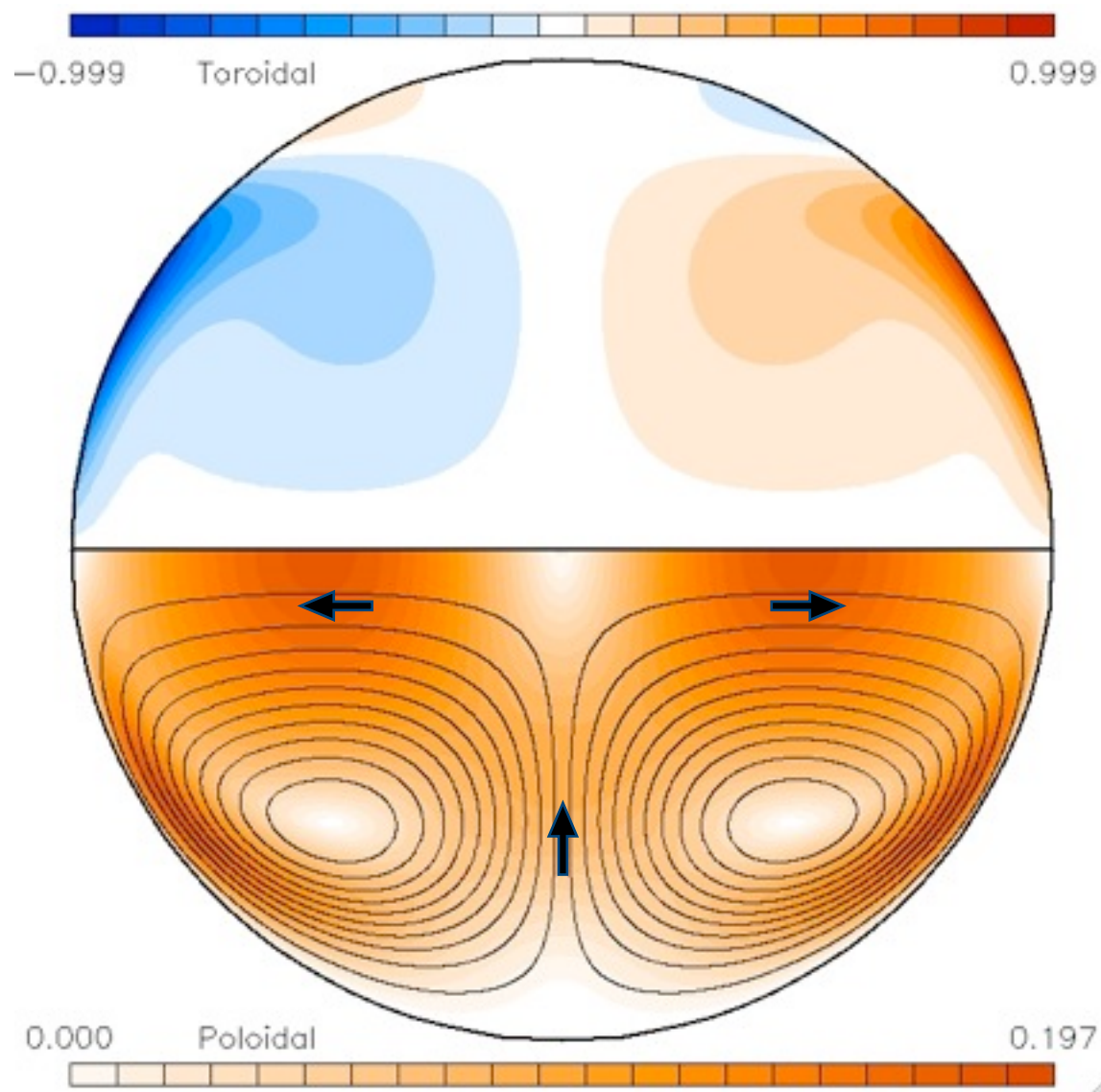
$$T_e = 9 \text{ eV}, \quad n = 8 \times 10^{17} \text{ m}^{-3}$$

$$U_{\text{max}} = 5 \text{ km/s}, \quad \text{Helium}$$

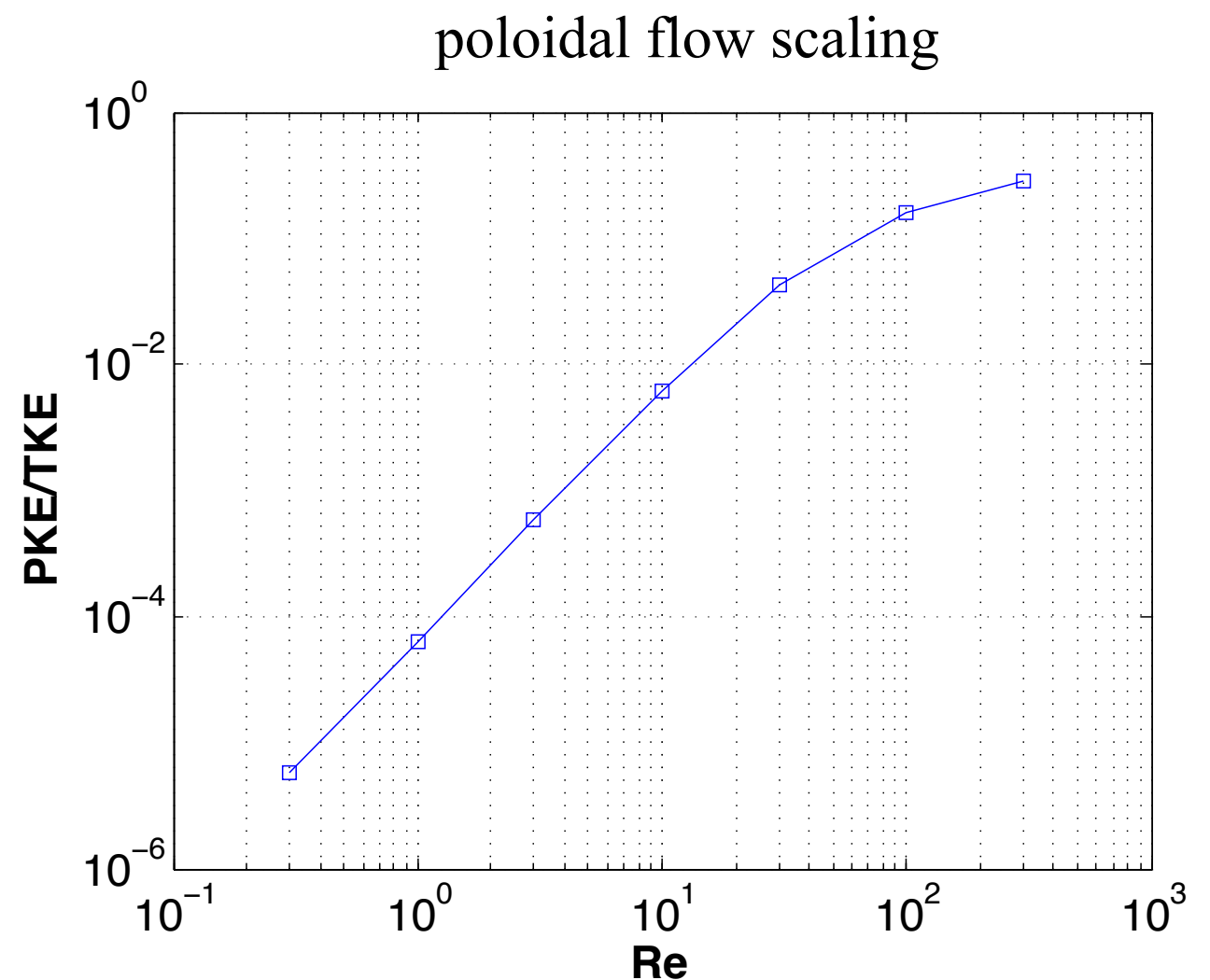
Spence, Reuter, and Forest, *A Spherical Plasma Dynamo Experiment*, The Astrophysical Journal **700** 470 (2009).

Hydrodynamics: toroidal forcing drives poloidal flow

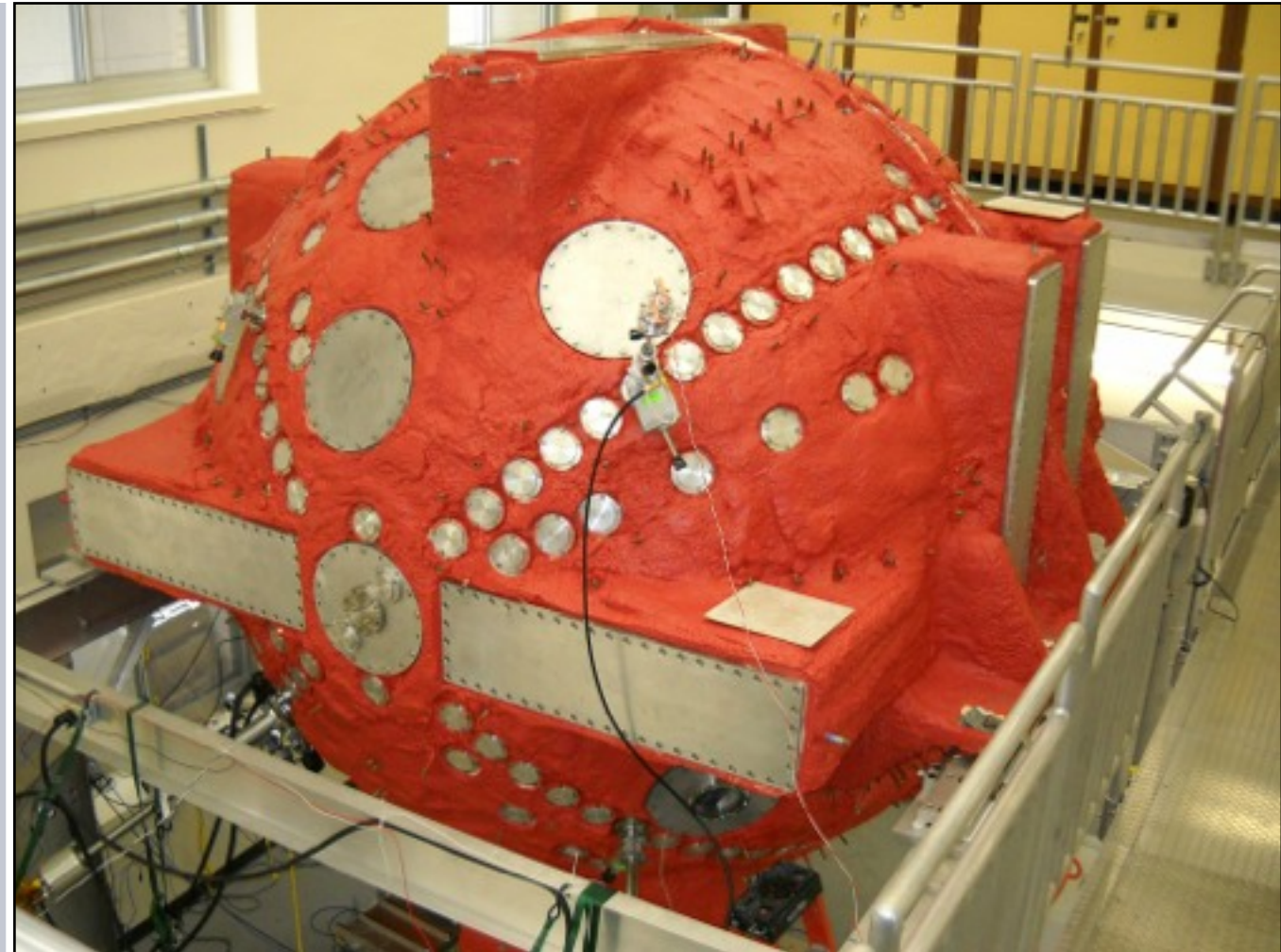
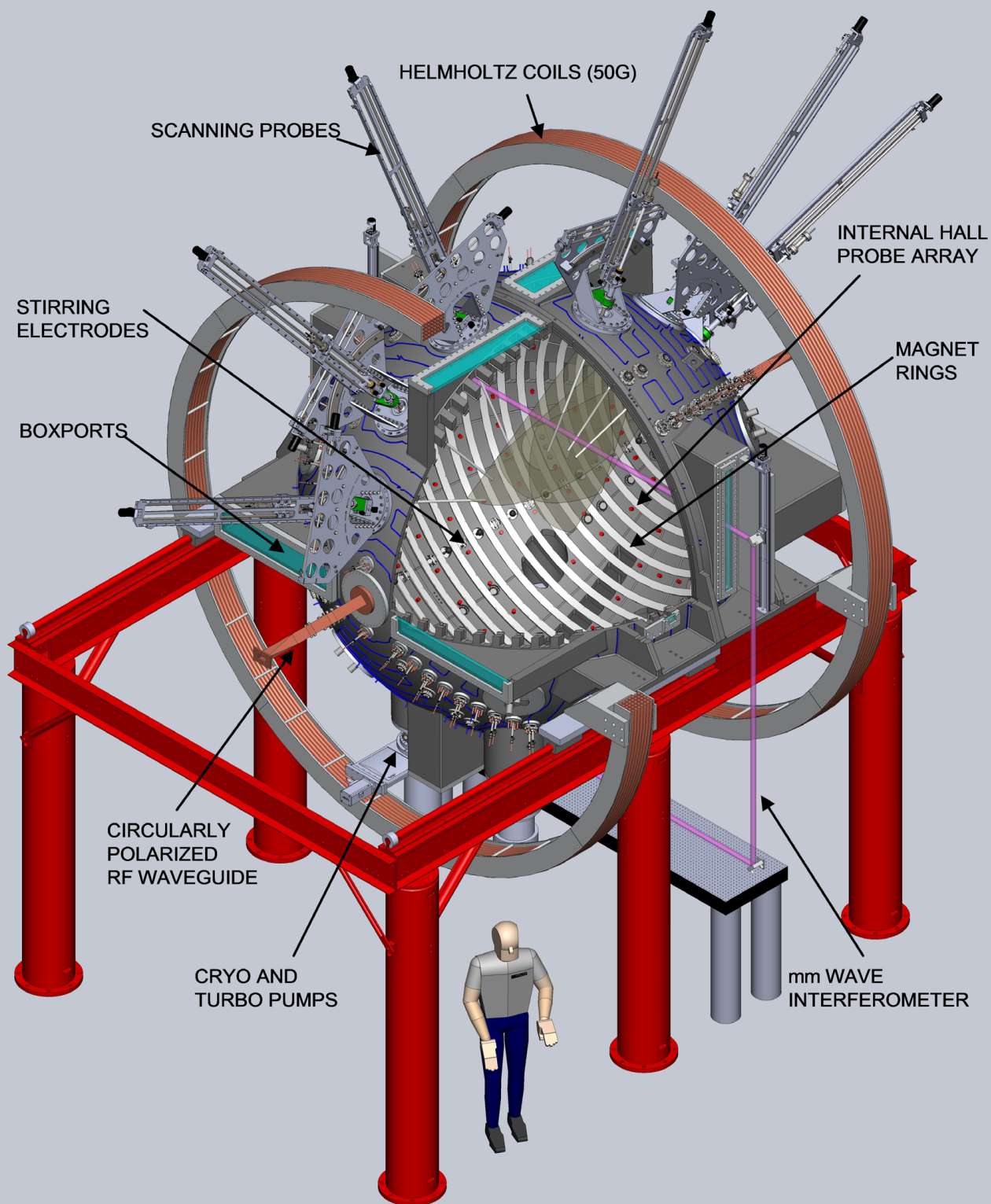
$R_m=300, Re=300$



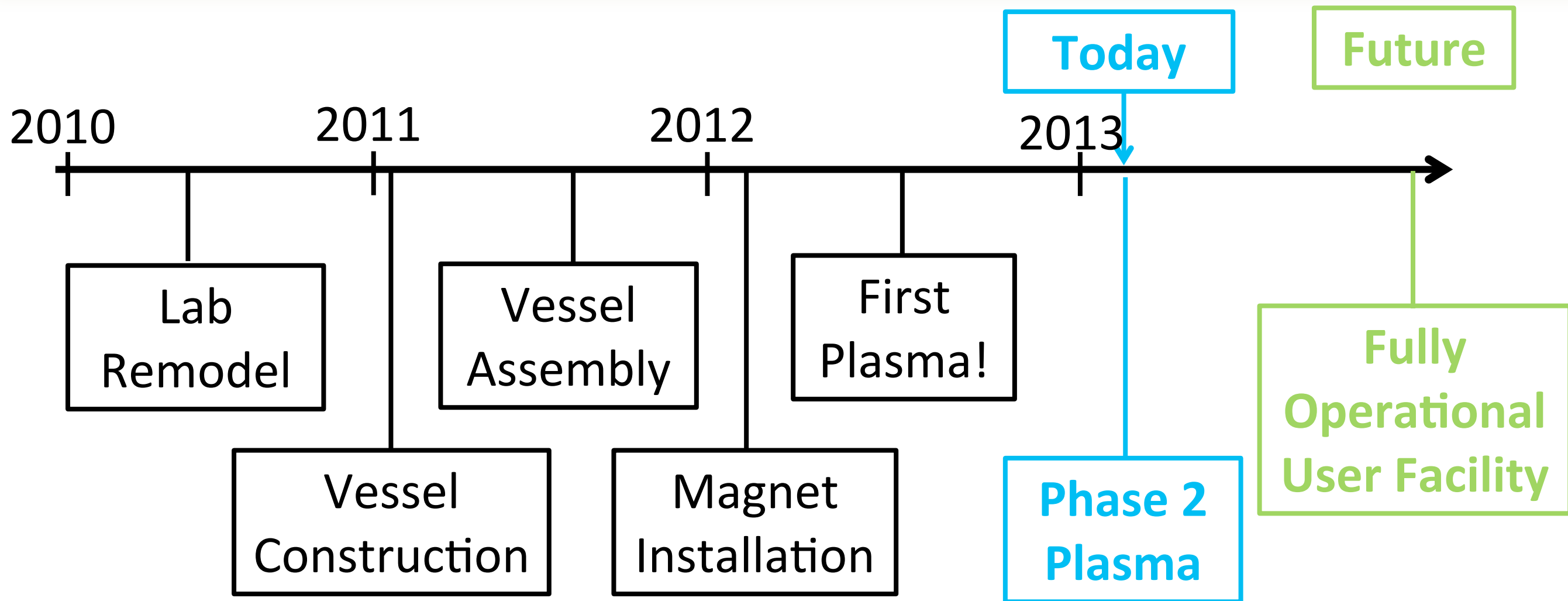
$T_e = 9 \text{ eV}, \quad n = 2.5 \times 10^{18} \text{ m}^{-3}$
 $U_{\max} = 5 \text{ km/s}, \quad \text{Helium}$



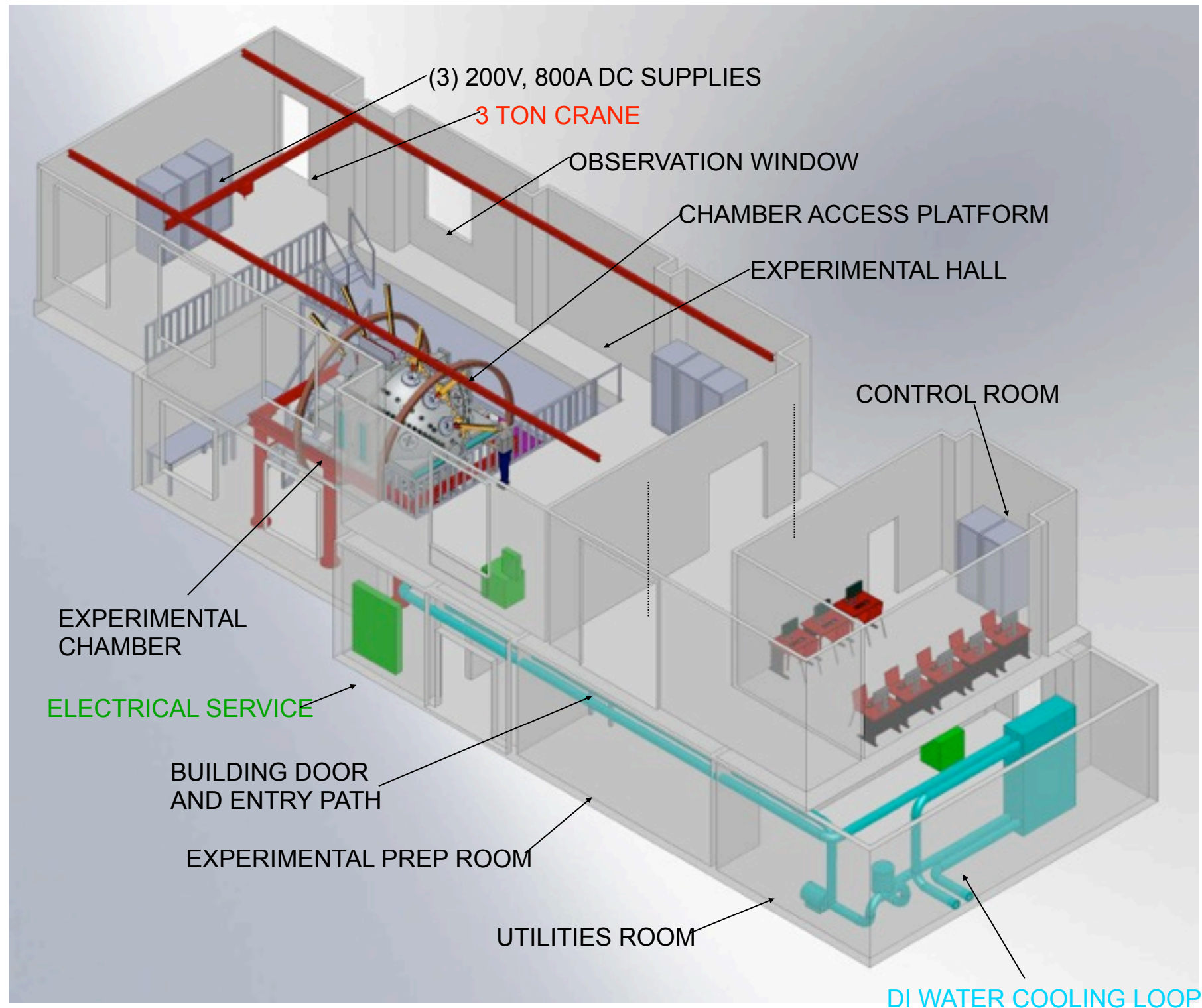
The Plasma Dynamo Experiment



MPDX Timeline



MPDX Facility



Remote operation
from control room

steady-state
operation possible
with installed
cooling

Vessel Construction





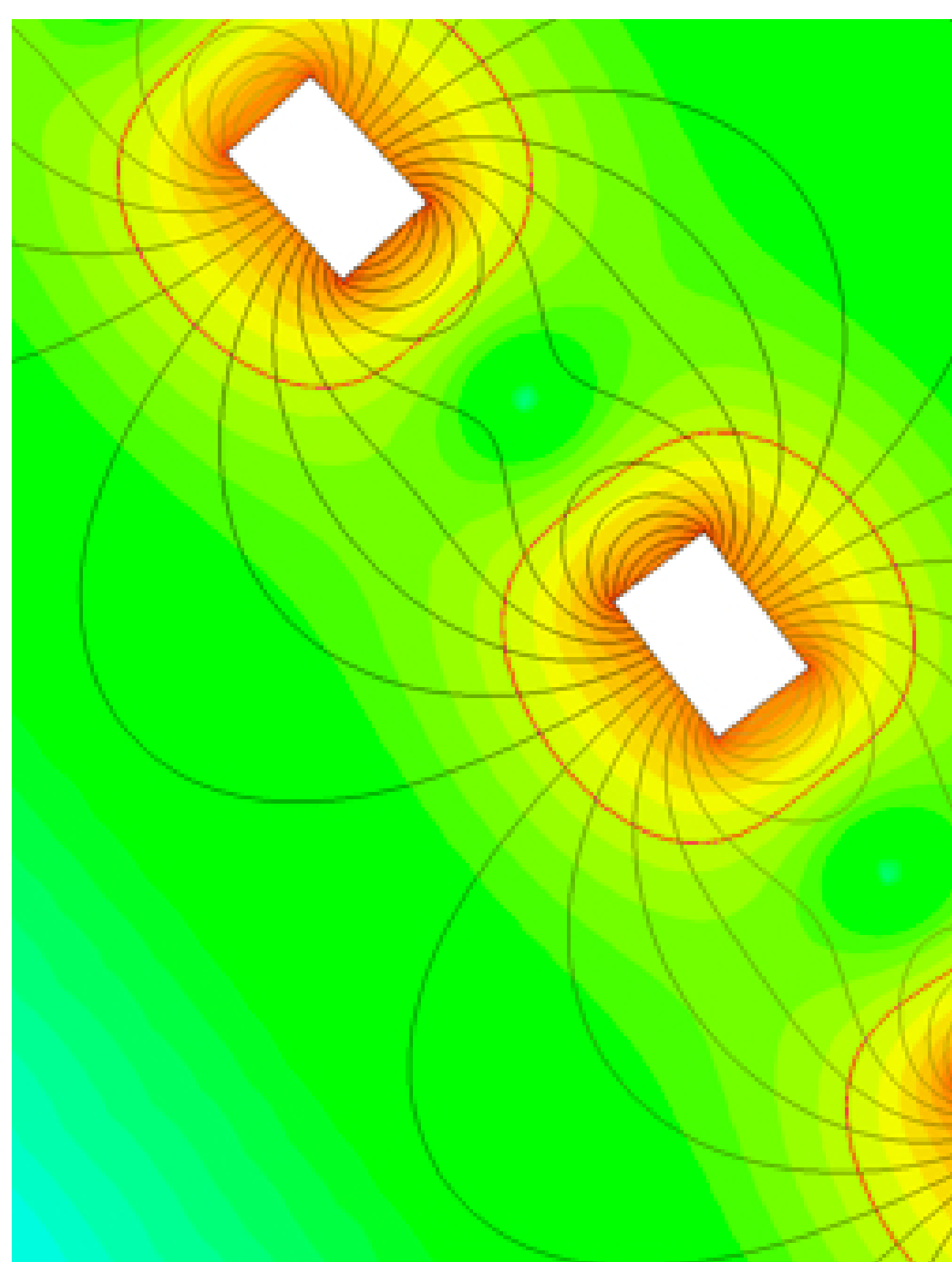
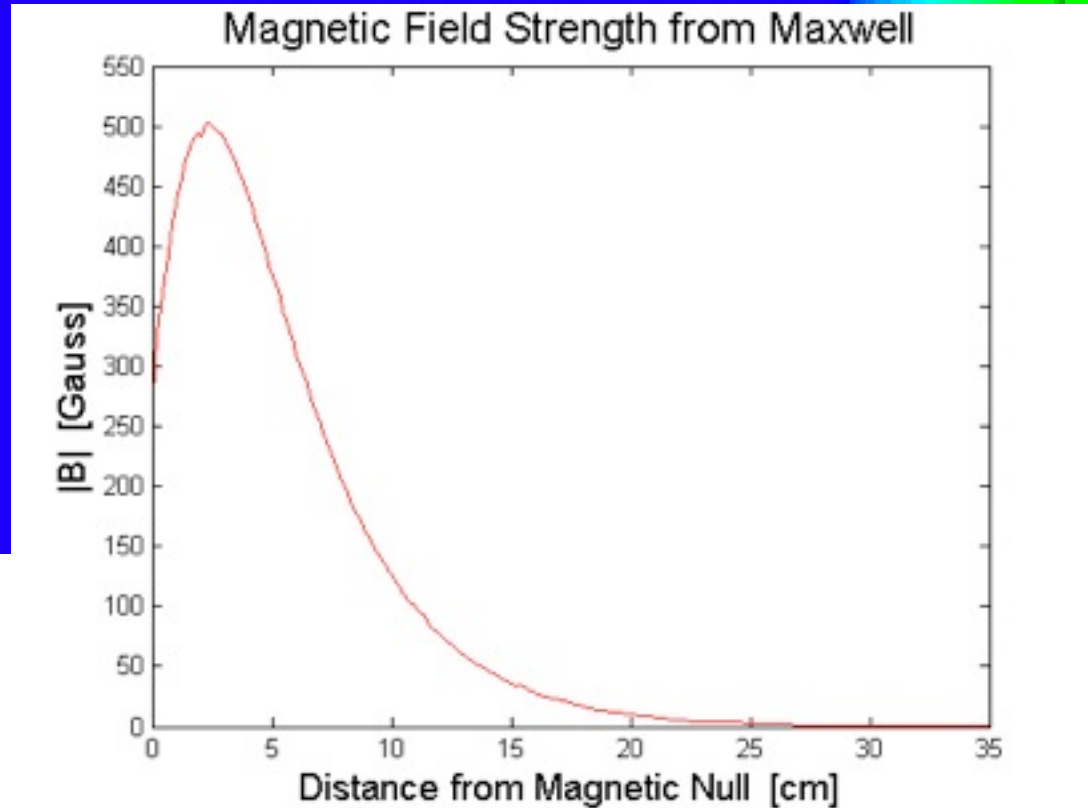
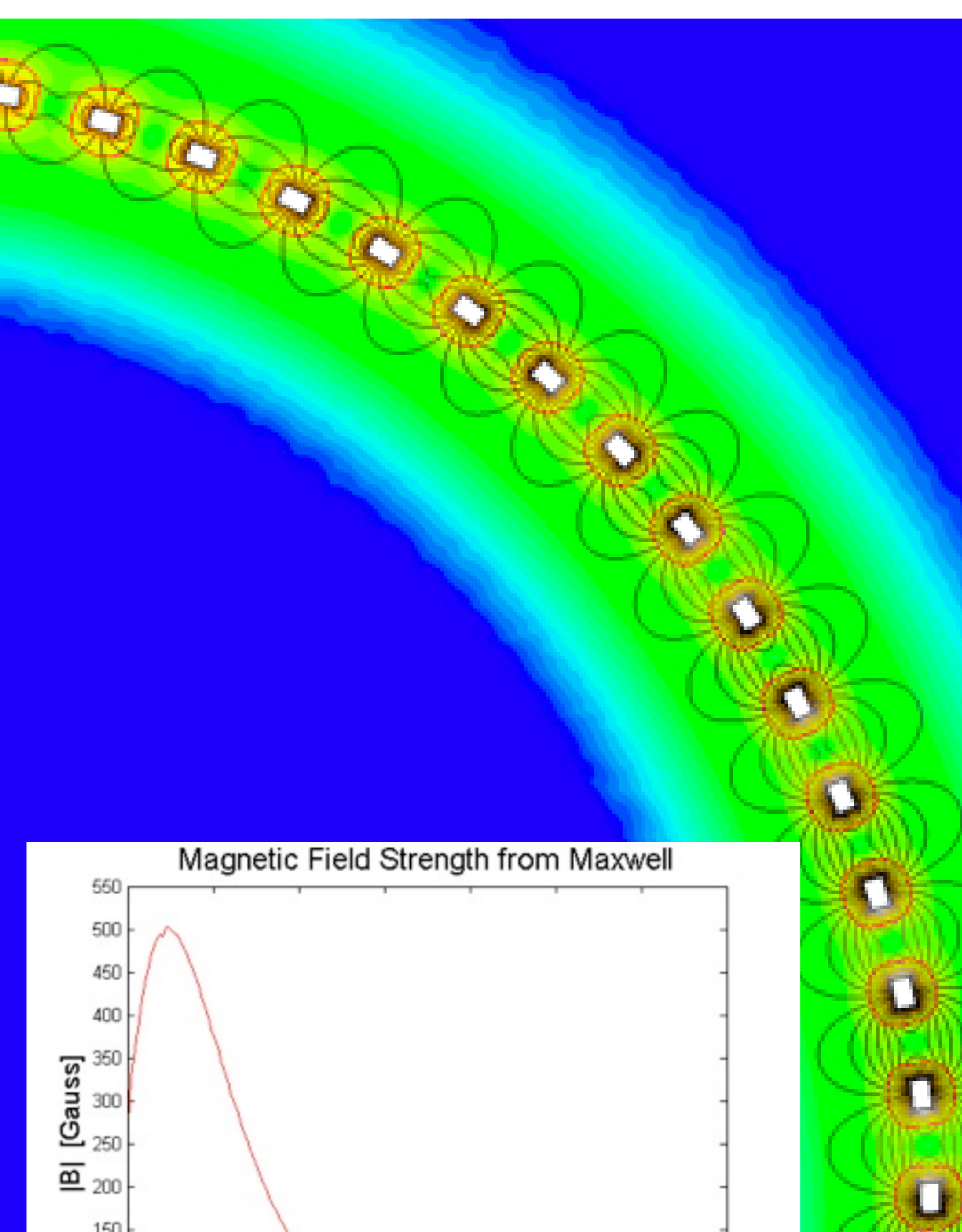


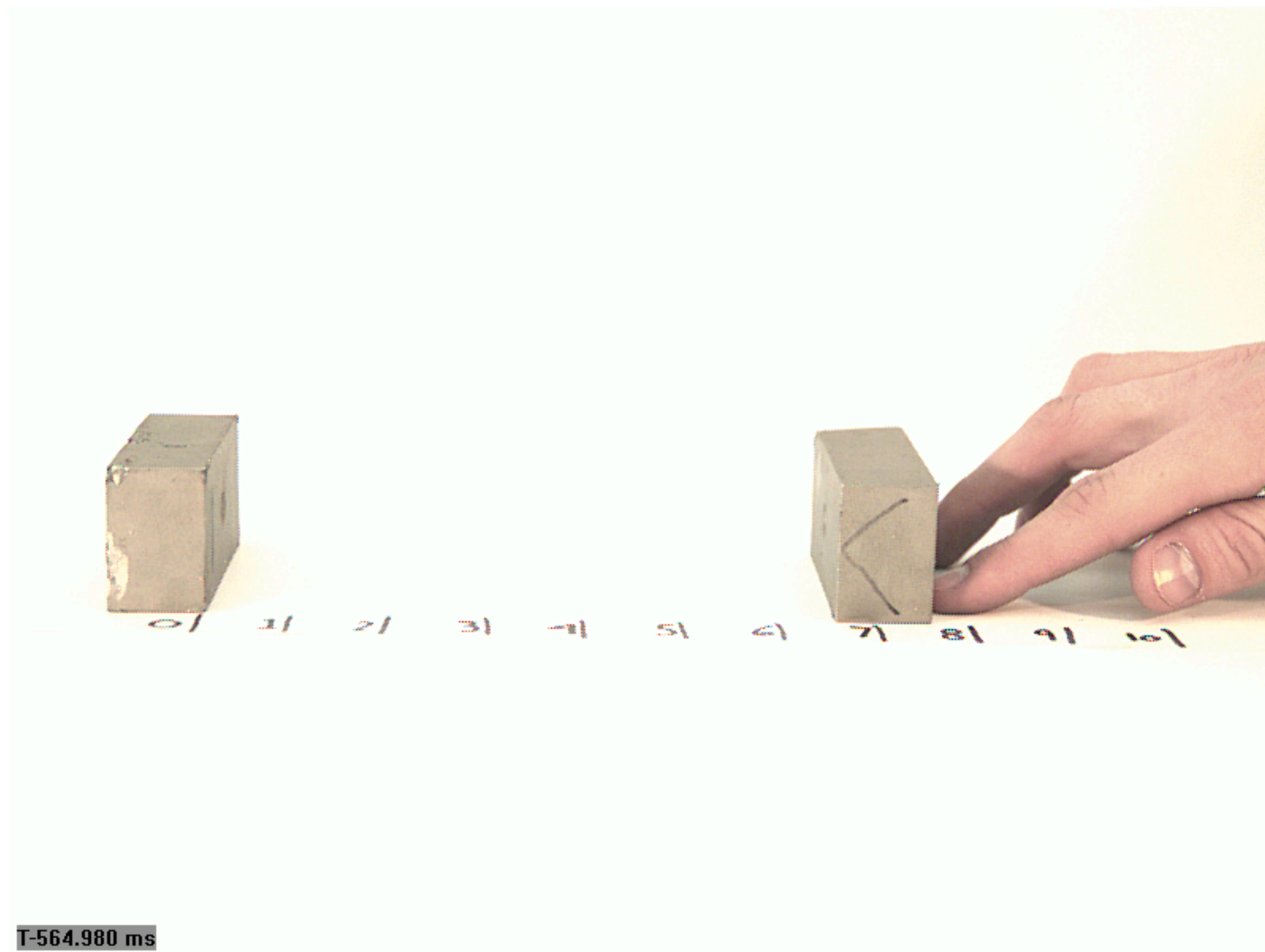


Monday, March 4, 13

3000 4 kG SmCo magnets installed MPDX



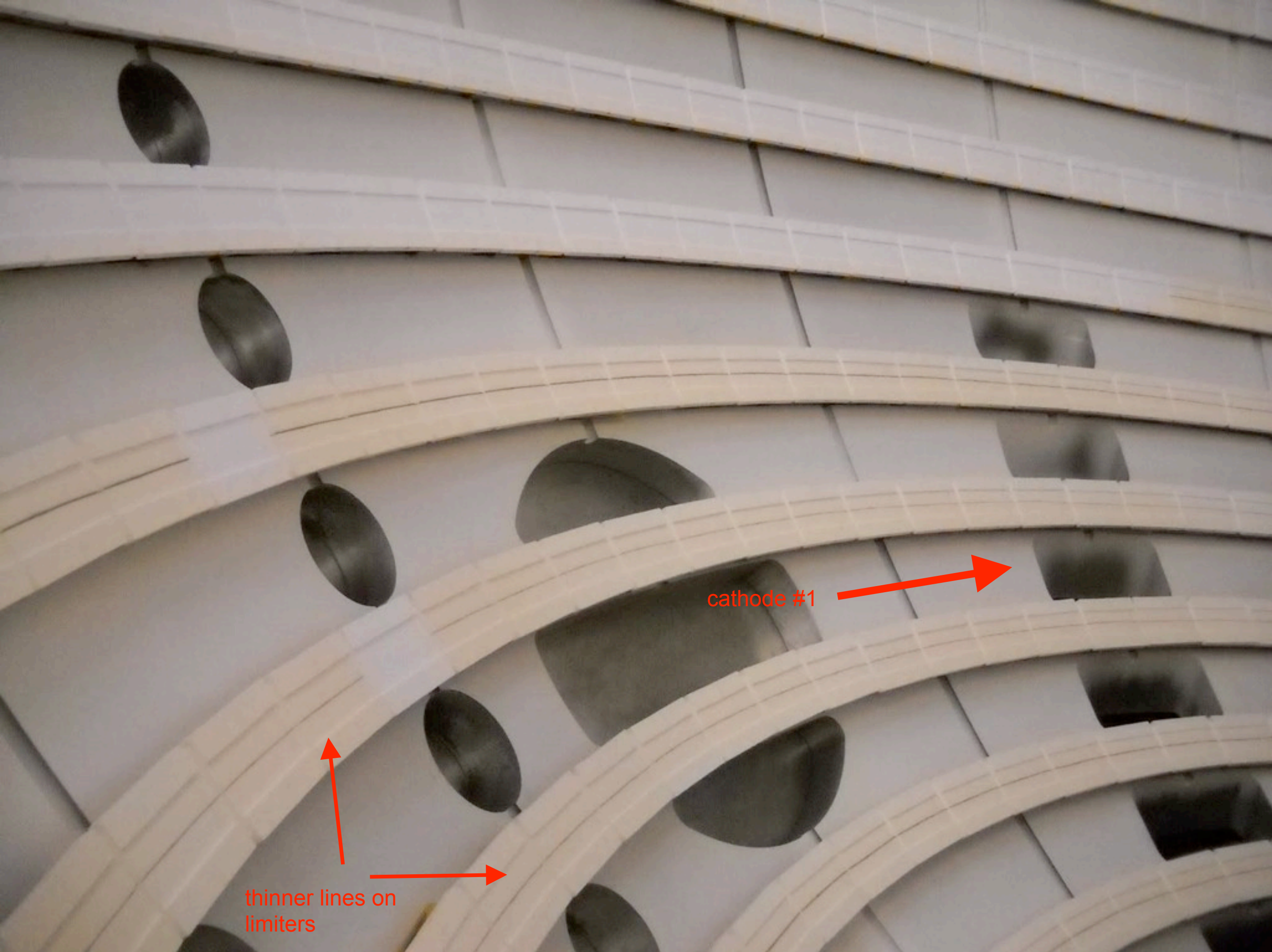




Magnet Installation in MPDX



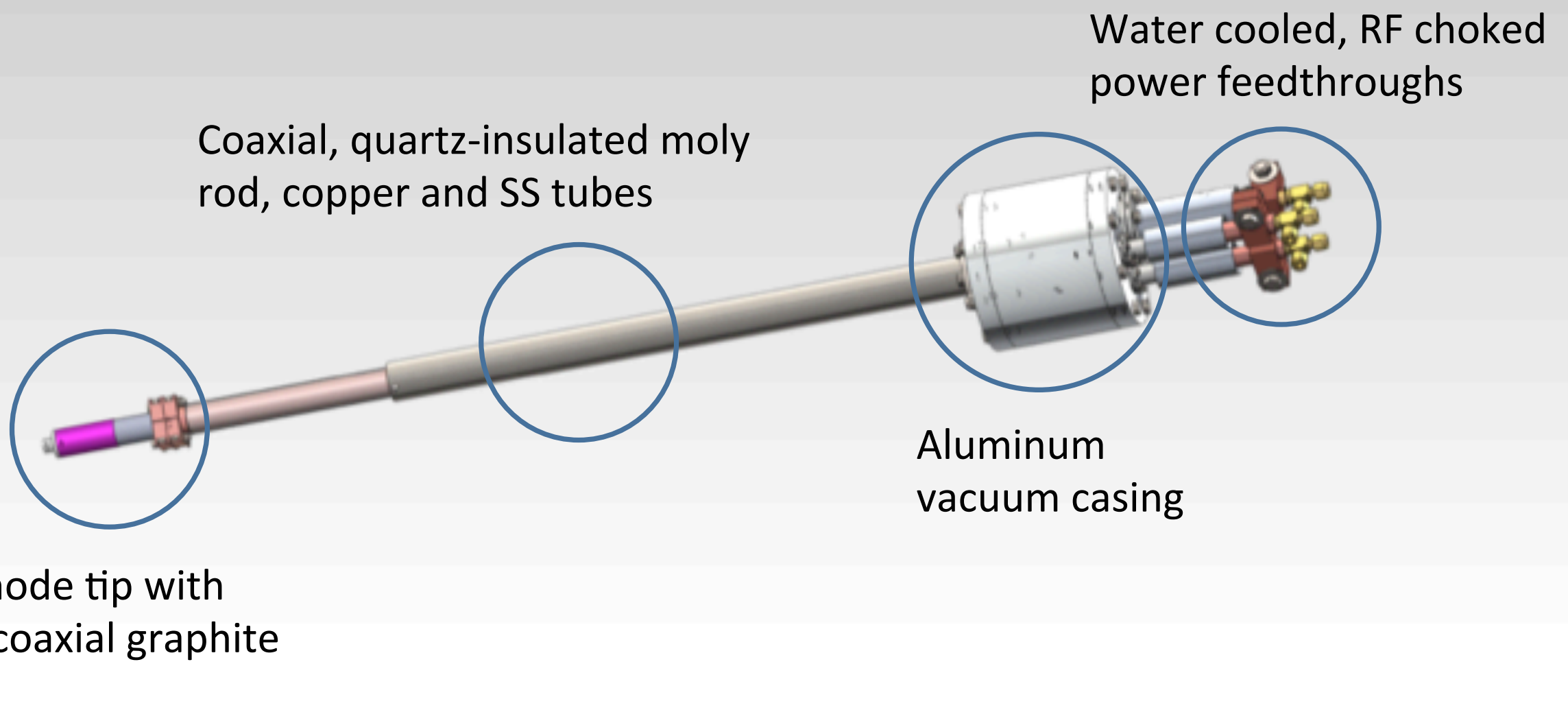
Lever
designed
by John
Wallace



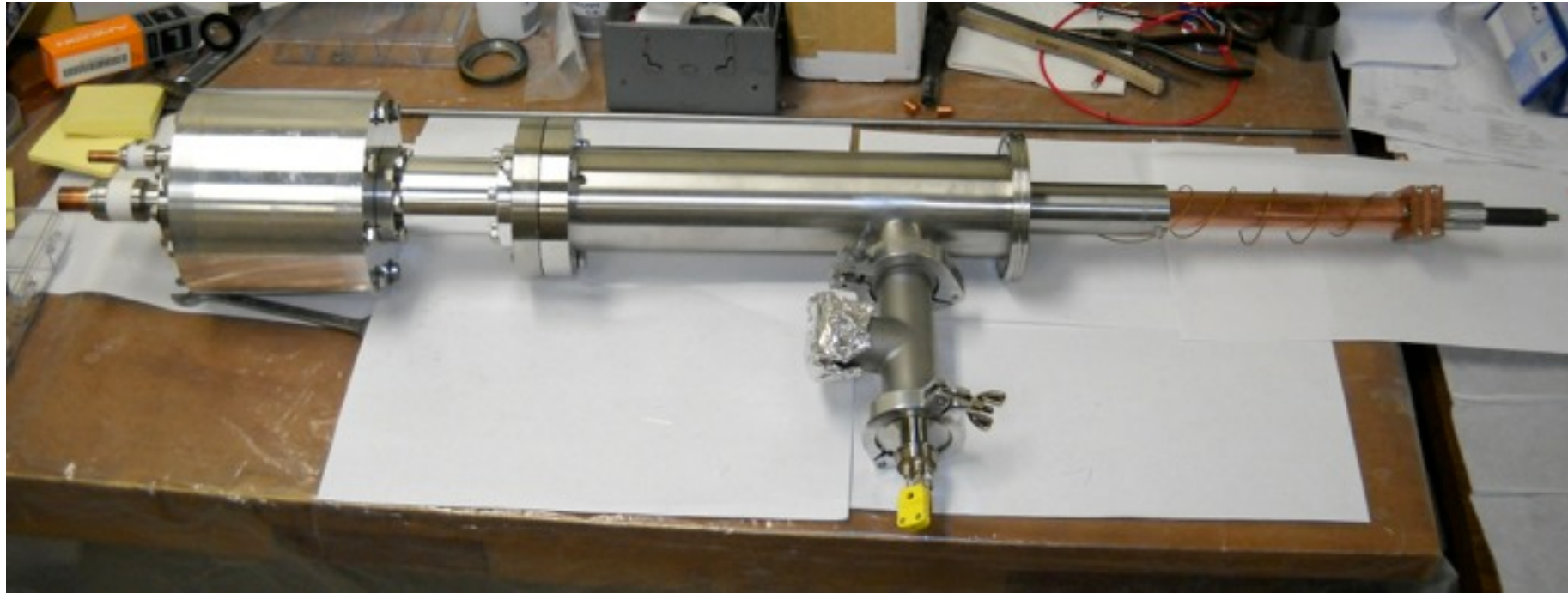
cathode #1

thinner lines on
limiters

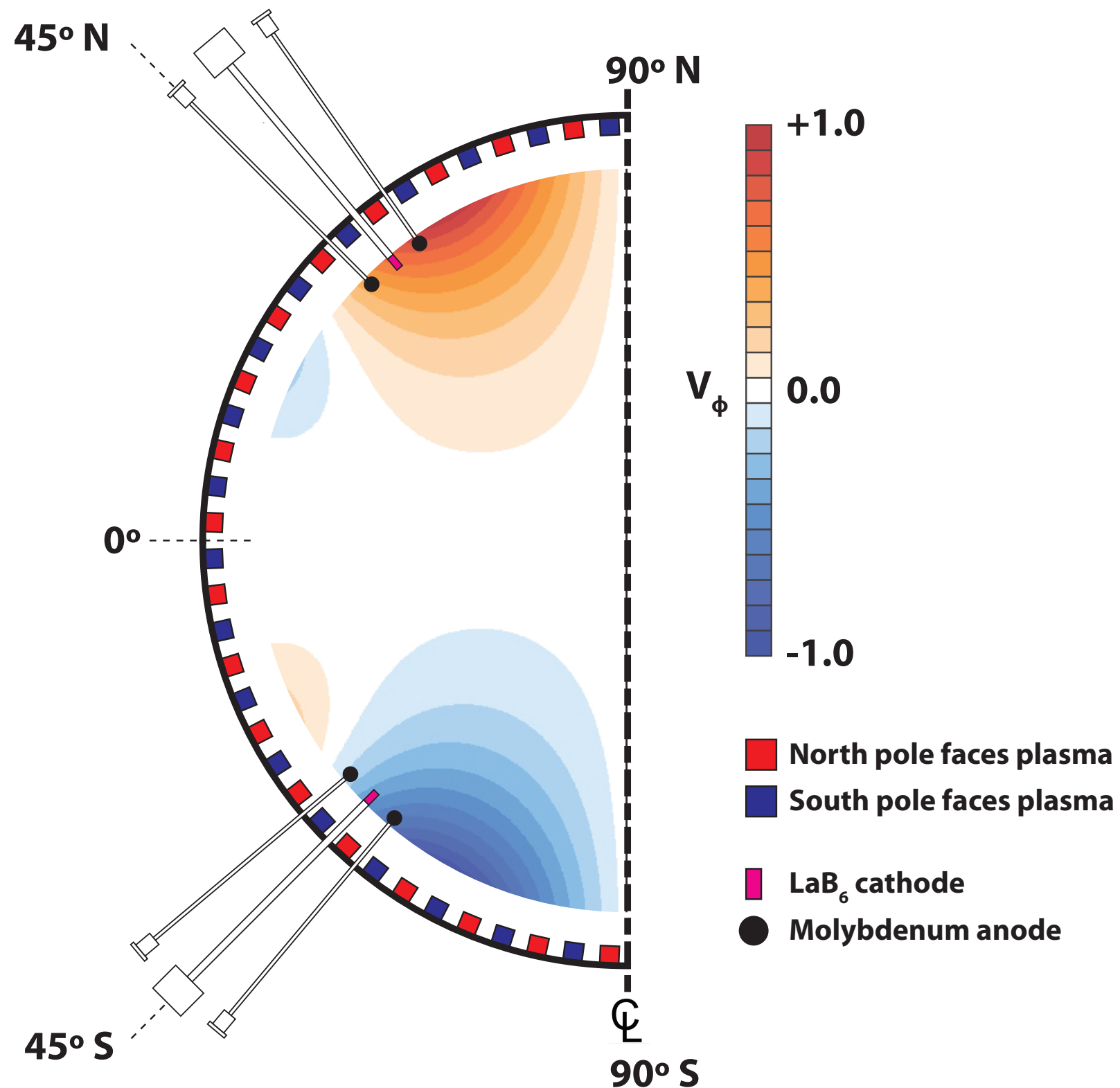
Hot LaB₆ cathodes for heating and stirring plasma



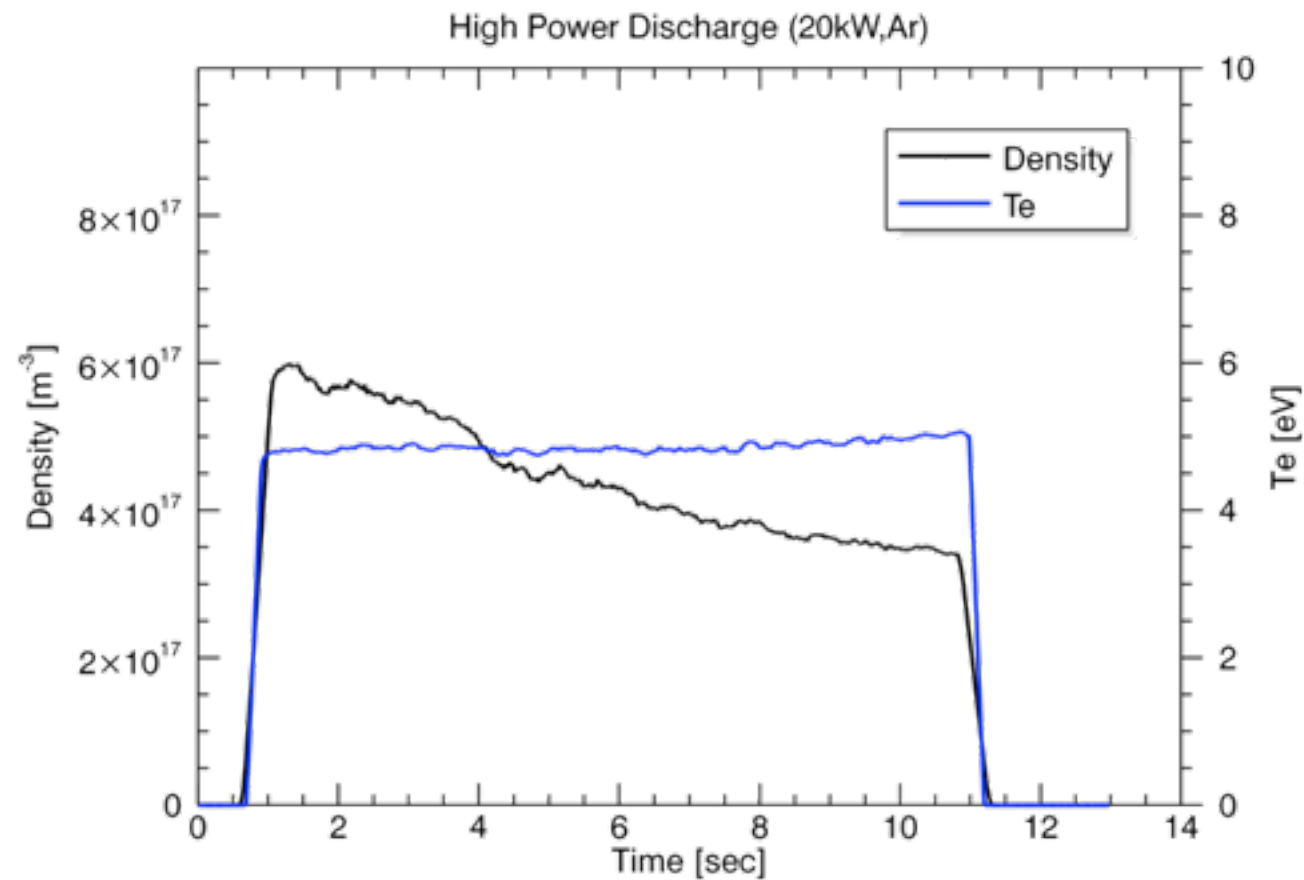
Two cathodes have been constructed



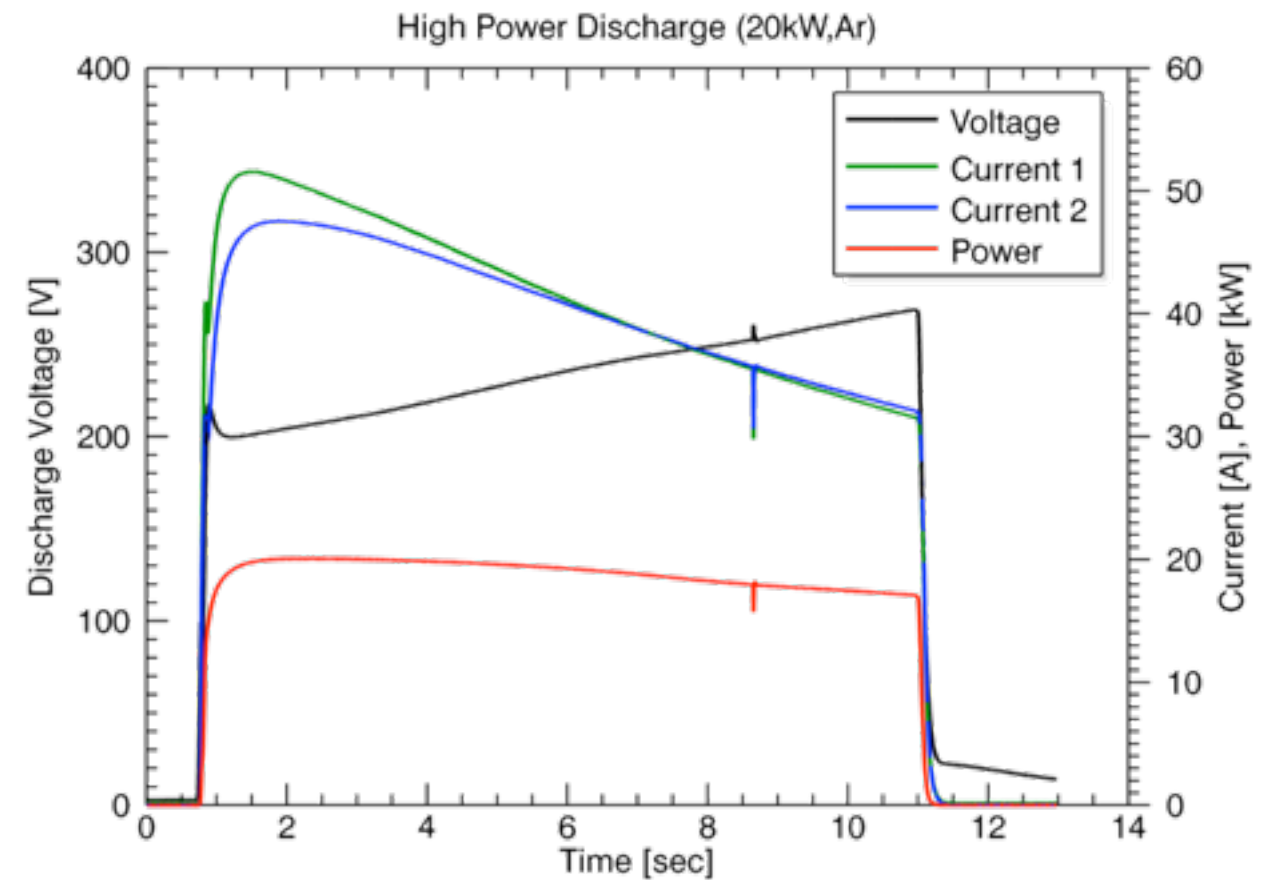
current setup



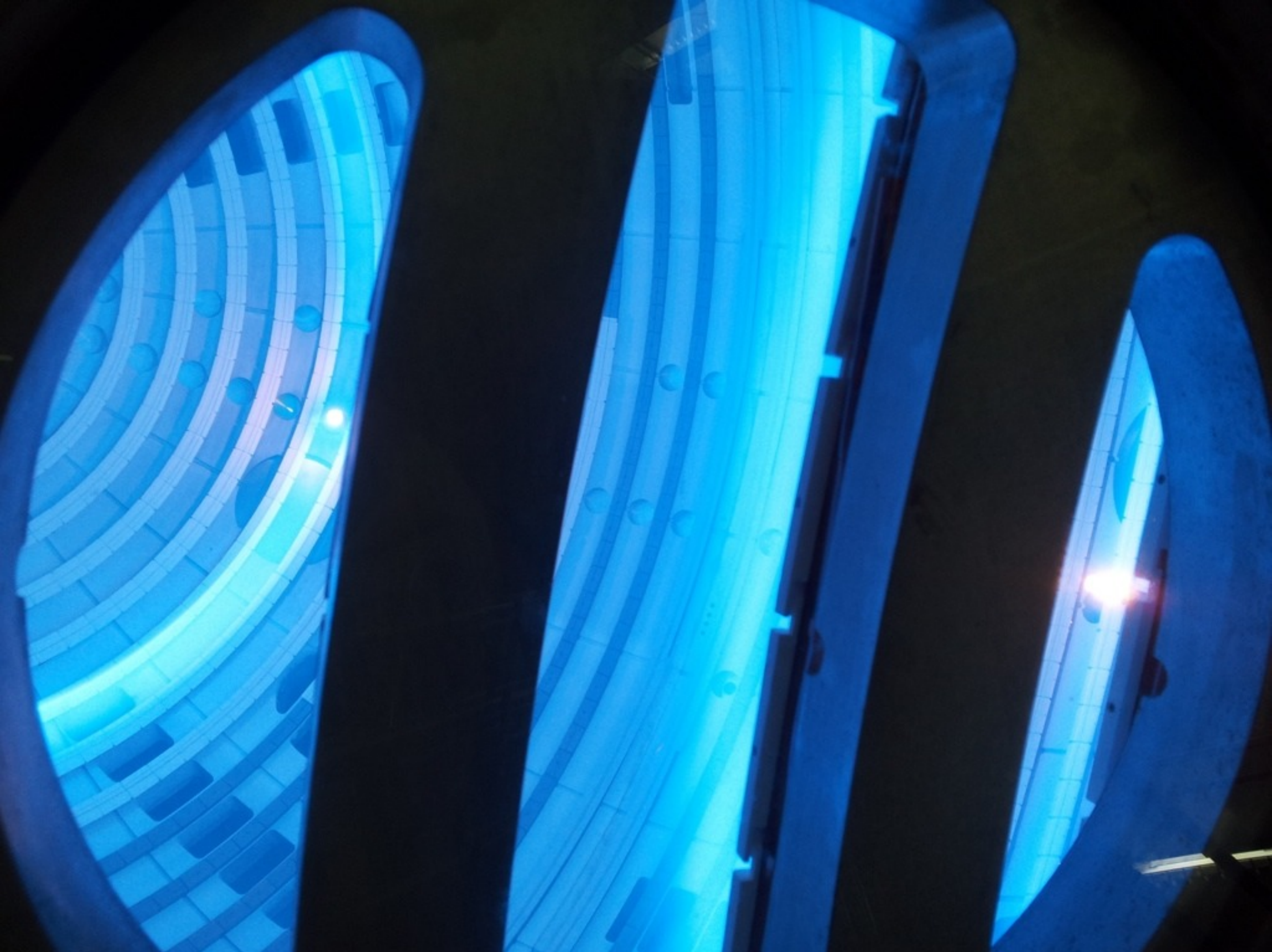
Even without microwave heating, LaB₆ stirring cathodes create hot, dense plasmas



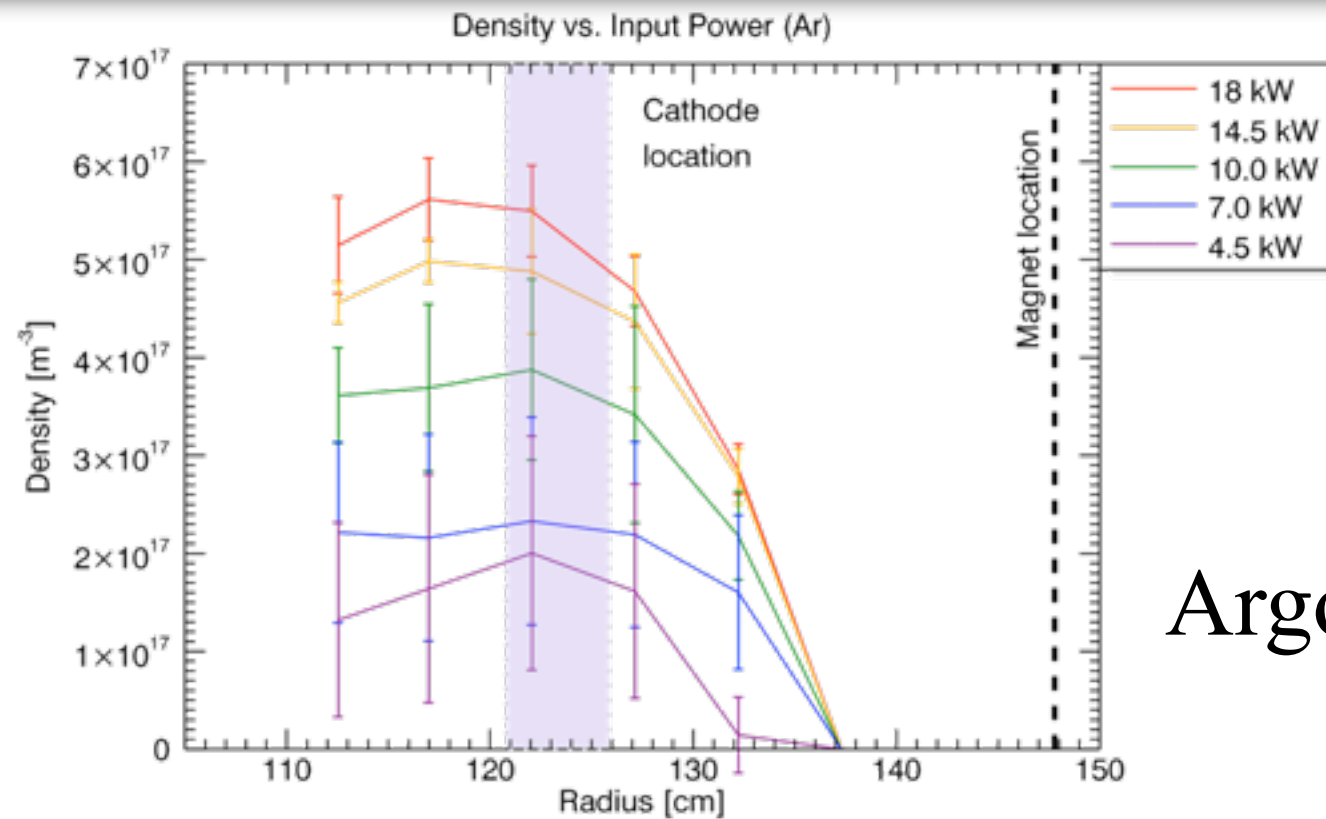
- Discharge powers of up to 20 kW result in $n_e = 5 \times 10^{17} \text{ m}^{-3}$, $T_e = 5 \text{ eV}$ plasmas.



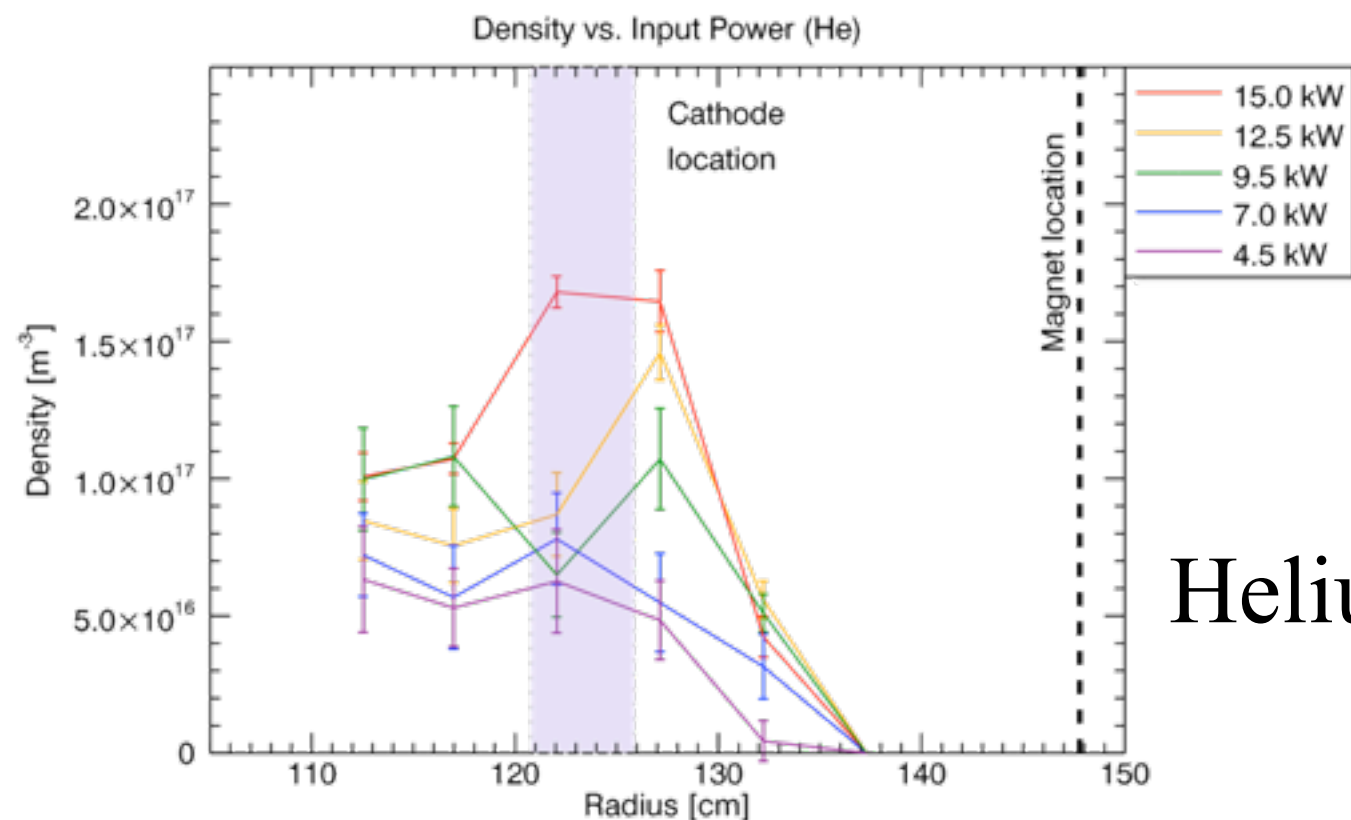
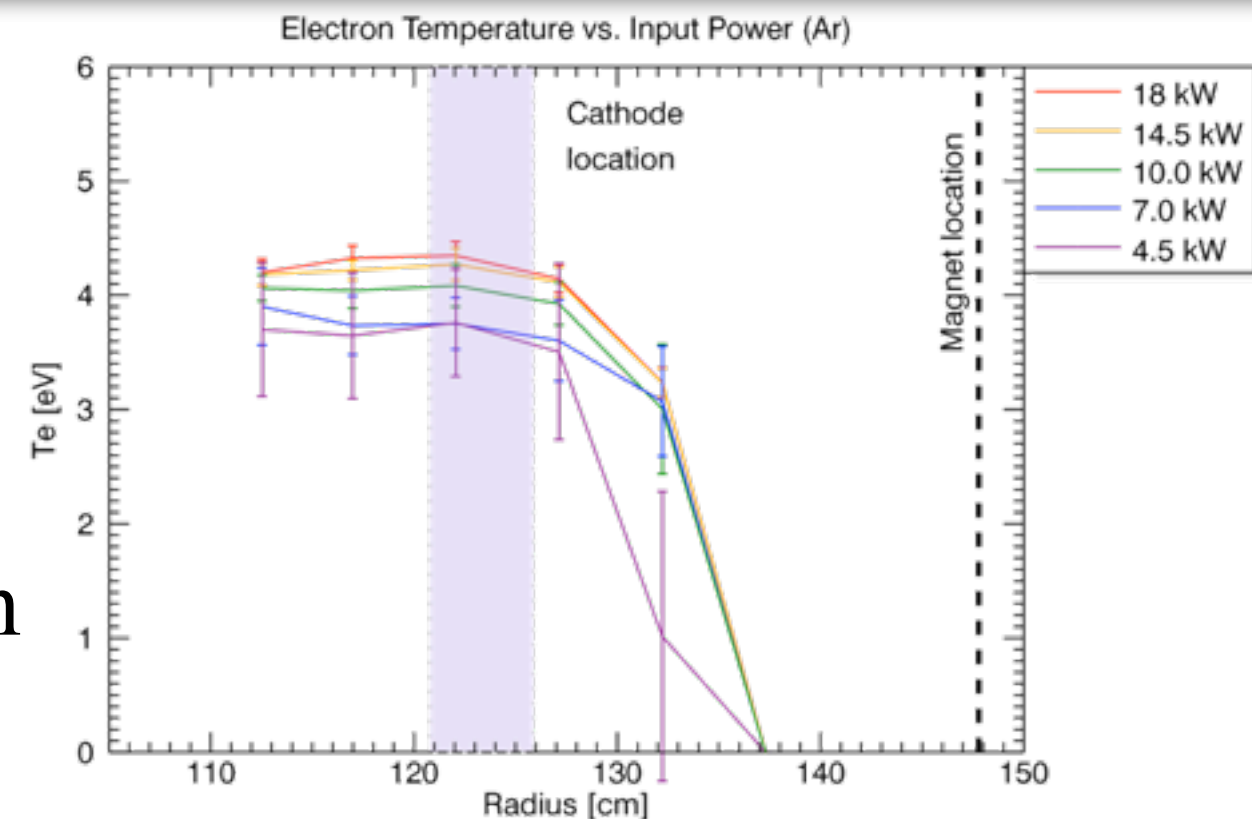
- LaB₆ cathodes draw up to 50A each during each 10 second plasma discharge.



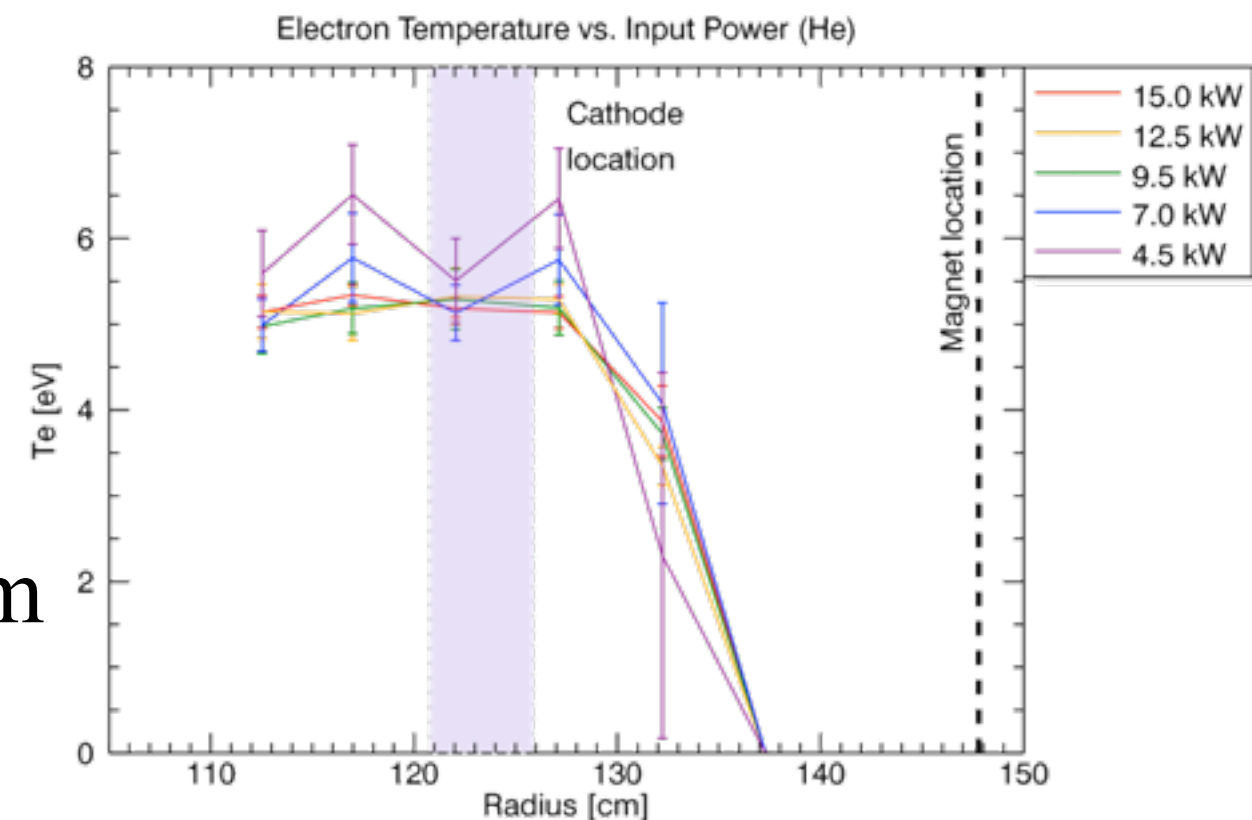
Argon and Helium plasmas are well confined by cusp field



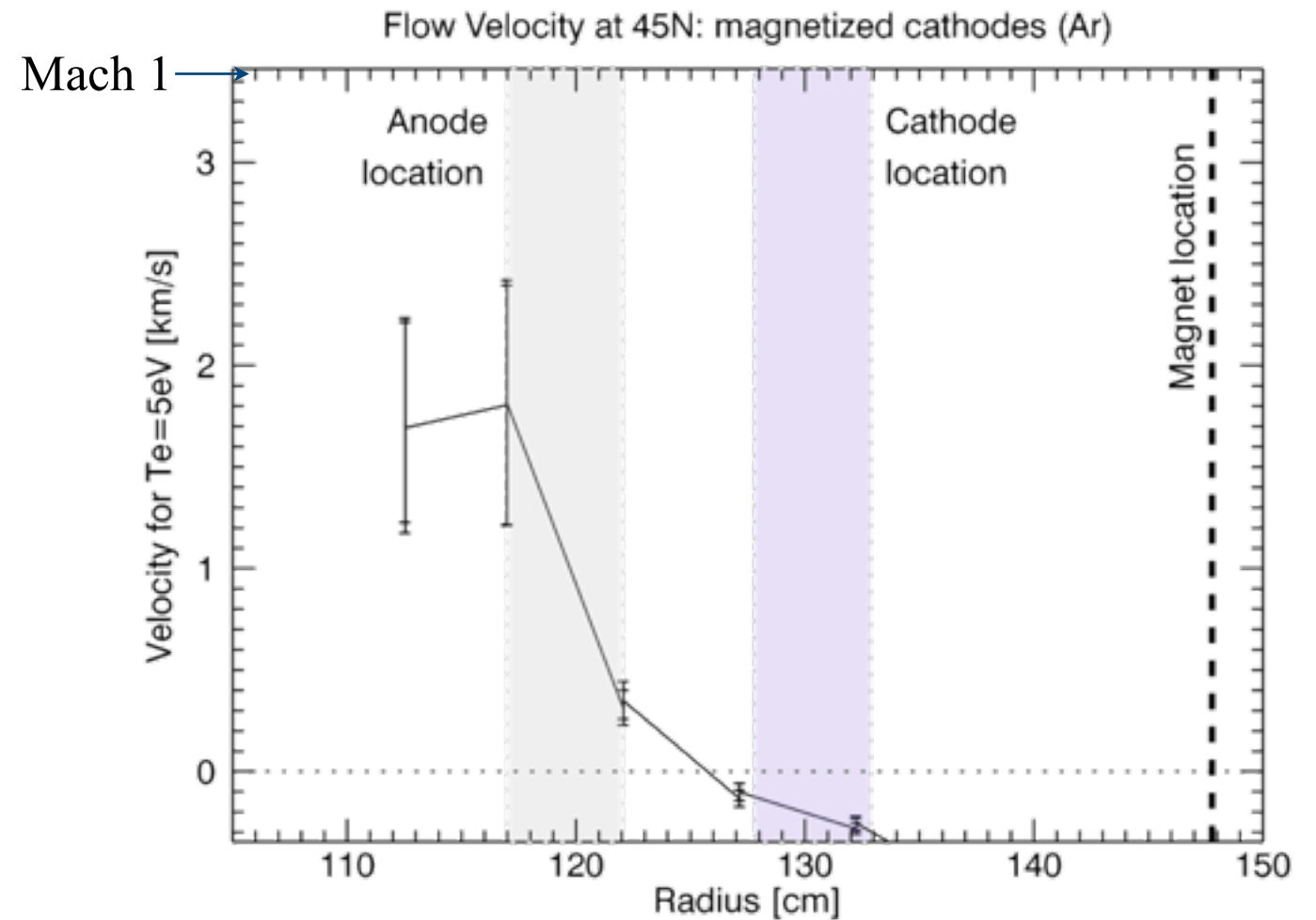
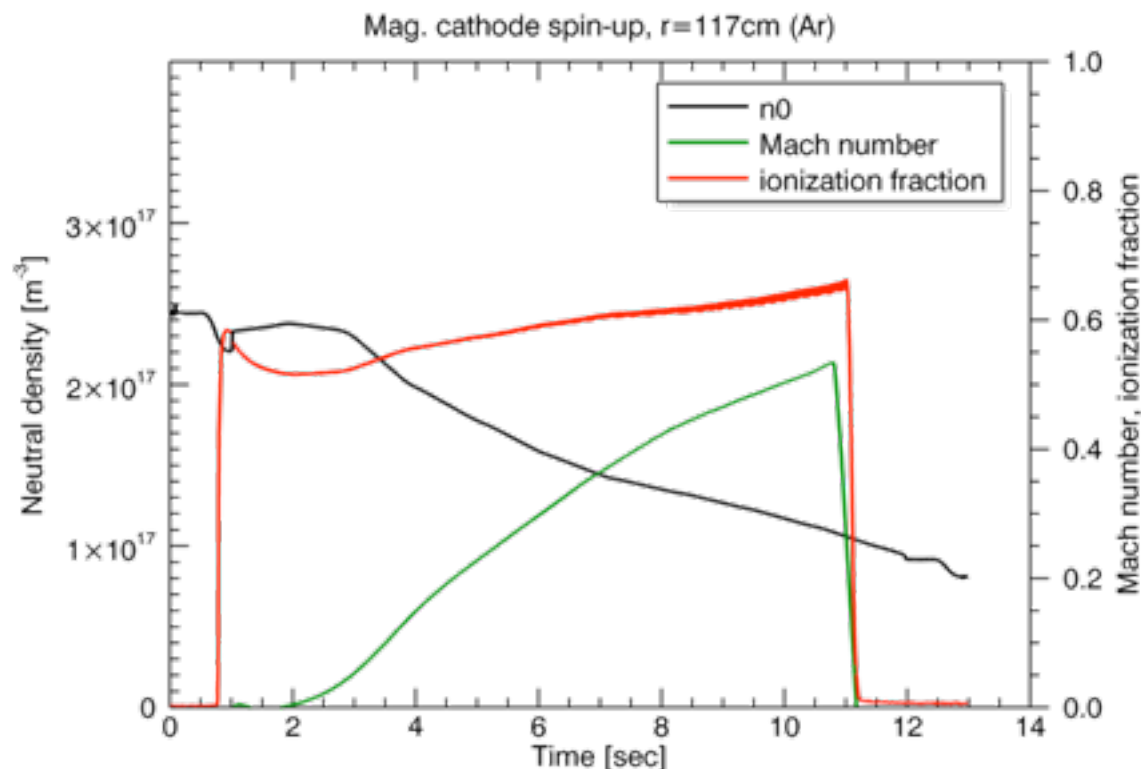
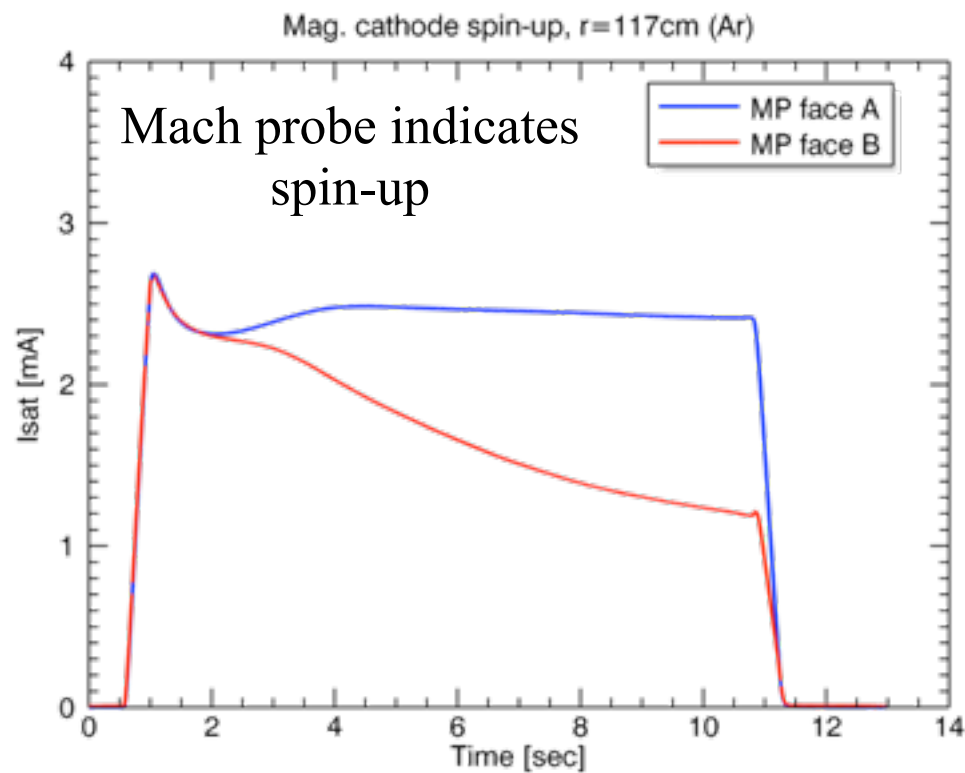
Argon



Helium



Flow observed with cathodes withdrawn into magnetized region



- Max velocity measured near anode location.
- Plasma spins up on neutral pump-out timescale.

$$mn \frac{du_\phi}{dt} = (\mathbf{J} \times \mathbf{B})_\phi - \frac{mn u_\phi}{\tau_{i0}} + mn \nu_{ii} [\nabla^2 \mathbf{u}]_\phi$$

Particle balance

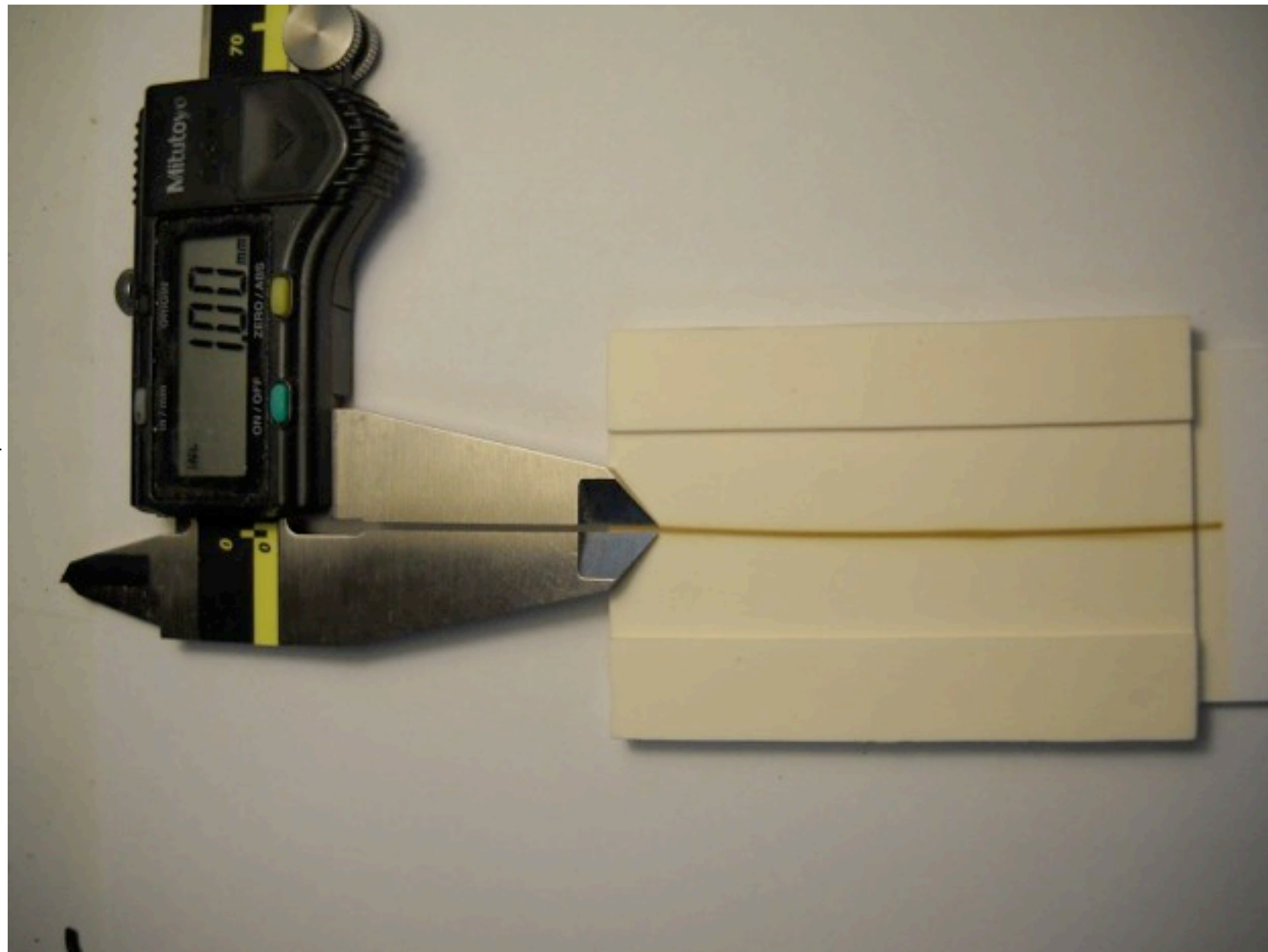
$$\frac{dn}{dt} = \frac{1}{2} \frac{c_s n A_{loss}}{V_{plasma}}$$

$$V_{plasma} = 10 \text{ m}^3$$

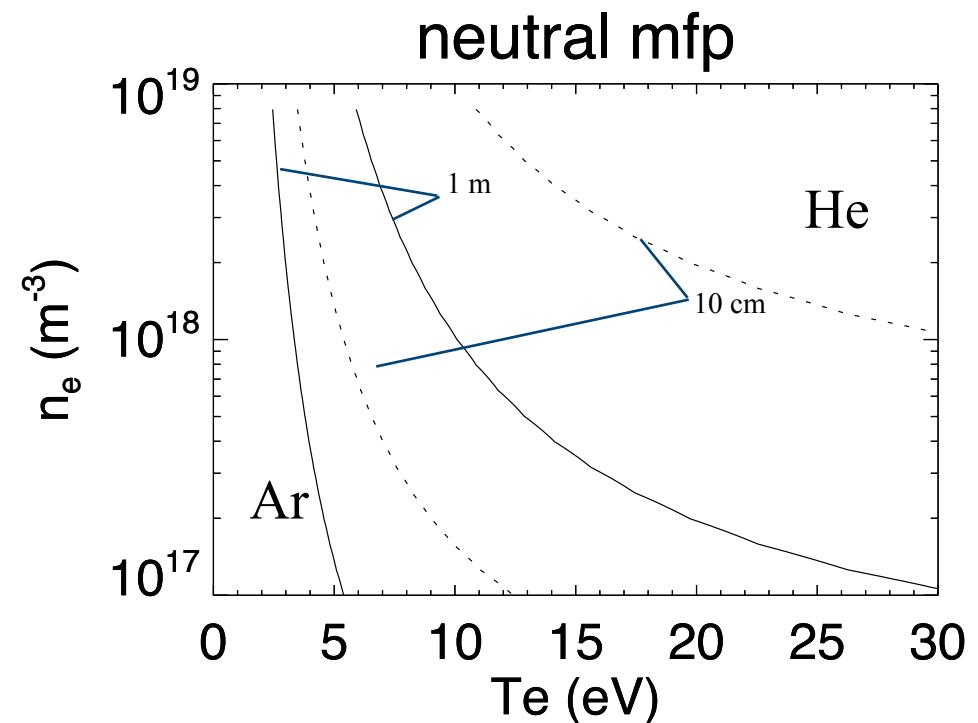
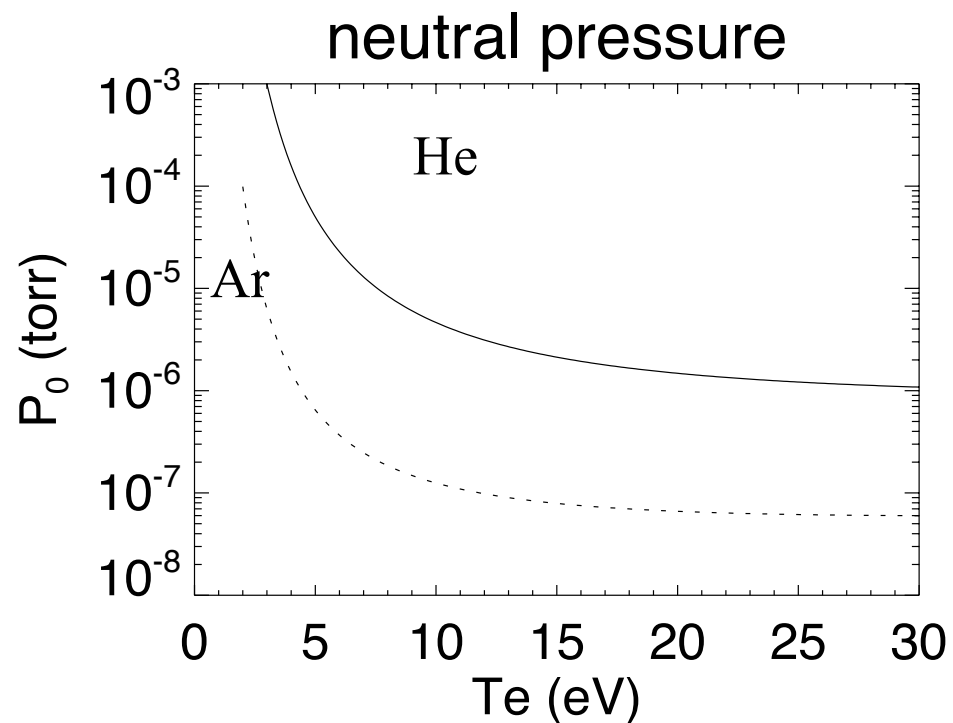
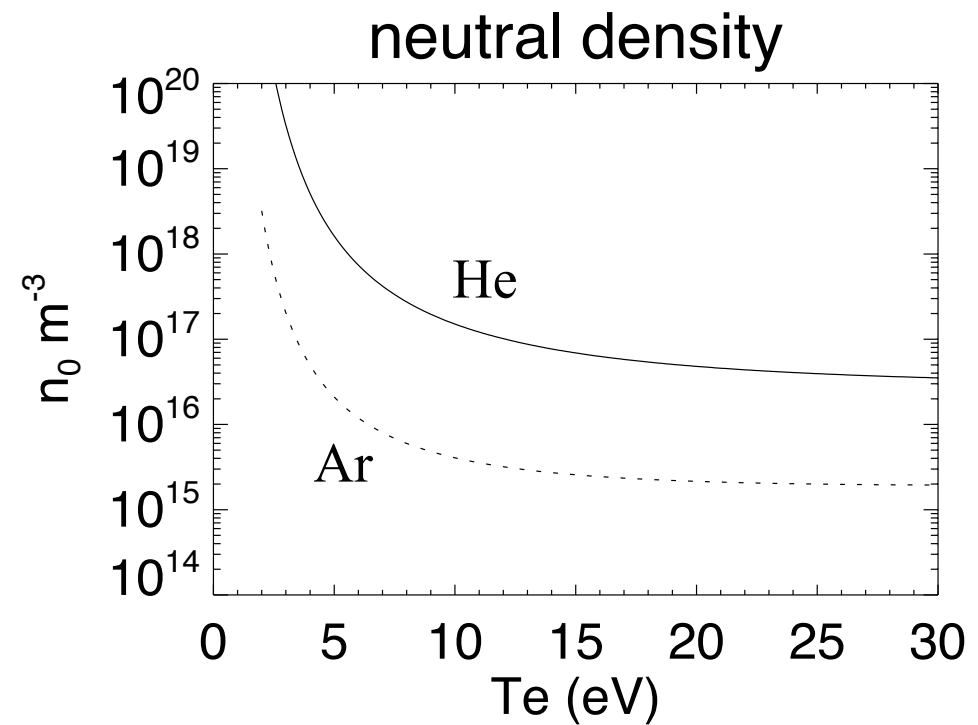
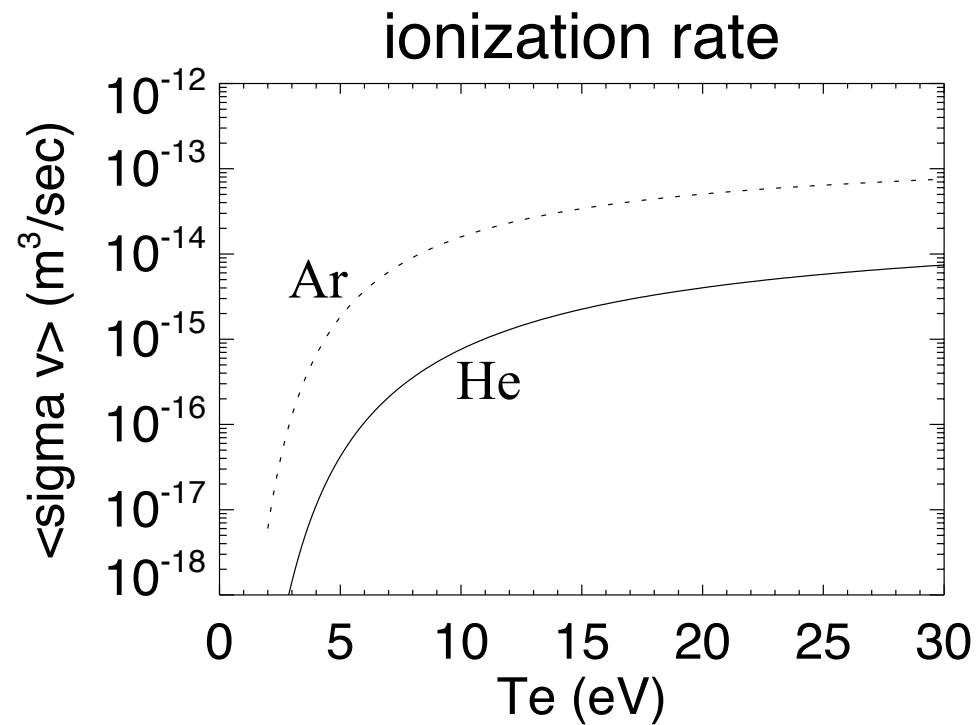
$$A_{loss} = W_{cusp} \times L_{cusp}$$

$$W_{cusp} = 4\sqrt{\rho_e \rho_i} \sim 1 \text{ mm}$$

$$L_{cusp} = 220 \text{ m}$$

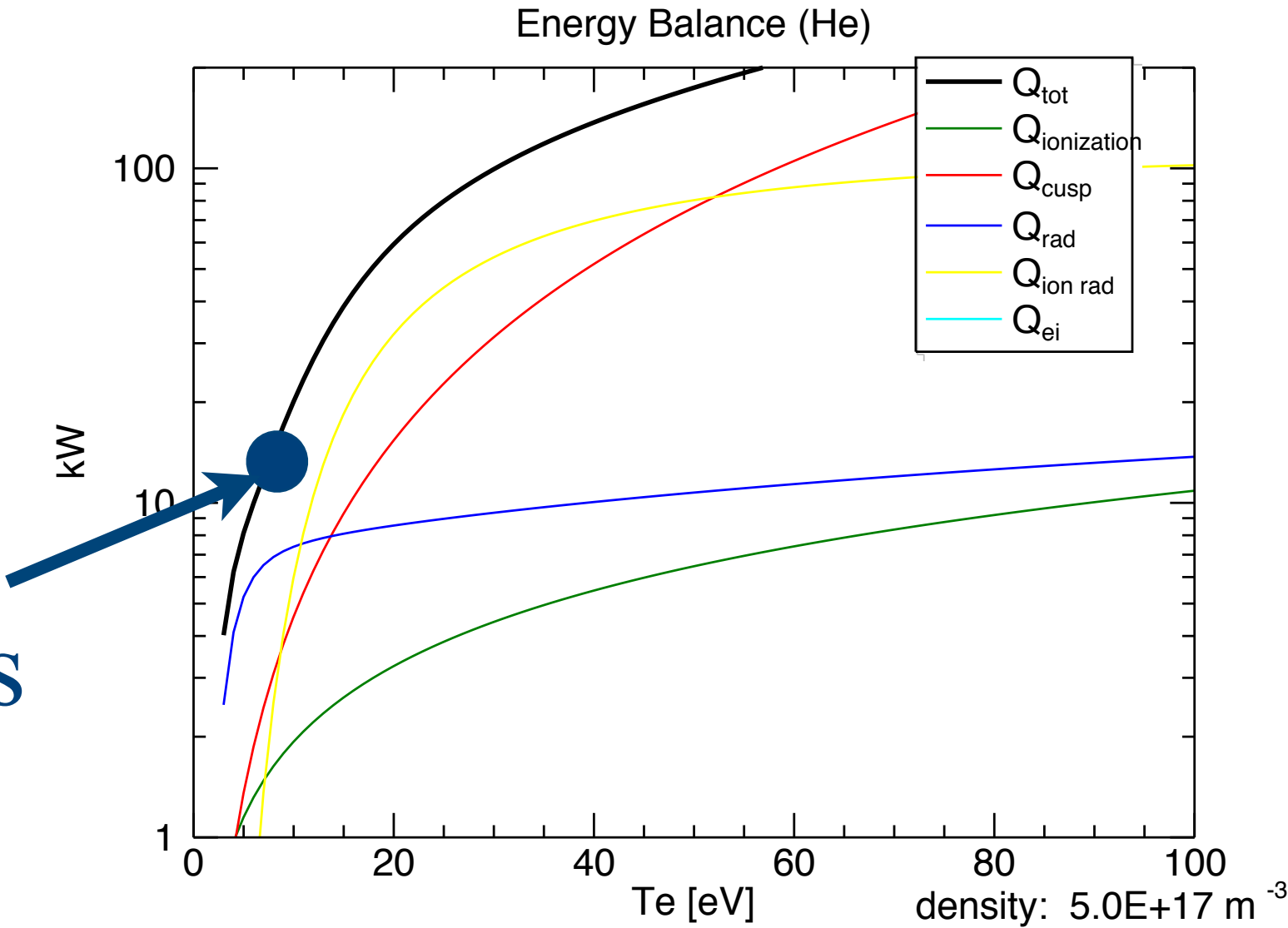


particle balance is $\frac{1}{2}n_e c_s A_{loss} = \langle \sigma_i v \rangle n_e n_0 V_{plasma}$
 so $n_0 = \frac{1}{2} \frac{c_s}{\langle \sigma_i v \rangle} \frac{A_{loss}}{V_{plasma}}$ is just a function of T_e



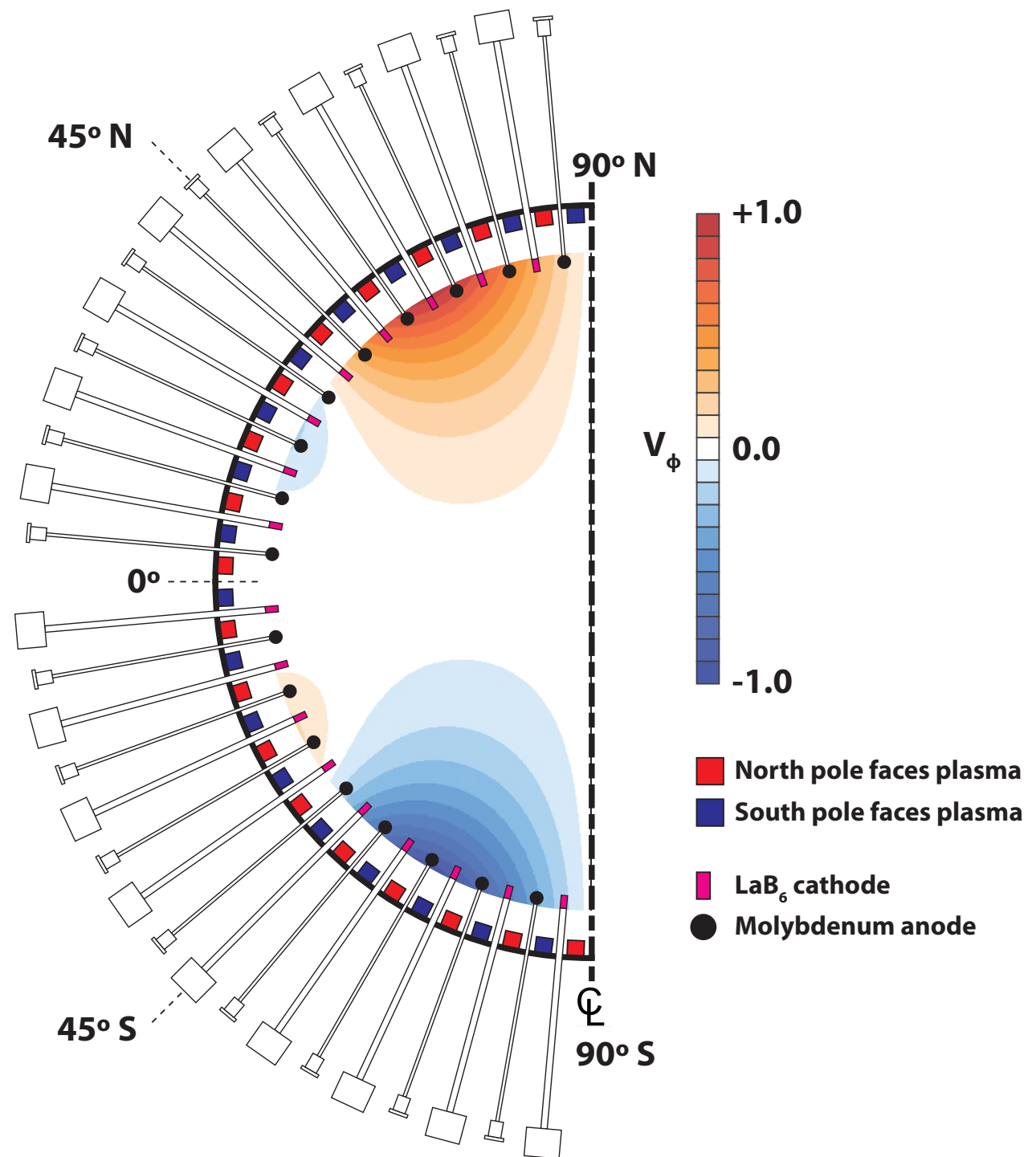
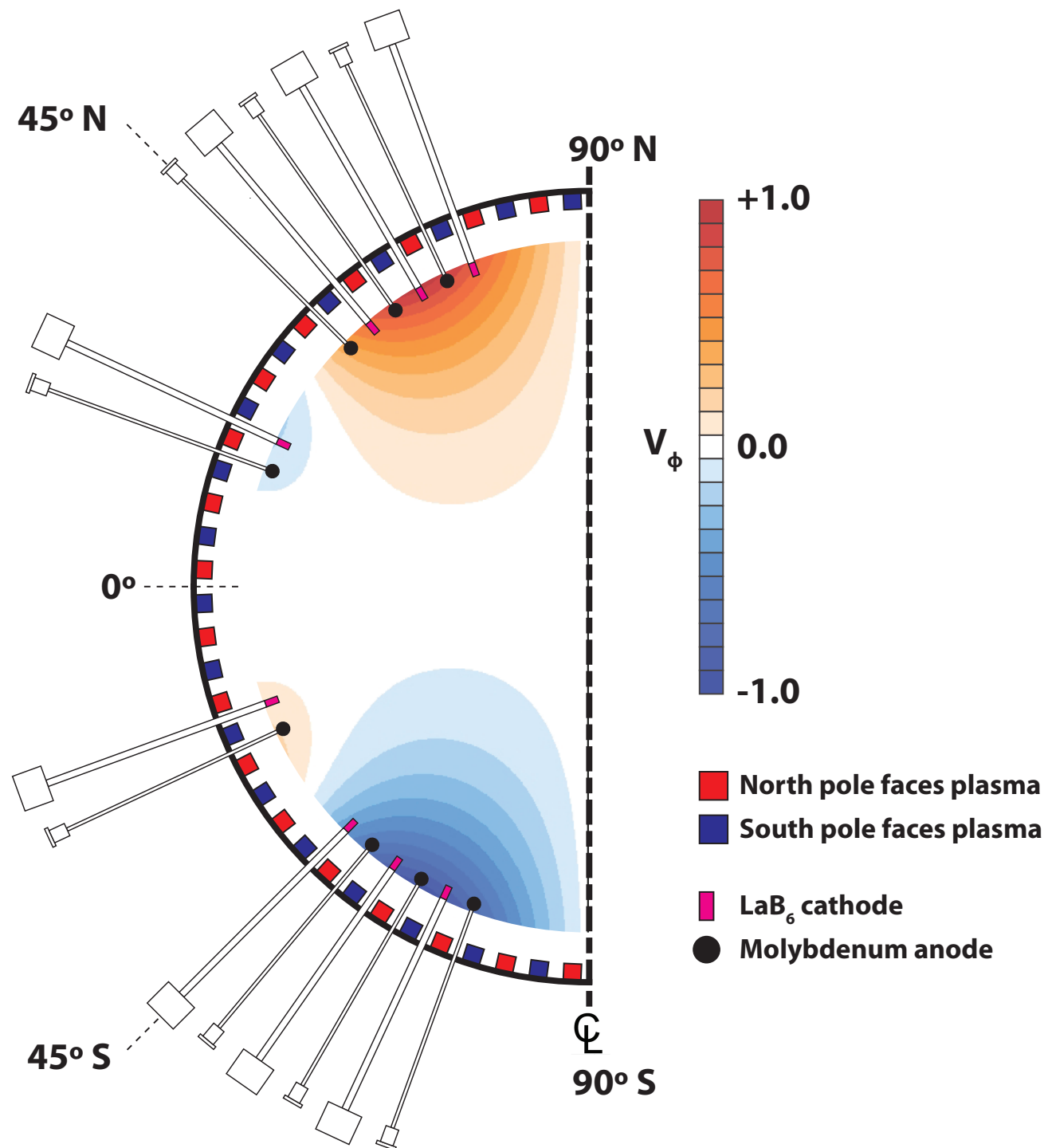
Power Balance:

First Results



Loss Mechanism	Expression [Energy/Time]
Ion losses at cusps	$\Gamma A_{\text{cusp}}(e\Delta V + \frac{3}{2}T_i)$
Electron losses at cusps	$\Gamma A_{\text{cusp}}(\frac{3}{2}T_e)$
Replacement ionization	$Q_{\text{ioniz}} \equiv \Gamma A_{\text{cusp}}E_{\text{ioniz}}$
Neutral radiation	$M(T_e)Q_{\text{ioniz}}$
Charge-exchange collisions	$\frac{3}{2}nn_0\langle\sigma_{cx}v\rangle e(T_i - T_0)Vol$
Ion radiation	$n^2R^i(T_e)Vol$

Soon (May, 2013) ... eventually (Fall, 2013)



Additional Heating from Microwaves



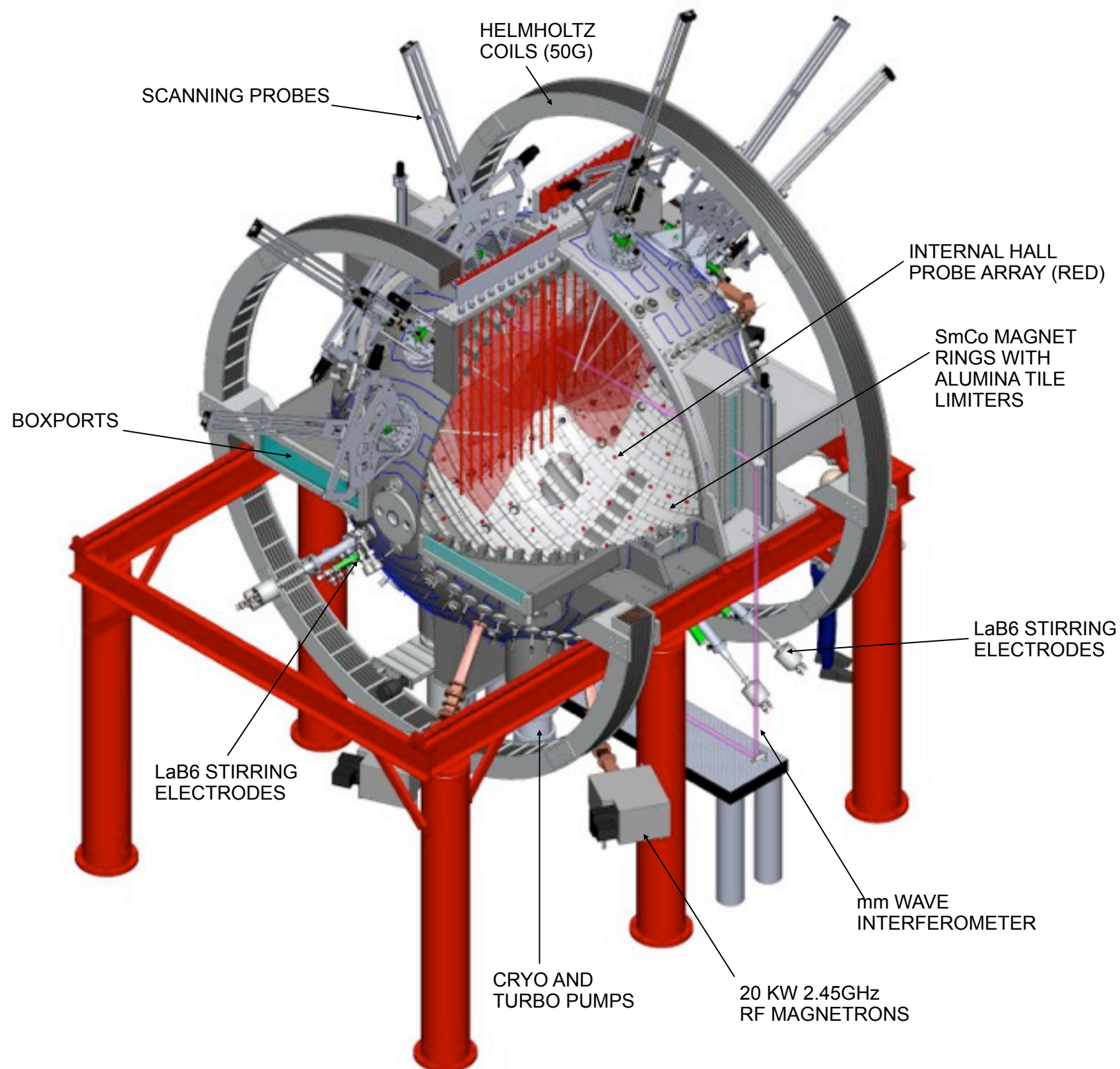
- Magnetrons produce Microwave radiation at 2.45 GHz (like at home)
- May, 2013



- Total heating power:

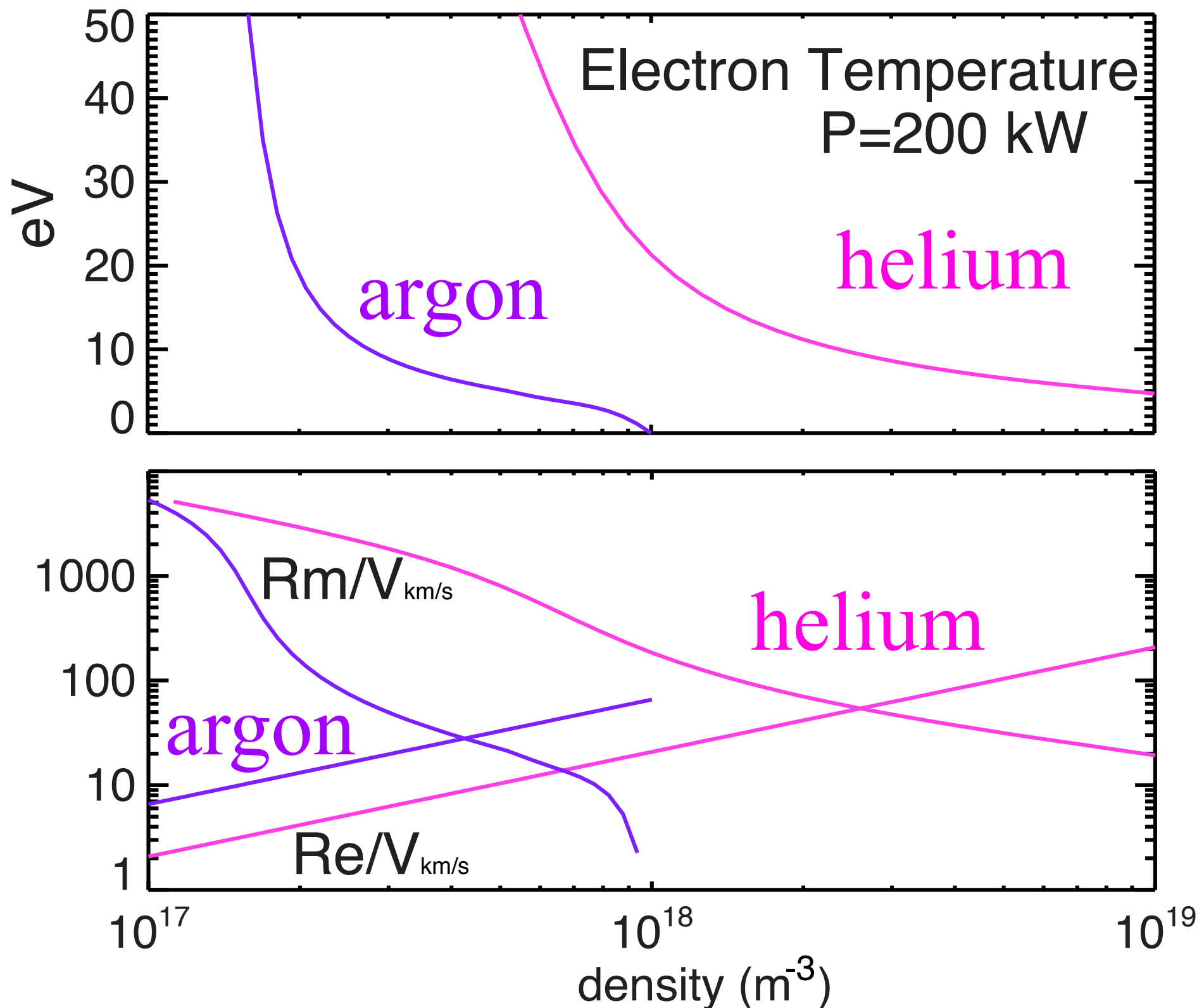
$$5 \times 20 \text{ kW} = \underline{100 \text{ kW}}$$

Diagnostics

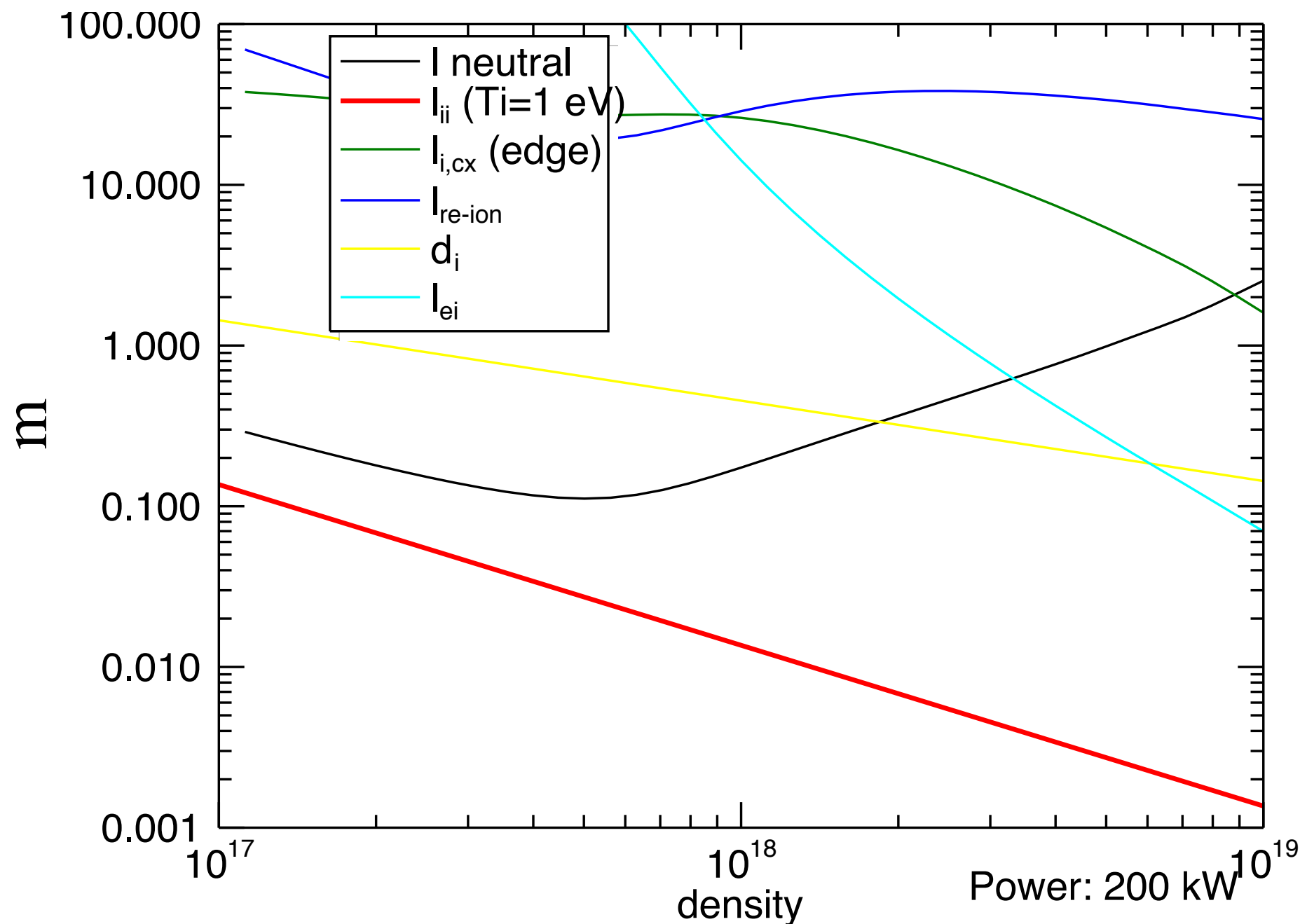


- 3-D Probes (n , V , B)
- 1 mm inteferometer/ polarimeter (n , B)
- line-ratio spectroscopy (Te , n)
- Fabry-Perot Ion Doppler Spectroscopy (Ti , V)
- Laser induced Fluorescence (Ti , V)
- Neutral Particle Analyzer
- Fast imaging camera
- edge magnetic array $B(a, \theta)$
- edge Mach Probe Array $V_{\phi}(a, \theta)$

Operational space set by power, density, and ion species



At low density, low collisionality plasma effects will play role



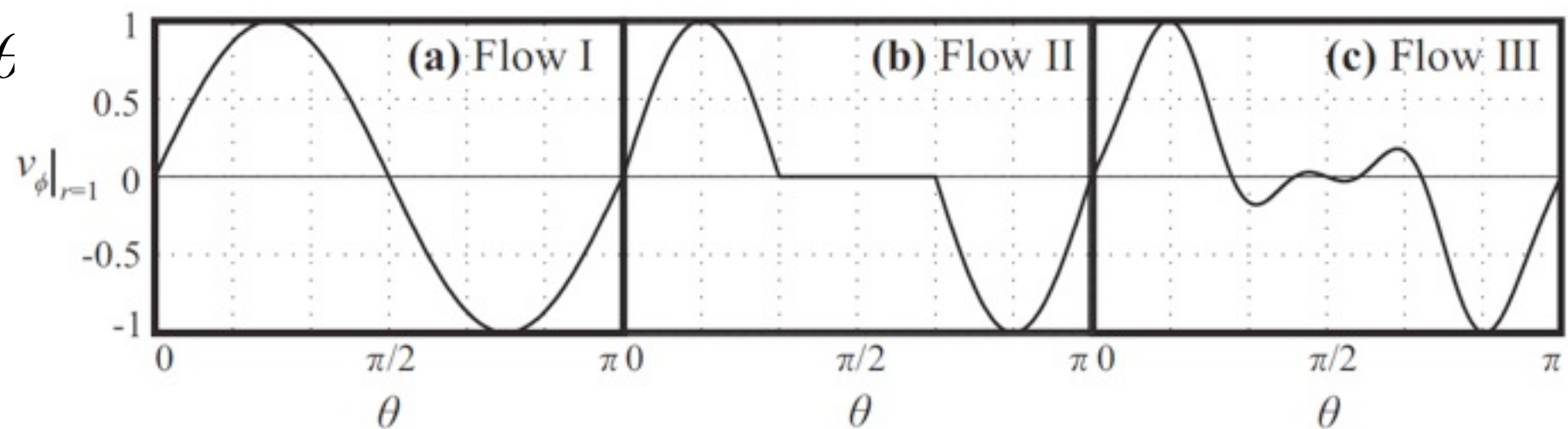
$$\rho_i = 1.02 \frac{\sqrt{\mu T_{i,\text{eV}}}}{B_g} \text{ m}$$

$$\rho_e = 0.023 \frac{\sqrt{T_{e,\text{eV}}}}{B_g} \text{ m}$$

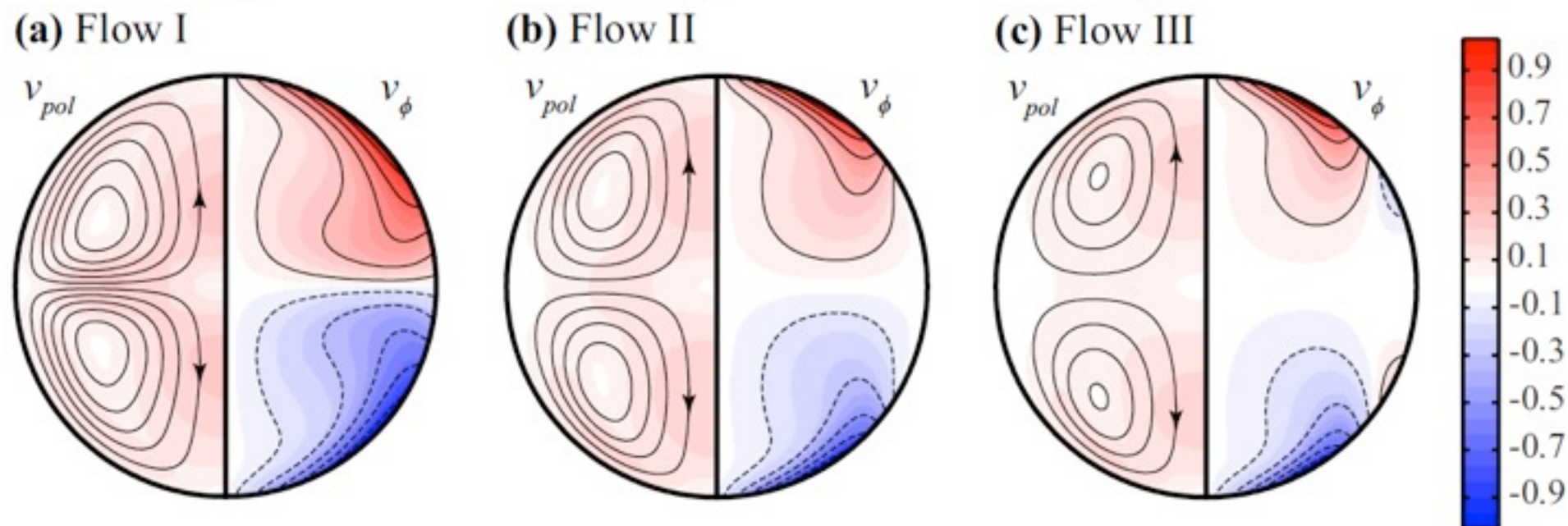
Boundary-driven flows

$$(\mathbf{v} \cdot \nabla) \mathbf{v} + \nabla p - \frac{1}{Re} \nabla^2 \mathbf{v} = 0, \quad \nabla \cdot \mathbf{v} = 0$$

Driving velocity $v_\phi(t)$
at the boundary



Flow structure ($Re=150$)

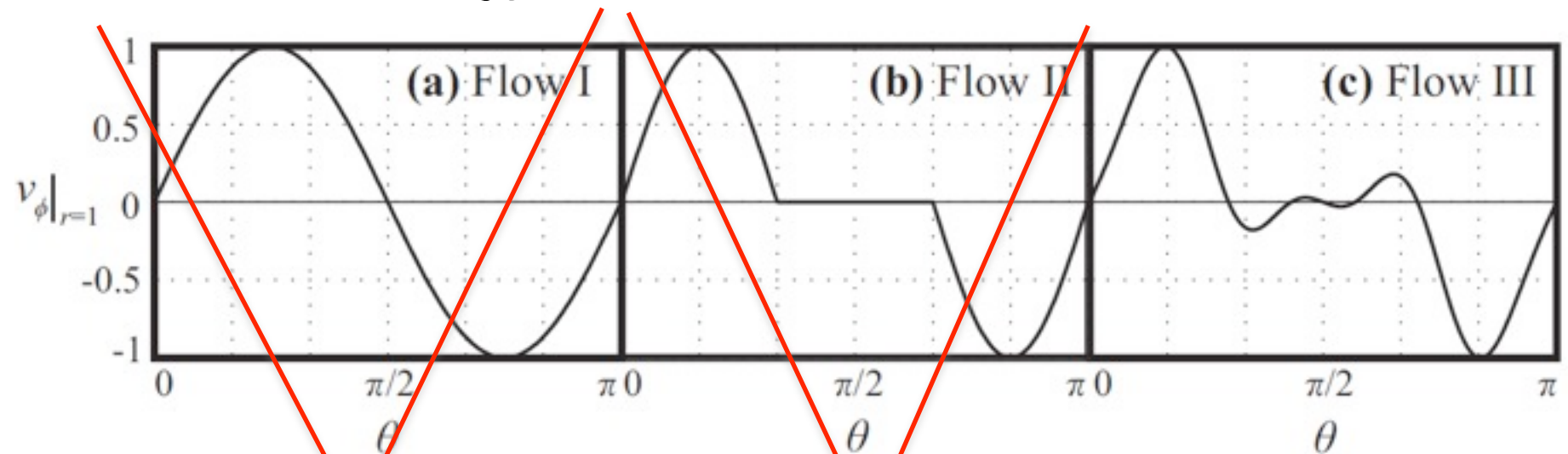


- Which flow can result in dynamo?

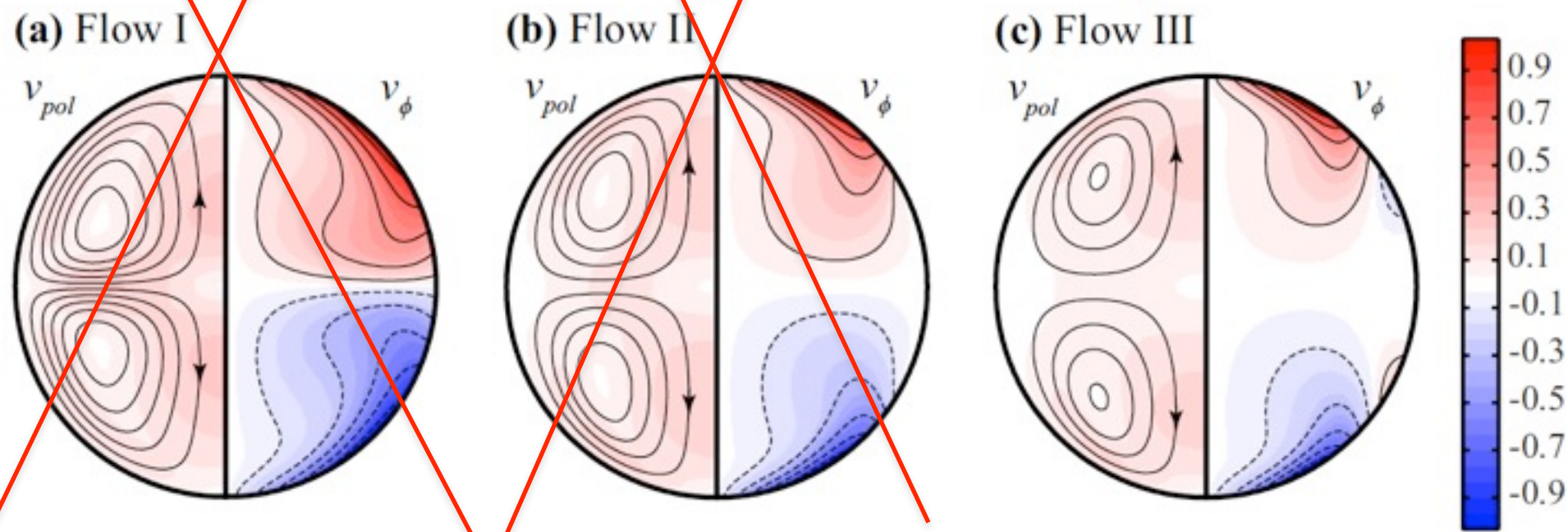
Boundary-driven flows

$$(\mathbf{v} \cdot \nabla) \mathbf{v} + \nabla p - \frac{1}{Re} \nabla^2 \mathbf{v} = 0, \quad \nabla \cdot \mathbf{v} = 0$$

$v_\phi(\theta)$
at the boundary



Flow structure ($Re=150$)

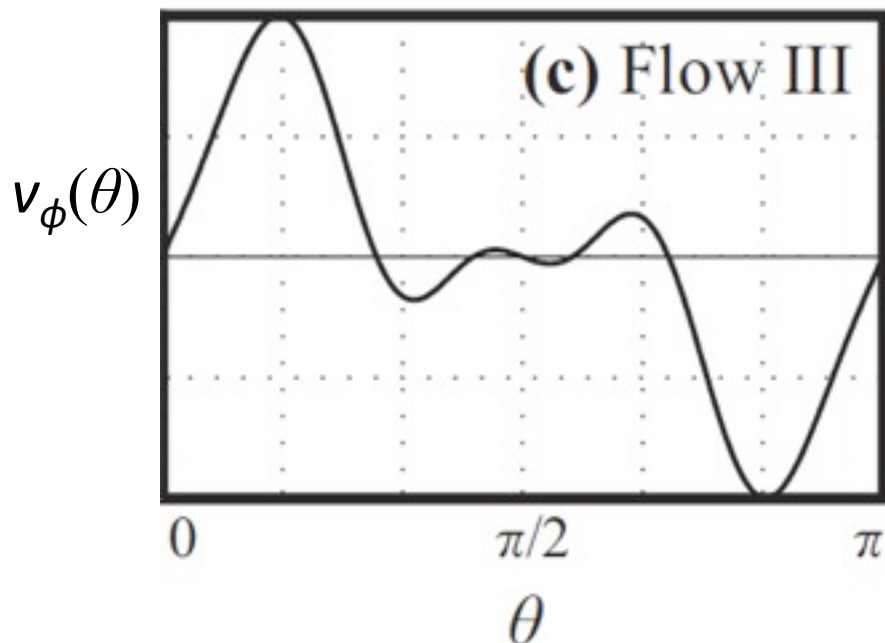


Dynamo!

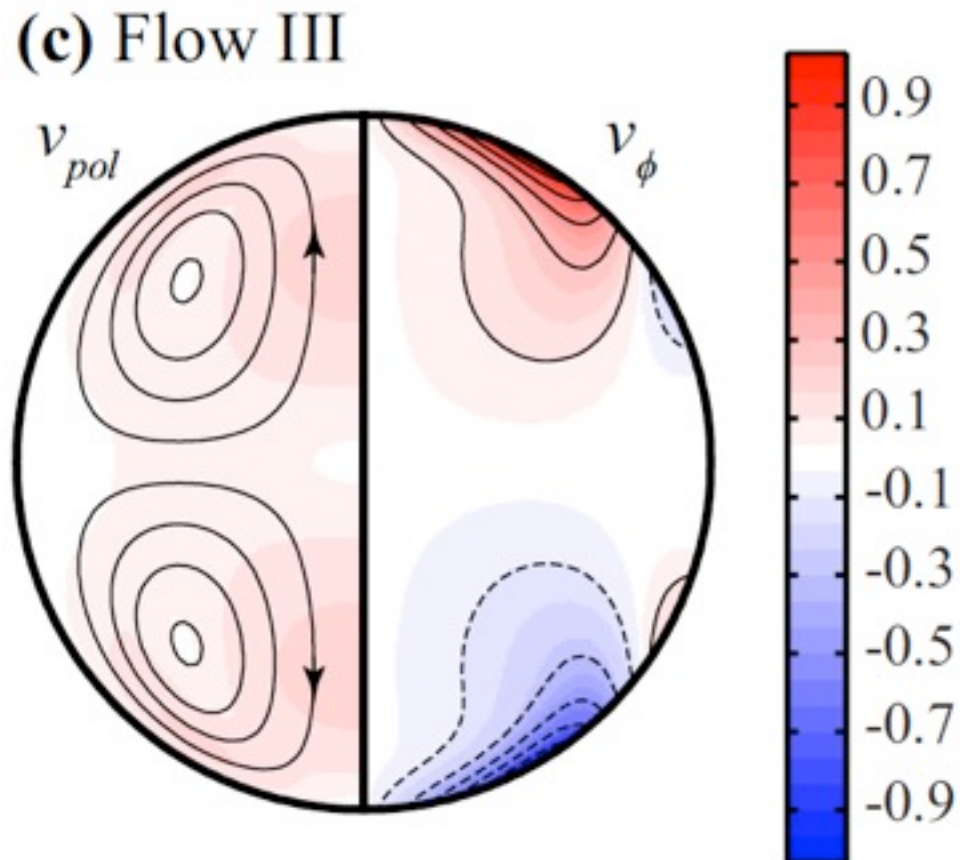
- Which flow can result in dynamo?
- Only flow III [from Spence et al., 2009]!

Dynamo in flow III

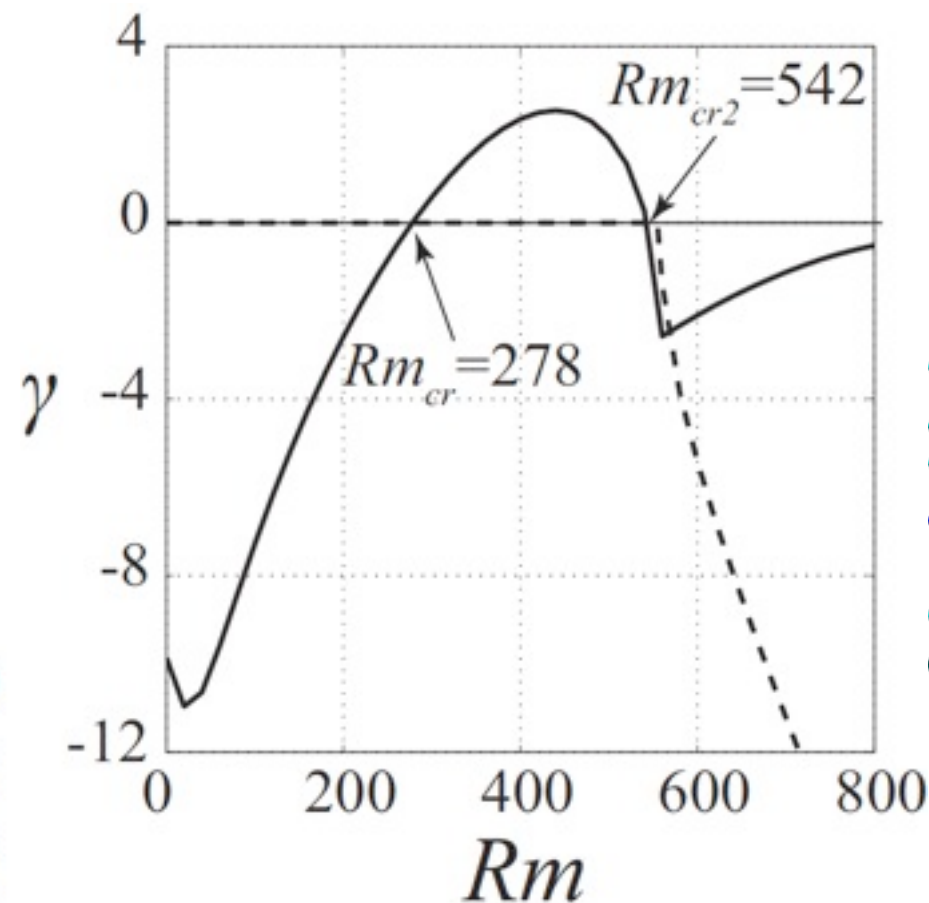
Boundary drive



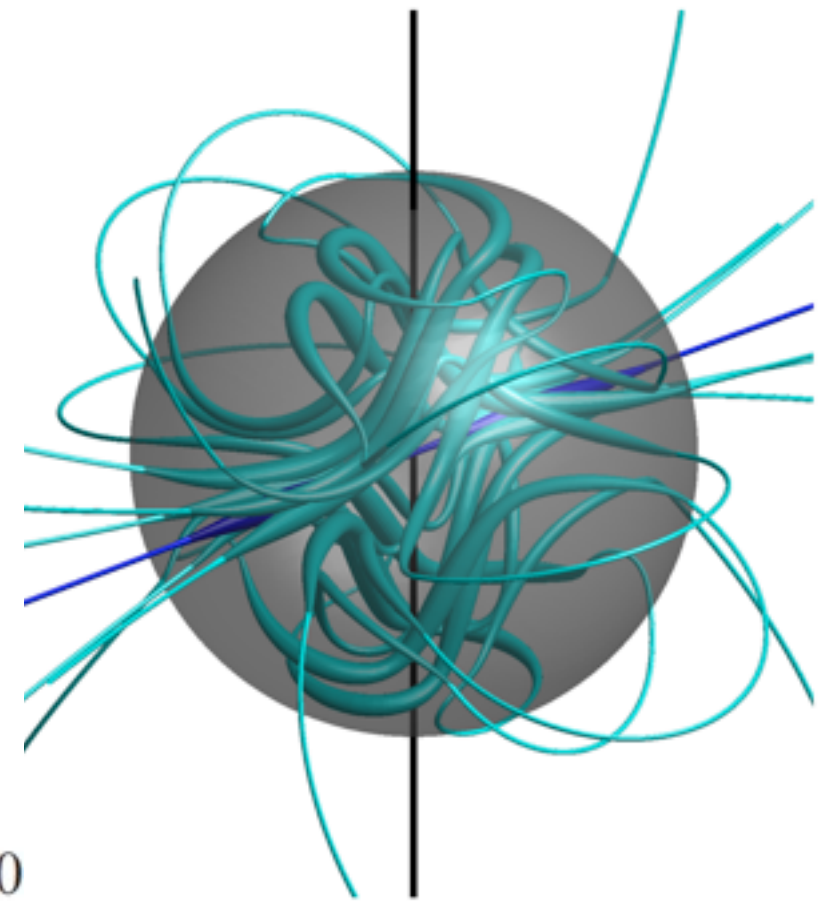
Flow structure, $Re=150$



Growth rate vs. Rm
(mode $m=1$)



Dynamo field, $Rm=400$

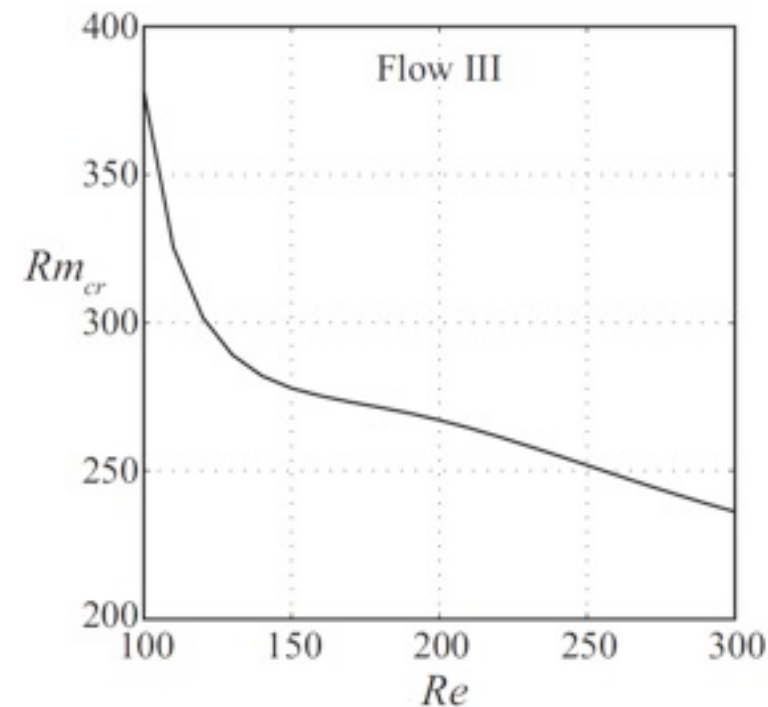
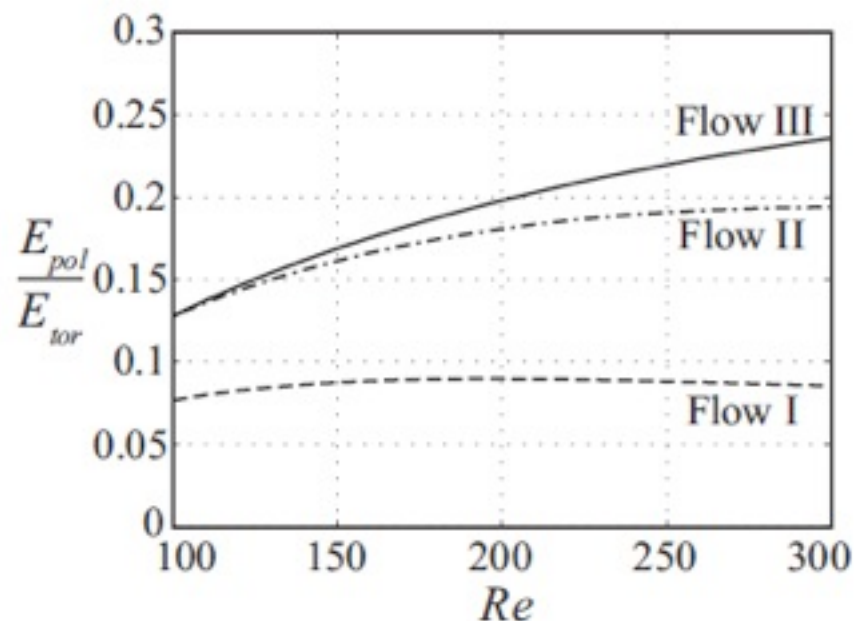


- This is achievable in MPDX with argon plasma, $V_0=5$ km/s, $n_0=10^{18}$ m⁻³, $T_e=10$ eV, $T_i=1$ eV
- Growth rate γ is normalized by resistive time

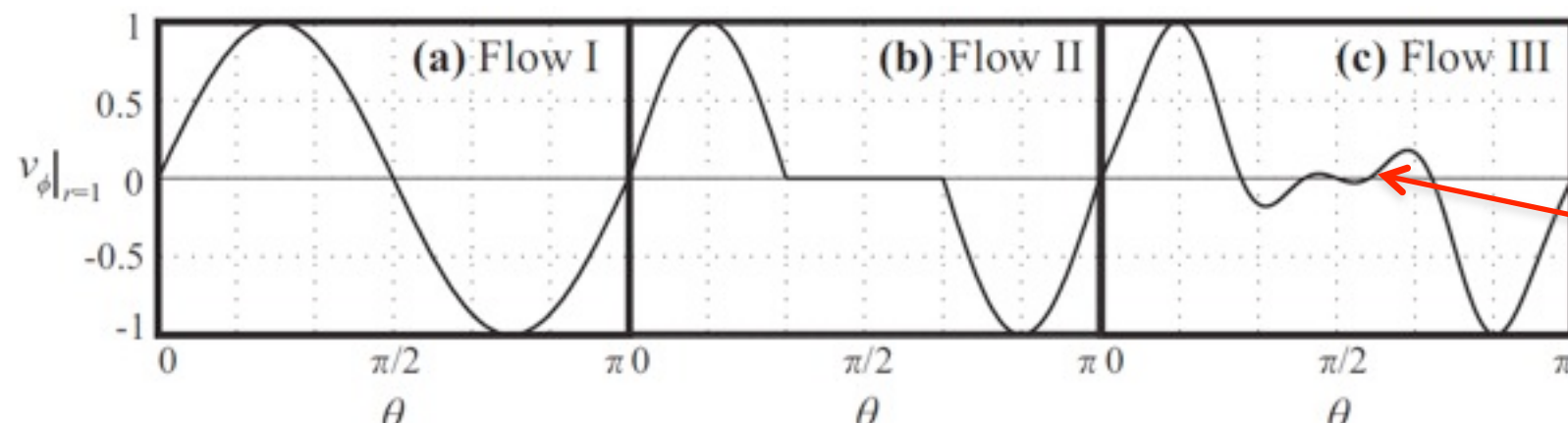
$$\tau_\sigma = \frac{R^0}{\eta} \approx 0.13 \text{ sec}$$

What determines dynamo onset?

- Pure toroidal flow cannot support dynamo [Elsasser, 1946]
- Dynamo onset depends on poloidal flow



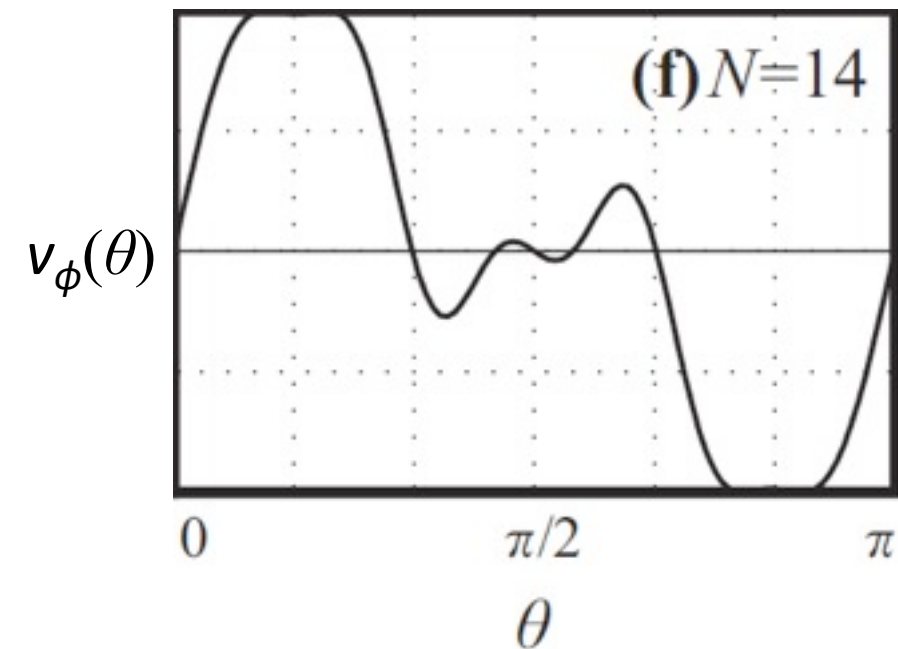
- Poloidal flow is necessary for dynamo, but not sufficient
- Details of flow structure play important role:



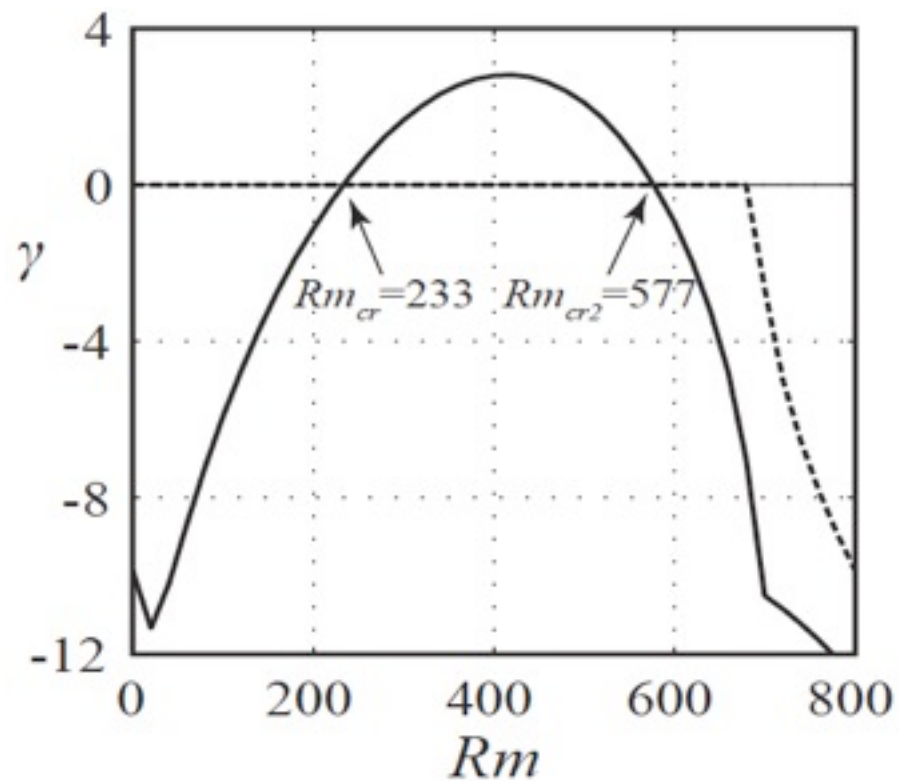
Flow reversals

Optimized flow

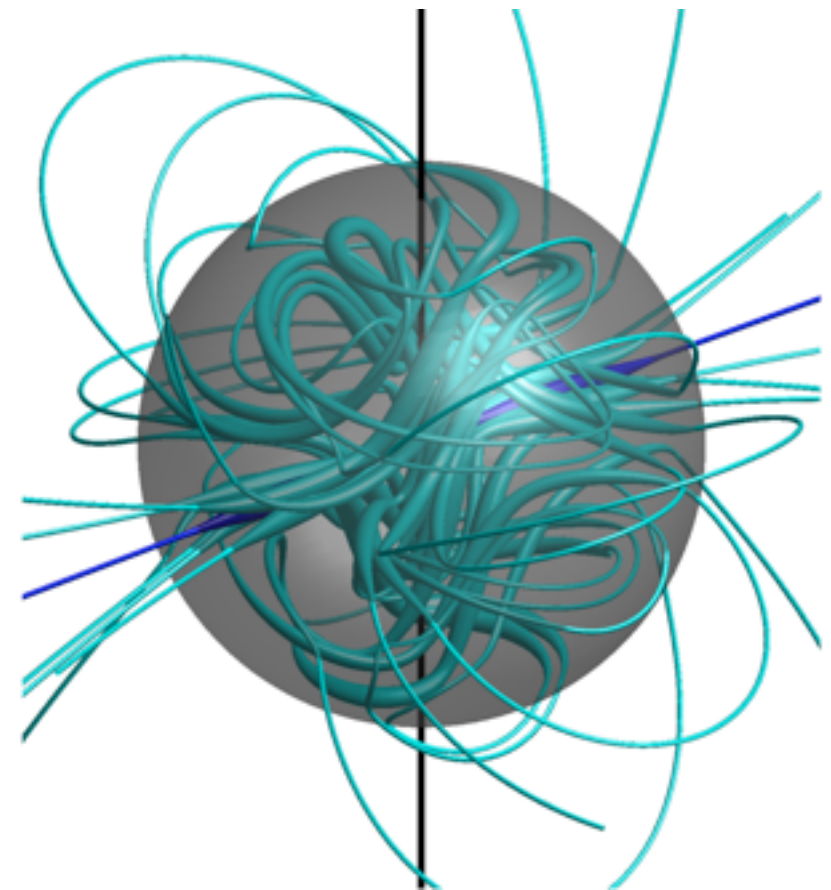
Boundary drive



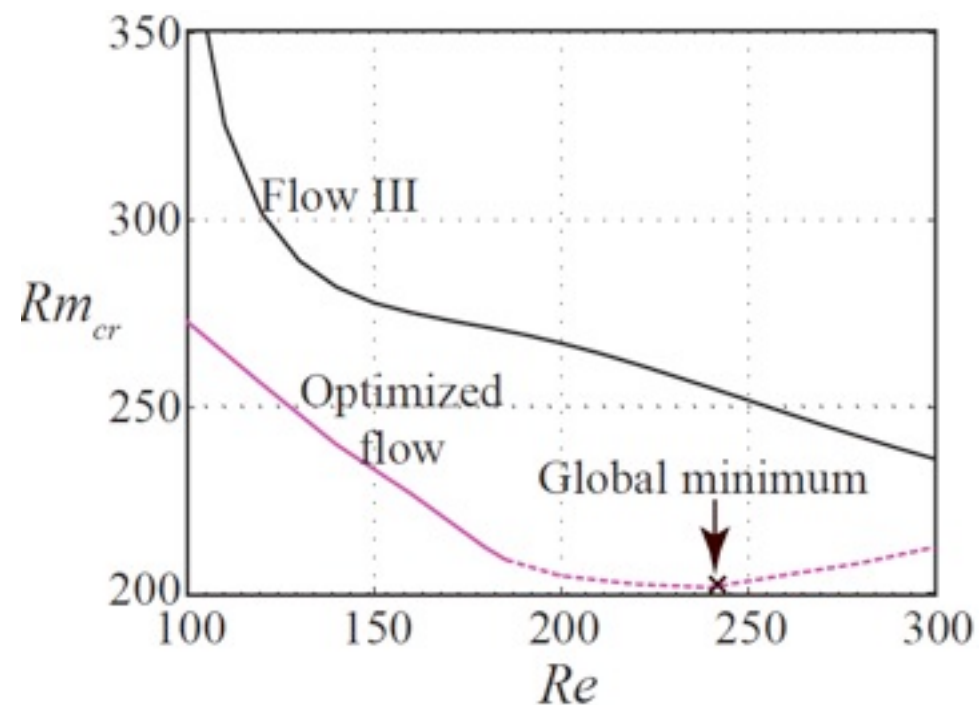
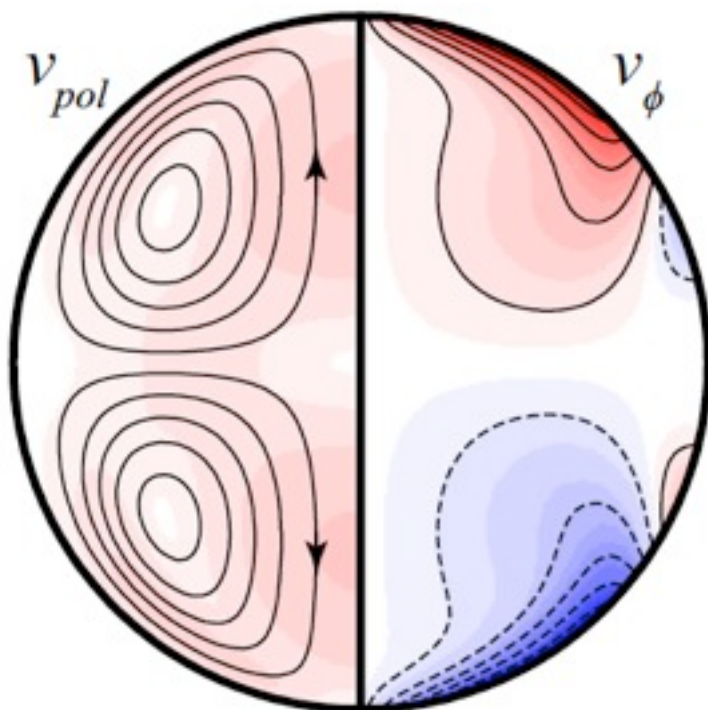
Growth rate vs. Rm



Dynamo field, $Rm=400$



Flow structure, $Re=150$



- The lowest $Rm_{cr} \approx 200$ achieved at $Re \approx 240$

Spherical von Karman flow

- Idea: create a hydrodynamically unstable flow, so non-axisymmetric modes are induced
- Solve full MHD system:

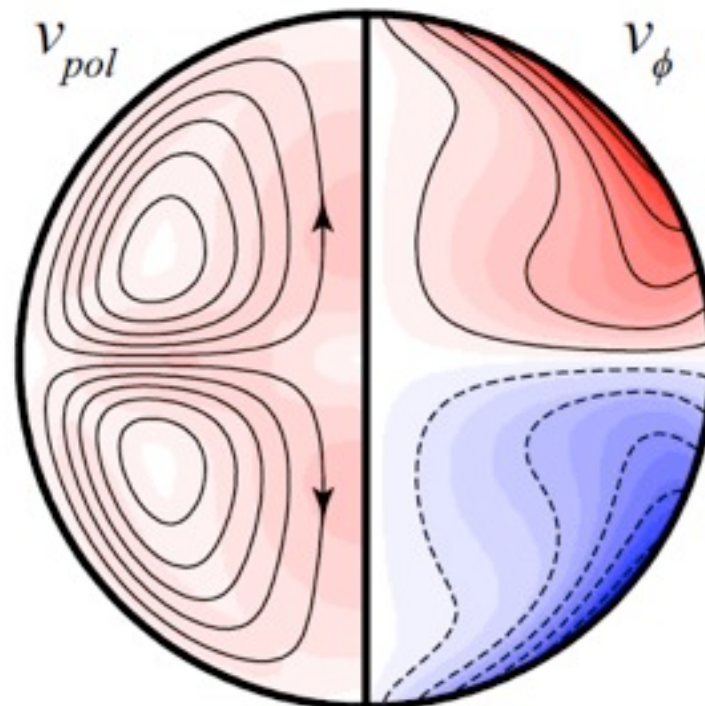
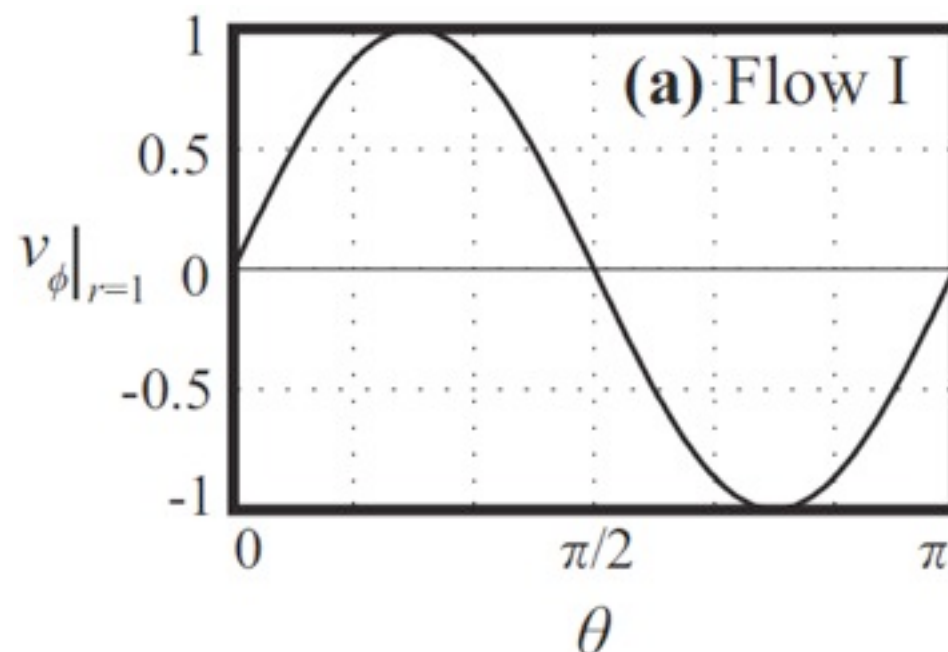
$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + \nabla p - \frac{1}{Re} \nabla^2 \mathbf{v} = 0, \quad \nabla \cdot \mathbf{v} = 0, \quad Re = \frac{R_0 V_0}{\nu}$$

Boundary drive

$$v_\phi(\theta) = \sin(2\theta)$$

Axisymmetric flow structure,

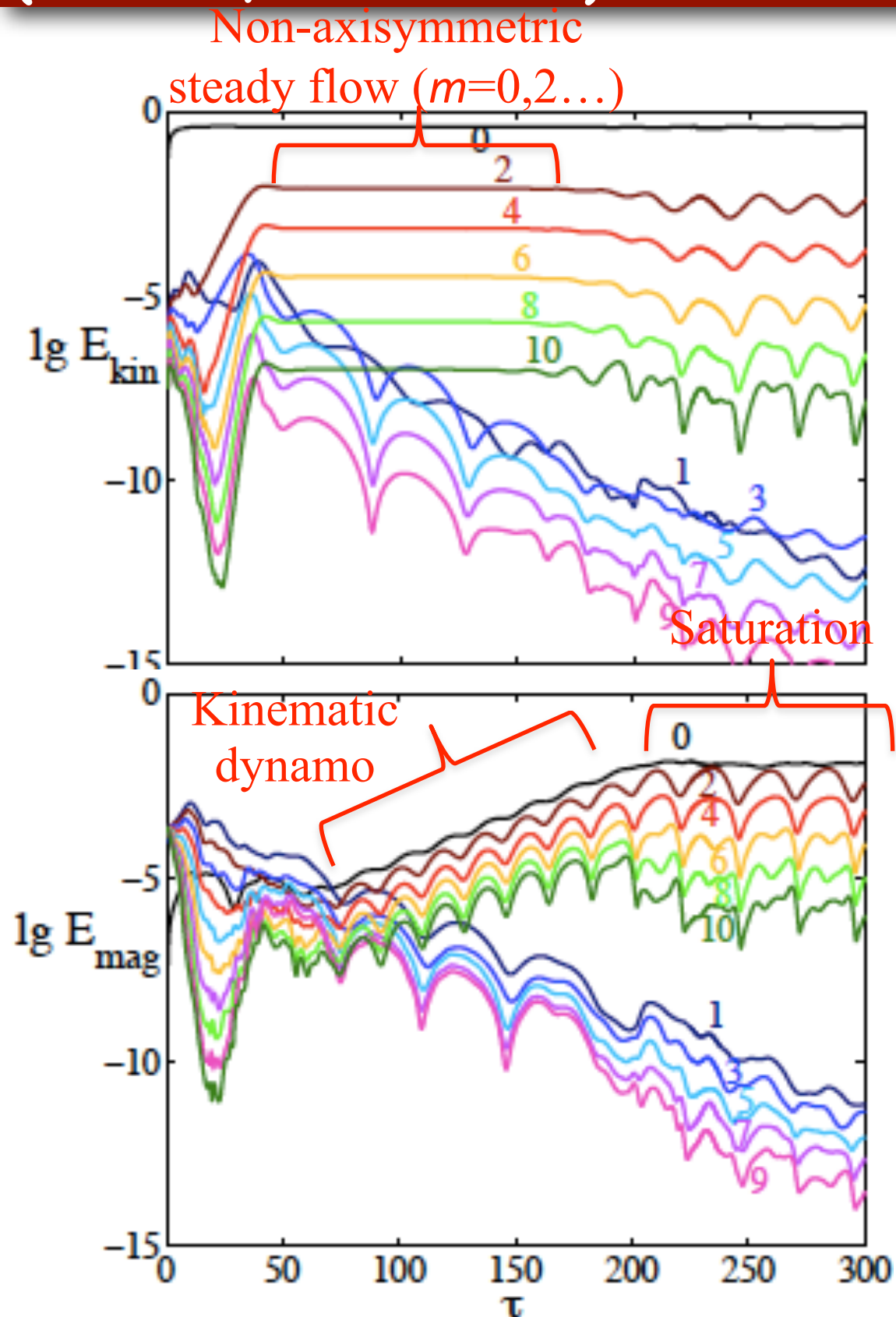
$$Re = 150$$



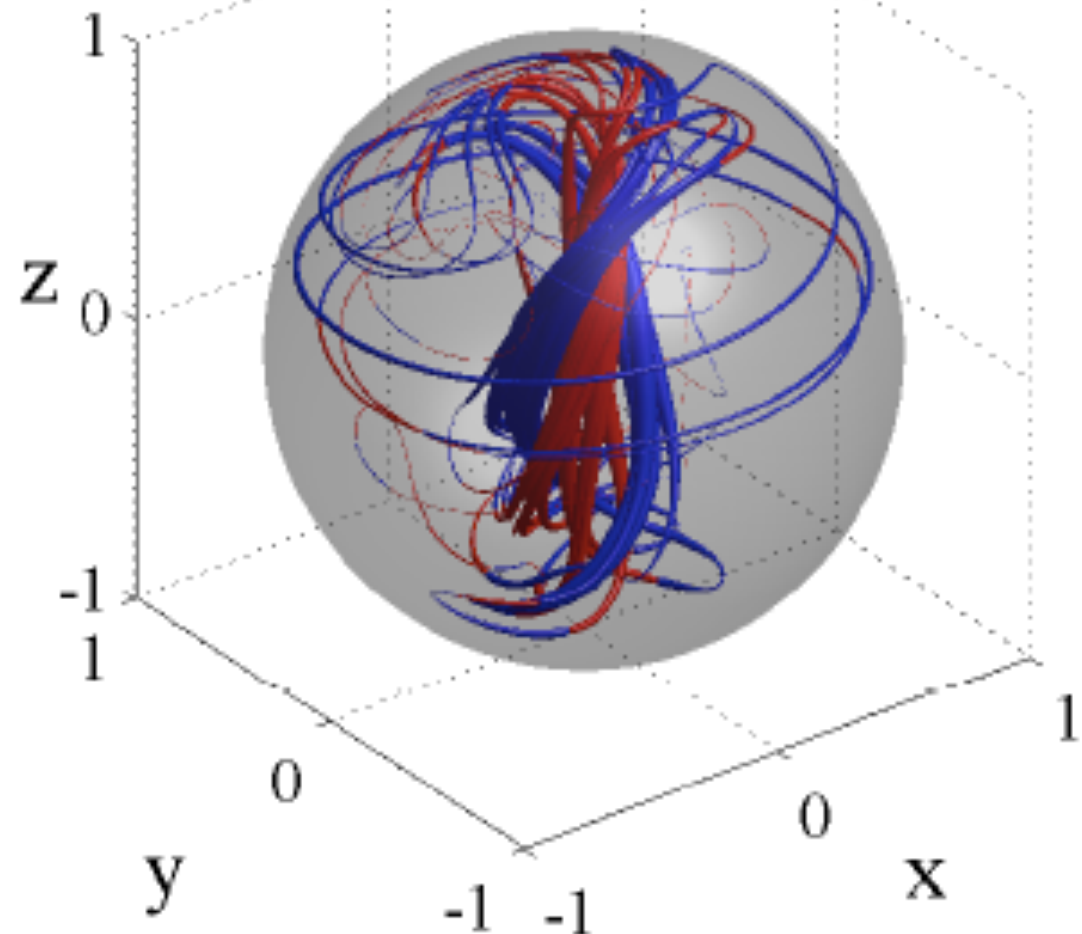
- This flow is unstable hydrodynamically with respect to mode $m=2$ when $Re > 115$

Dynamo in non-axisymmetric flow

($Re=150$, $Rm=2000$)

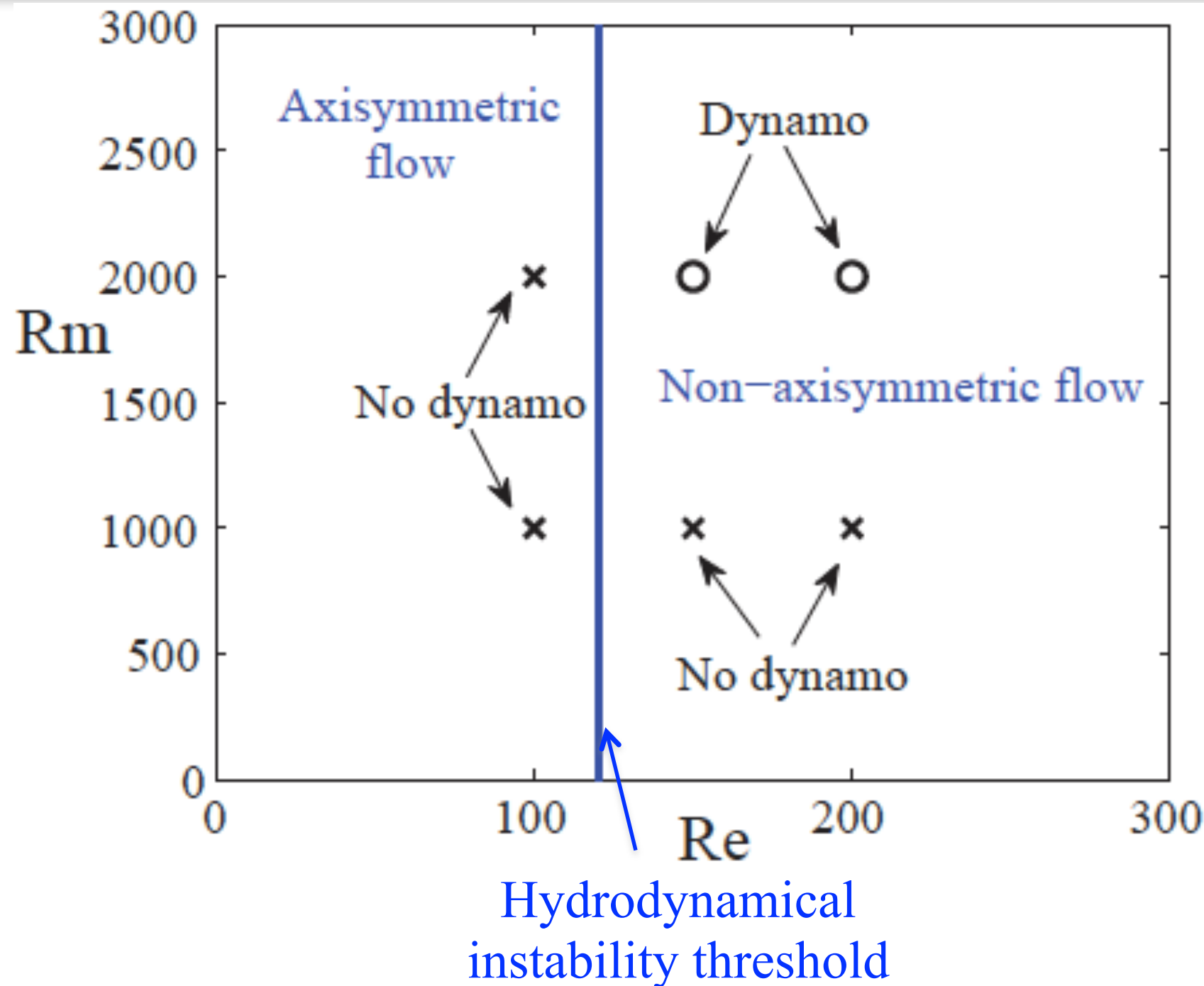


Magnetic field lines at $t=237.9$



- Saturated state is time-periodic
- Both flow and field have only even harmonics with $m=0, 2, 4, \dots$
- Regime corresponds to argon plasma with $V_0=10$ km/s, $n_0=10^{18} \text{ m}^{-3}$, $T_e=18 \text{ eV}$, $T_i=1.3 \text{ eV}$

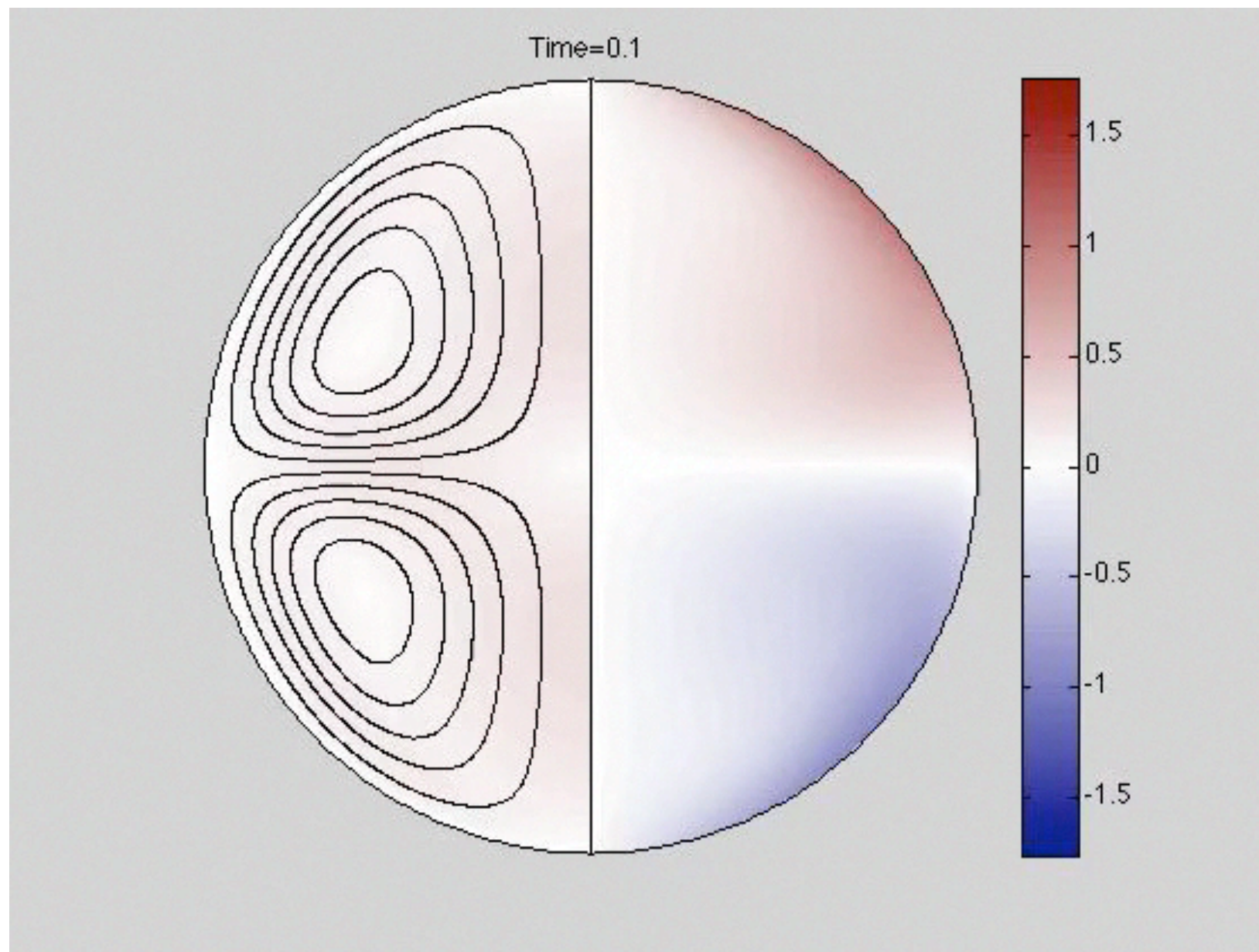
High Pm Domain of dynamo parameters (conducting wall, steady von Karman flows)



- Dynamo excited in non-axisymmetric von Karman flows ($Re > 115$) for sufficiently high Rm ($Rm > 1500$)

Galloway-Proctor flow in sphere driven by time dependent boundary flow

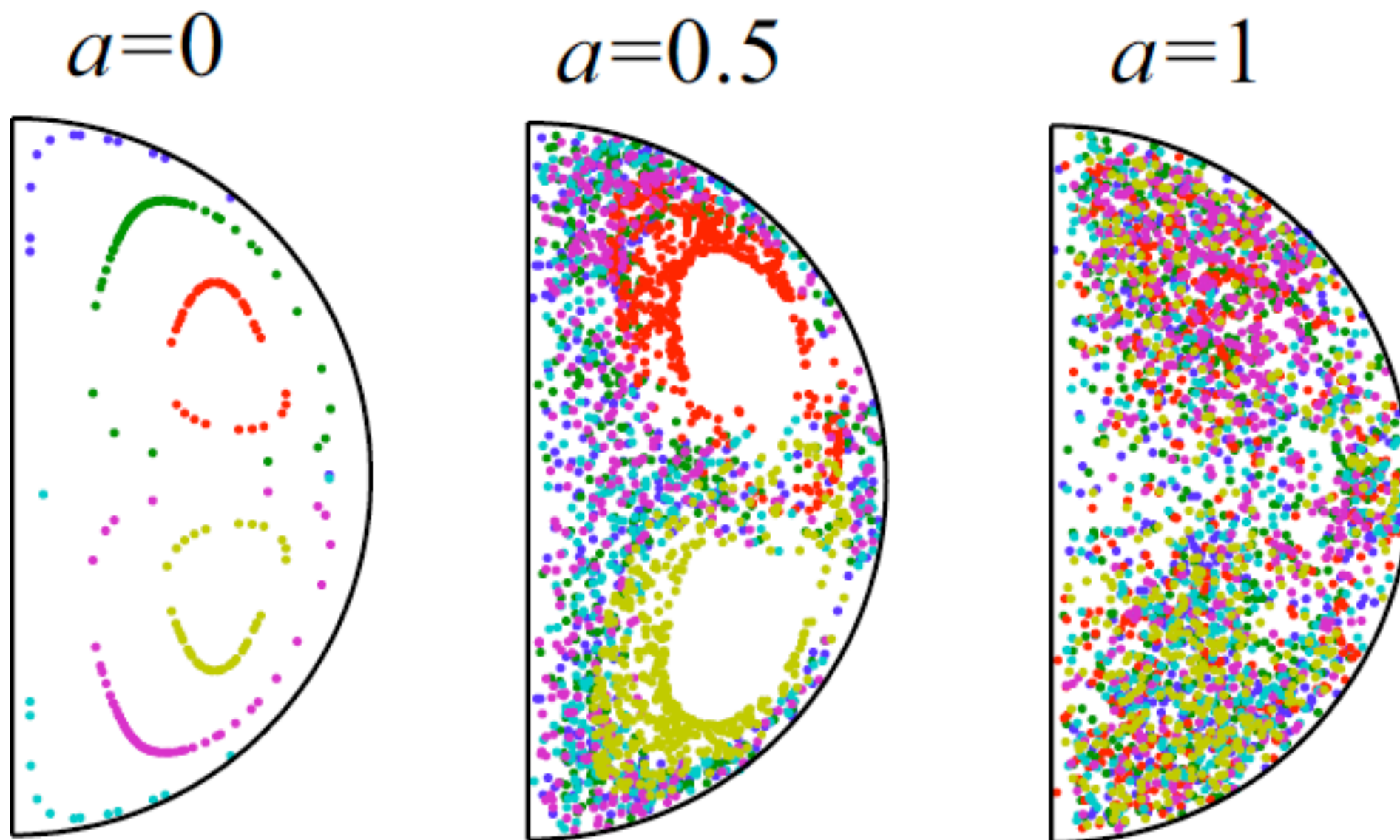
- Consider driving velocity: $V_\phi(\theta, t) = \sin 2\theta + a \sin \theta \sin \omega t$
 $Re=100, a=1, \omega=1$



solves Navier-Stokes:
$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + \nabla p = \frac{1}{Re} \nabla^2 \mathbf{v}$$

Chaos in time-periodic flow

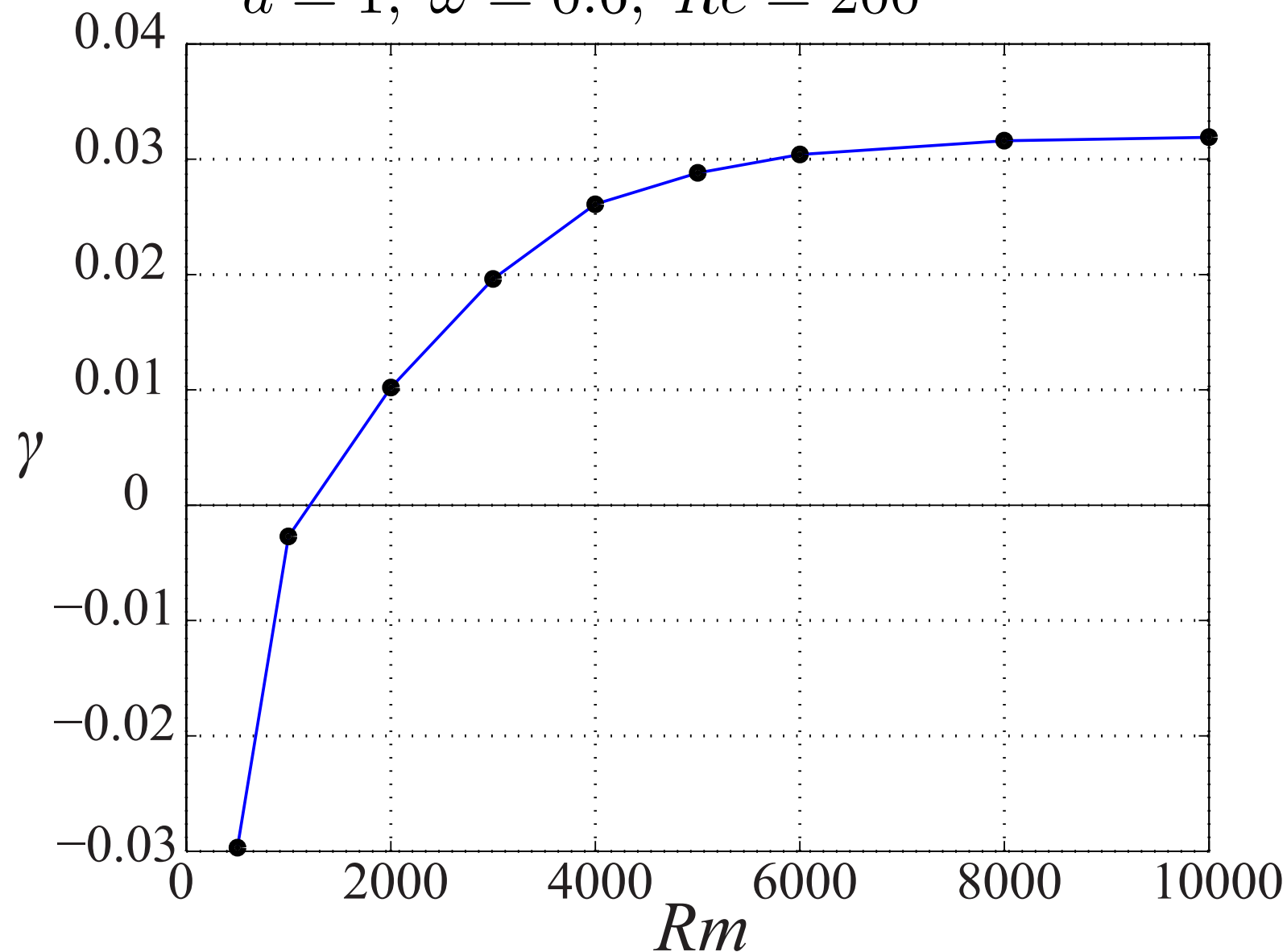
- Boundary: $Re=200$, $\omega=0.6$ $V_\phi(\theta, t) = \sin 2\theta + a \sin \theta \sin \omega t$



Growth rate vs. Rm shows dynamo is a fast dynamo

Drive: $V_\phi(\theta, t) = \sin 2\theta + a \sin \theta \sin \omega t$ Unit of time: $t_0 = R_0/V_0$

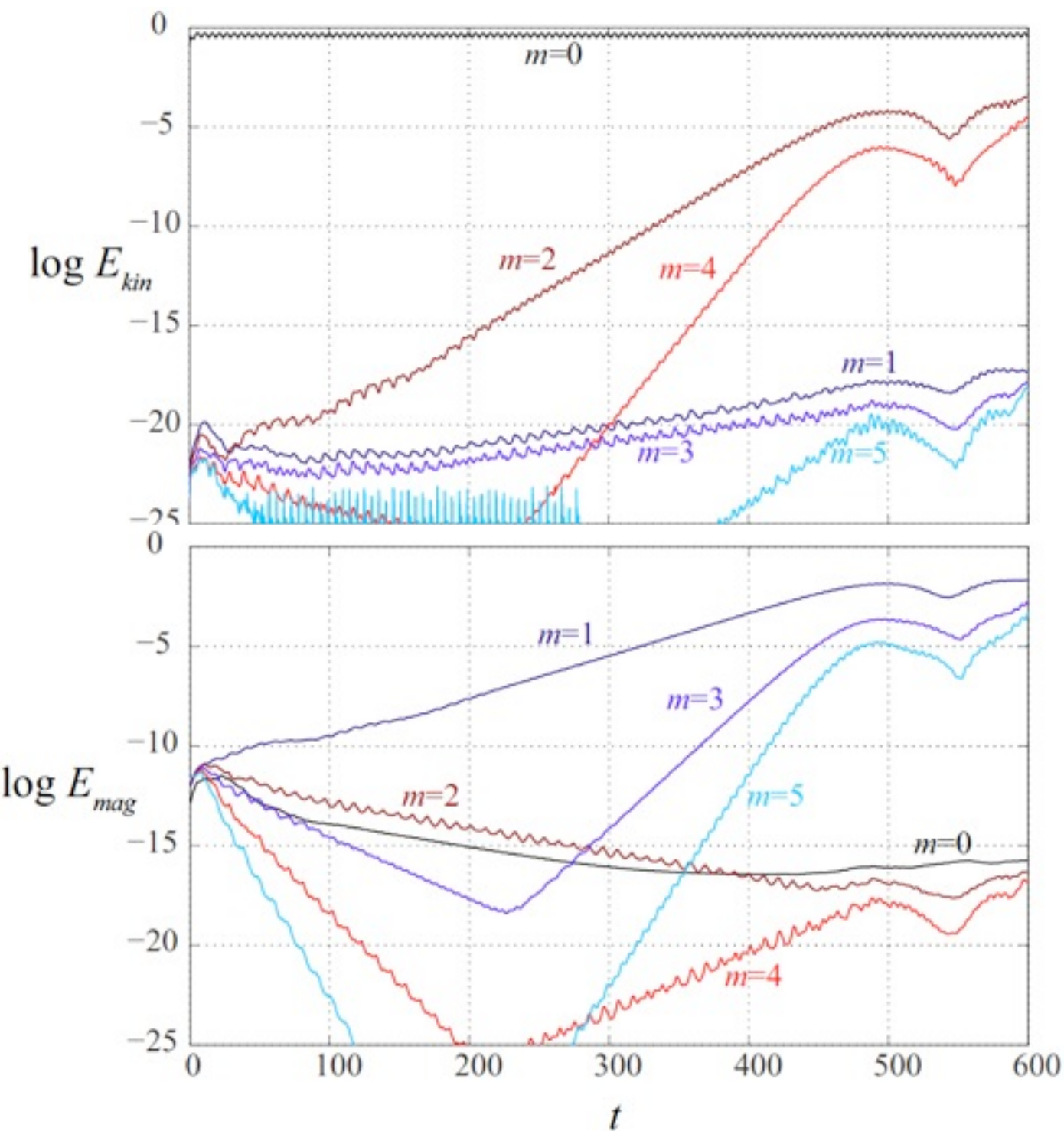
$a = 1, \omega = 0.6, Re = 200$



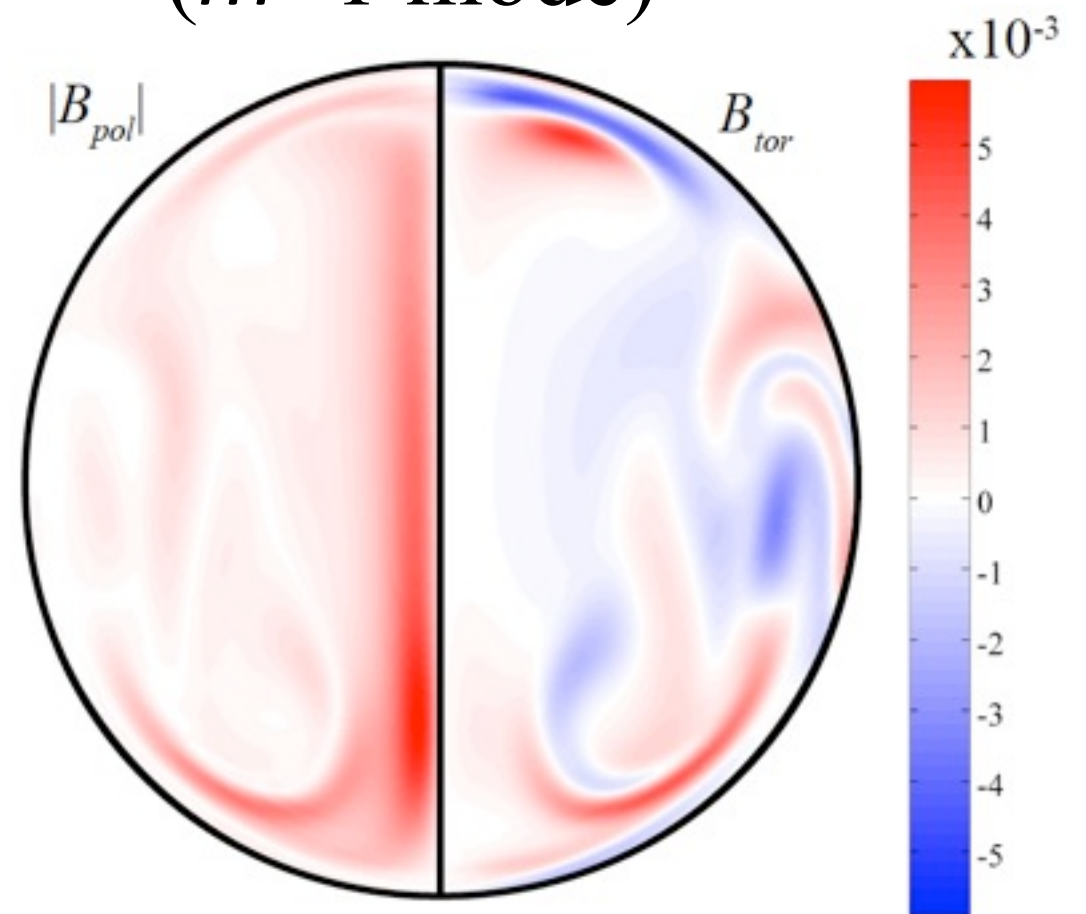
- Floquet Theory Eigenmode code (Khalzov)

Small-scale dynamo

Energies of modes vs. time

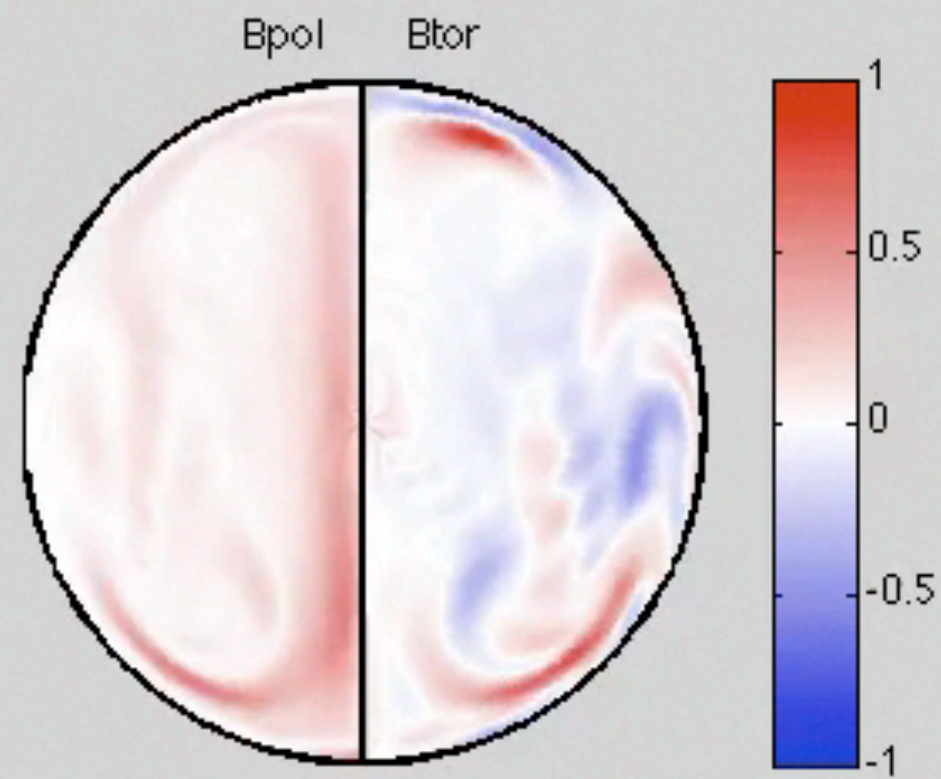
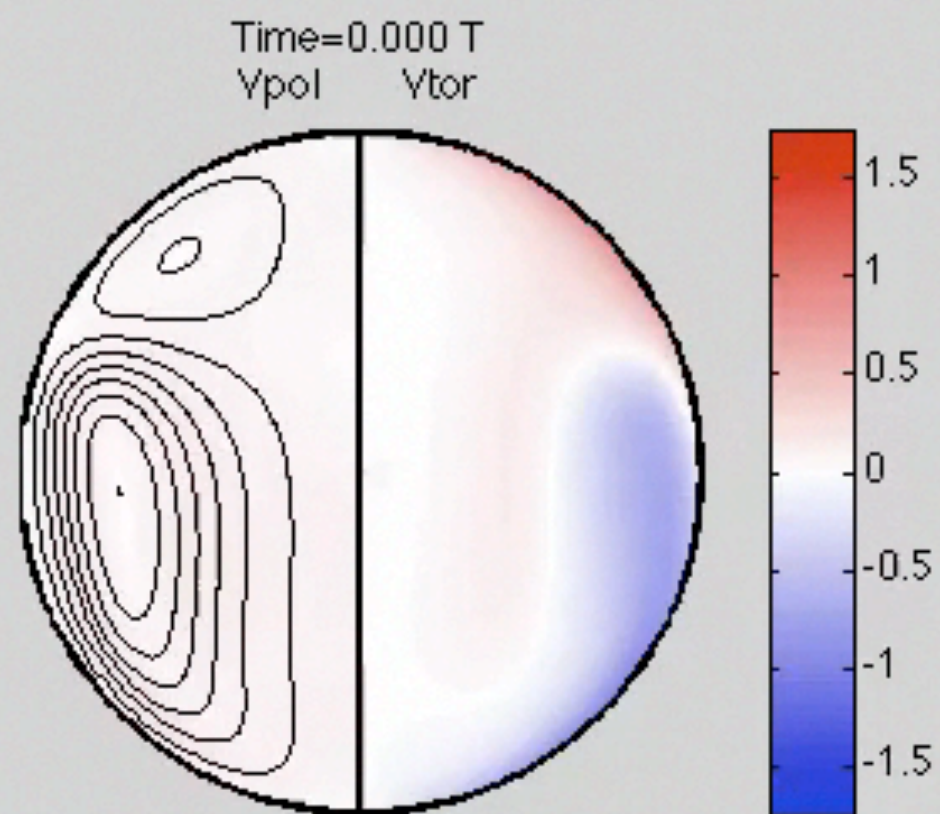


Dynamo field at $t=300$
($m=1$ mode)



$$V_{\phi}(\theta, t) = \sin 2\theta + a \sin \theta \sin \omega t$$

$$a = 1, \omega = 0.6, Re = 200, Rm = 5000$$



Fast dynamo eigen-mode structure

($Re=200$, $Rm=5000$, $m=1$)

- Dynamics during one period $T=2\pi/\omega=10.47$ (unit of time: $t_0=R_0/V_0$)

